
Parameters Tuning of Hadronic Generator LUARLW

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Abstract The measurements of R value needs Monte Carlo hadronization generator to determine the hadronic efficiency. The generator LUARLW based on the Lund area law can simulate the productions and decays of the continuous and resonant hadronic states from $2m_\pi$ up to 5 GeV. The related phenomenological parameter set in LUARLW are tuned using the data samples between 2.0-4.6 GeV taken at BESIII. The some important distributions of final states particles of data and MC are compared at the detector level. The estimated systematic error of hadronic efficiency are about 2%.

Key words Monte Carlo generator, Lund area law, hadronic efficiency.

1 Motivation

The directly measured hadronic cross section from the data analysis is the so-called observed cross section written in following formula

$$\sigma_{obs}^{tot}(s) = \frac{N_{obs} - N_{bg}}{L} \quad (1)$$

where N_{obs} is the number of hadronic event selected from the data sample with the integrated luminosity L , and N_{bg} the number of the remnant backgrounds. the quantity of $\sigma_{obs}^{tot}(s)$ is experiment dependent and is not physics cross section. If one may know the efficiency factor $\bar{\epsilon}$ which reflect the ratio that of number of hadronic event produced at the collision vertex and the number of the observed hadronic event after raw data reconstruction and event selection, the quantity presented by

$$\sigma^{tot}(s) \equiv \sigma_{phys}^{tot}(s) = \frac{N_{obs} - N_{bg}}{L\bar{\epsilon}} \quad (2)$$

is the physics and experiment independent cross section which corresponds to the value calculated in theoretic calculation. In principle, the hadronic efficiency $\bar{\epsilon}$ is unknown in data analysis, but it can be estimated by the Monte Carlo method. The Monte Carlo consists two parts, the simulation of hadronic production at collision vertex and all processes of the long-life particles in detector. It is always assume that the detector simulations are perfect, and the difference between data and MC simulations are caused by the faultiness of event generator, and the coincidence between data and MC can be improved by tuning the phenomenological parameters of the event generator. In the experimental measurements of R value at BESII, the hadronic generator LUARLW which based on the Lund string fragmentation area law was used and the parameters has been tuned, the combined systematic errors of event selection and hadronic detection efficiency were estimated as about 2%. But in BESIII R value measurement, the detector simulation adopts the GEANT4 scheme instead of the BGS scheme in BESII. So the parameters in LUARLW need to tune again using BESIII data. In addition, the LUARLW itself has also some improvements.

2 The Lund string fragmentation hadronization model

In the modern point of view, the hadron production in the high energy electron-positron collision may be described as the following processes depicted in Fig.1: the electron and positron (e^+ and e^-) annihilate into a

virtual photo γ^* (Z^0 can be neglected at BEPC energies), then the γ^* splits into a pair of initial quark-antiquark pairs $q\bar{q}$, and follows a series of perturbative QCD cascade evolutions of quarks and gluons (the parton shower), finally the preliminary hadrons produce via nonperturbative hadronization processes, the unstable hadrons (if have) will decay into the stable final state. The inclusive hadronization processes (the stage represented by the gray region in the Fig.1) are usually described by phenomenological models.

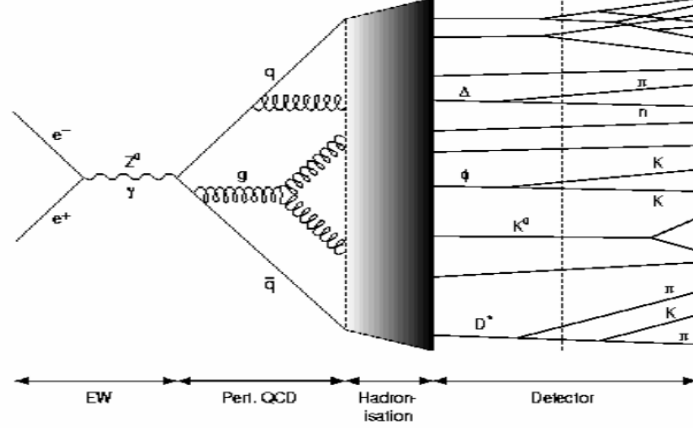


Fig.1 The hadron production processes via virtual photo annihilation in the electron-positron collision.

In the R value measurement at BESII and BESIII, the Lund string fragmentation model is studied and used. The well-known Monte Carlo simulation packet JETSET is not built in order to describe few-body states (in particular at the few GeV level in e^+e^- annihilation as in BEPC energy region). For the purpose of R value measurement with BEPC/BES, the formalism of the Lund model was tried to solve strictly starting from the basic assumptions (relativity, quantum mechanics, causality) and did not adapt the extreme high energy approximation $s \rightarrow \infty$ in the derivation of the formula like did in JETSET, the obtained form of the string fragmentation function is the Lund area law, instead of the fragmentation function $f(z)$ which is used in JETSET. So the Monte Carlo packet LUARLW is an hadronic generator based on the Lund Model area law directly^[2].

In general, any hadronic generator contains a few phenomenological parameters which have definite physical meaning reflecting the dynamical mechanisms of the hadronization, but their values are unknown and need to tune using the experimental data. The tactic to tune these parameters is to choose a set of distributions of the hadronic final states (such as, the charged and neutral multiplicities, the fractions of the all preliminary and final hadrons produced in the simulation, momentum, etc.), compare the corresponding distributions between data and generator at detector level and tune the related values of the parameters iteratively, untill all the chosen distributions of data and MC are acceptable agreement. In the parameters tuning, one always assume that the detector simulations are perfect, and any difference between data and simulation are due to the shortcoming of the hadronization model and/or the values of the parameters are unreasonable.

In the R value measurement with the BESII and BESIII data, the generator LUARLW is used. Historically, the LUARLW was ever tuned twice. The first time, the data samples with small statistics (about 1000 hadronic events for each energy point) between 2-5 GeV at about 100 energy points taken in 1999 are used, the error of the hadronic efficiency ϵ_{had} are about $2 \sim 5\%$ ^[3, 4]. Later on, LUARLW was tuned again using the data samples with larger statistics taken at four energy points below open charm threshold collected in 2004, these large data samples provided more information and small statistic fluctuation, which are helpful to tune the parameters of LUARLW effectively, the error of the ϵ_{had} decreased to about 3% ^[5]. Now, the higher luminosity of BEPCII provides data samples with even large statistics, and the better performance of BESIII detector and

more accurate detector simulations with GEANT4 will be helpful to tune the parameters of LUARLW.

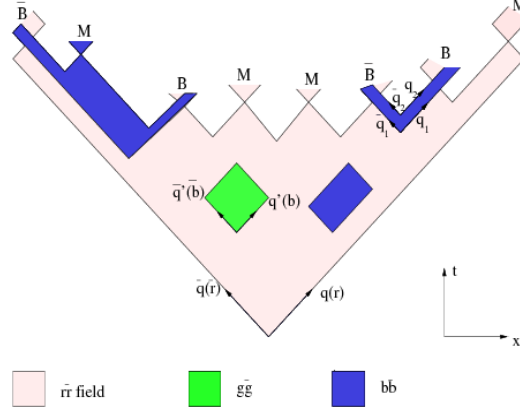


Fig.2 The hadron production processes in string fragmentation.

3 The formalism of Lund area law

The inclusive hadronic generator LUARLW are used for BEPCII/BESIII experiments, the nominal center of mass energy E_{cm} range is 2-4.6 GeV, due to the bremsstrahlung of the initial beam, the effective energy E'_{cm} may decrease to the lowest hadronic production threshold (i.e. the state $\pi^+\pi^-$). The resonances with quantum number $J^{PC}=1^{--}$ will be produced except for the continuous states.

In BEPC/BESII running period, the processes can be simulated by LUARLW are not complete, they are described in Ref.[4]^[3] and listed below

In Lund area law, the lowest cross section for an exclusive process $e^+e^- \rightarrow \text{string} \rightarrow m_1+m_2\ldots+m_n$ can be written as

$$d\sigma(e^+e^- \rightarrow m_1, m_2 \ldots m_n) = d\sigma(e^+e^- \rightarrow q\bar{q}) \cdot d\mathcal{P}(q\bar{q} \rightarrow m_1, m_2 \ldots m_n; s), \quad (3)$$

the differential cross section for $e^+e^- \rightarrow q\bar{q}$ can use QED expression

$$d\sigma(e^+e^- \rightarrow q\bar{q}) = N_c \frac{\alpha^2}{4s} \sum_q \beta_q [1 + \cos^2\theta + (1 - \beta_q^2) \sin^2\theta] = N_c \frac{\alpha^2}{4s} \sum_q \beta_q [(2 - \beta_q^2) + \beta_q^2 \cos^2\theta], \quad (4)$$

and the probability for string fragmentation into n hadrons reads

$$d\mathcal{P}(q\bar{q} \rightarrow m_1, m_2 \ldots m_n; s) = \delta(1 - \sum_{j=1}^n \frac{m_{\perp j}^2}{sz_j}) \delta(1 - \sum_{j=1}^n z_j) \delta^{(2)}(\sum_{j=1}^n \vec{k}_j) \overline{\sum} |\mathcal{T}^{(n)}|^2 d\Phi_n, \quad d\Phi_n = \prod_{j=1}^n d^2\vec{k}_j \frac{dz_j}{z_j} \quad (5)$$

where, the factor like the element of matrix in Lund area law can be written as

$$\mathcal{T}^{(n)} = N^n \mathcal{T}_{\perp}^{(n)} \cdot \mathcal{T}_{//}^{(n)} \quad (6)$$

the part corresponding to the transverse momentum has the Gaussiann form

$$\mathcal{T}_{\perp}^{(n)} = \exp(-\sum_{j=1}^n \vec{k}_j^2) \quad \vec{k}_j = \frac{\vec{p}_{j\perp}}{2\sigma}, \quad (7)$$

and the part corresponding to the longitudinal momentum takes the Lund area law

$$\mathcal{T}_{//}^{(n)} = \exp(i\xi \mathcal{A}_n) \quad \xi = \frac{1}{2\kappa} + i\frac{b}{2} \quad (8)$$

where $\mathcal{A}_n$ is the light-cone area enclosed by the light-cone of the n hadrons.

Starting from the Lund area law, one may obtain a approximation expression of a poisson-like multiplicity distribution of the fragmentation hadron

$$P_n(s) = \frac{\mu^n}{n!} \exp[c_0 + c_1(n - \mu) + c_2(n - \mu)^2], \quad (9)$$

where, n is the number of the fragmentation hadron, μ can be understood as the average multiplicity which can quote the QCD prediction

$$\mu = \alpha + \beta \exp(\gamma \sqrt{s}), \quad (10)$$

where, $c_1, c_2, c_3, \alpha, \beta$ and γ are free parameters and to be tuned by experimental data.

For the following string fragmentation and n hadrons production process

$$e^+e^- \Rightarrow q\bar{q} \Rightarrow \text{string} \Rightarrow m_1 + m_2 + \cdots + m_n,$$

the total matrix element can be written as

$$\mathcal{M} \equiv \mathcal{M}_{\text{QED}}(e^+e^- \rightarrow q\bar{q}) \mathcal{M}_{\text{LUND}}(q\bar{q} \rightarrow \text{string} \rightarrow m_1, m_2, \cdots m_n), \quad (11)$$

where $\mathcal{M}_{\text{QED}}$ is for the pure QED process $e^+e^- \rightarrow q\bar{q}$, and gives the differential cross section

$$\frac{d\sigma(e^+e^- \rightarrow q\bar{q})}{d\Omega_{q\bar{q}}} = \frac{3\alpha^2}{4s} e_q^2 \beta [1 + \cos^2\theta + (1 - \beta^2)\sin^2\theta]. \quad (12)$$

The matrix element for the string fragmentation into n hadrons can be factorized as two factors

$$\mathcal{M}_{\text{LUND}}(q\bar{q} \rightarrow \text{string} \rightarrow m_1, m_2, \cdots m_n) = C_n \mathcal{M}_\perp \mathcal{M}_{//}, \quad (13)$$

where $\mathcal{M}_\perp$ for the transverse momentum of the fragmentation hadrons and has the Gaussian form

$$\mathcal{M}_\perp = \exp(-\sum_{j=1}^n \mathbf{k}_j^2), \quad (14)$$

and $\mathcal{M}_{//}$ for the longitudinal momentum has the Lund area law form

$$\mathcal{M}_{//} = \exp(i\xi \mathcal{A}_n), \quad (15)$$

where $\mathcal{A}_n$ is the light-cone area in Fig.2, and ξ is the complex parameter reflecting the dynamics properties of the string fragmentation

$$\xi = \frac{1}{2\kappa} + i\frac{b}{2}. \quad (16)$$

The probability of n hadrons in the element of phase space $d\Phi_n$ is

$$d\mathcal{P}(q\bar{q} \rightarrow m_1, m_2, \cdots m_n) = (2\pi)^4 \delta(1 - \sum_{j=1}^n \frac{m_{\perp j}^2}{sz_j}) \delta(1 - \sum_{j=1}^n z_j) \delta^2(\sum_{j=1}^n \mathbf{k}_j) \sum |\mathcal{M}_{\text{LUND}}|^2 d\Phi_n, \quad (17)$$

where

$$d\Phi_n = \prod_{j=1}^n d^2\mathbf{k}_j \frac{dz_j}{z_j}, \quad (18)$$

and the transverse mass of hadron j with mass m_j is

$$m_{\perp j} = \sqrt{m_j^2 + p_{\perp j}^2} = \sqrt{m_j^2 + 4\sigma_\perp^2 \mathbf{k}_j^2}, \quad (j=1, 2, \cdots n). \quad (19)$$

The probability given in Eq.(17) can be solved for $n=2, 3 \cdots 6$ respectively^[2], which determine the momentum distribution of the n preliminary hadrons with masses $(m_1, m_2 \cdots m_n)$ in the string fragmentation. Finally, the total cross section of n hadrons production in the Lund area law reads

$$\sigma(e^+e^- \rightarrow m_1, m_2, \cdots m_n) = \int d\Omega_{q\bar{q}} \frac{d\sigma(e^+e^- \rightarrow q\bar{q})}{d\Omega_{q\bar{q}}} \cdot d\mathcal{P}_n(q\bar{q} \rightarrow \text{string} \rightarrow m_1, m_2, \cdots m_n). \quad (20)$$

$$Z_n = s \int d\Phi_n \exp(-b\mathcal{A}_n), \quad (21)$$

$$P_n = Z_n / \sum Z_n, \quad (22)$$

In the Monte Carlo simulation of an inclusive hadronic sample, the multiplicity of the preliminary hadrons is determined by following distribution

$$P_n(s) = \frac{\mu^n}{n!} \exp[c_0 + c_1(n - \mu) + c_2(n - \mu)^2]. \quad (23)$$

where could be understood as the average multiplicity, and expressed as the form predicted by QCD

$$\mu = \alpha + \beta \exp(\gamma \sqrt{s}), \quad (24)$$

the values of the free parameters c_0 , c_1 , c_2 , α , β and γ can be obtained by match the data.

4 Functions of LUARLW

The inclusive hadronic generator LUARLW are used for BEPCII/BESIII experiments, the nominal center of mass energy E_{cm} range can be 2-4.6 GeV, due to the bremsstrahlung of the initial beam, the effective energy E'_{cm} may decrease to the lowest hadronic production threshold (i.e. the state $\pi^+\pi^-$). The resonances with quantum number $J^{PC} = 1^{--}$ will be produced except for the continuous states. In addition, the leading order QCD correction is also considered. In order to simulate the hadronic production processes more completely, LUARLW are extended and improved compared to the version used in BESII work. The basic processes can be simulated by LUARLW are summarized below.

$$e^+e^- \Rightarrow \gamma^* \Rightarrow \begin{cases} q\bar{q} \Rightarrow \text{string} \Rightarrow \text{hadrons} \\ gq\bar{q} \Rightarrow \text{string} + \text{string} \Rightarrow \text{hadrons} \end{cases}$$

The vector mesons with $J^{PC} = 1^{--}$ can directly couple to virtual photon and production in ISR return process

$$e^+e^- \rightarrow \gamma^* \rightarrow \rho(770), \omega(782), \phi(1020) \cdots \rho(1700). \quad (25)$$

The production and decay of the charmonium adopt the standard pictures^[4].

$$e^+e^- \Rightarrow \gamma^* \Rightarrow J/\psi \Rightarrow \begin{cases} \gamma^* \Rightarrow e^+e^-, \mu^+\mu^- \\ \gamma^* \Rightarrow \bar{q}q \Rightarrow \text{string} \Rightarrow \text{hadrons} \\ ggg \Rightarrow \text{string} + \text{string} + \text{string} \Rightarrow \text{hadrons} \\ \gamma gg \Rightarrow \text{string} + \text{string} \Rightarrow \text{hadrons} \\ \gamma \eta_c \Rightarrow gg \Rightarrow \text{string} + \text{string} \Rightarrow \text{hadrons} \\ \gamma + \text{exclusive radiative decay channels} \end{cases}$$

$$e^+e^- \Rightarrow \gamma^* \Rightarrow \psi(3686) \Rightarrow \begin{cases} \gamma^* \Rightarrow e^+e^-, \mu^+\mu^-, \tau^+\tau^- \\ \gamma^* \Rightarrow q\bar{q} \Rightarrow \text{string} \Rightarrow \text{hadrons} \\ ggg \Rightarrow \text{string} + \text{string} + \text{string} \Rightarrow \text{hadrons} \\ \gamma gg \Rightarrow \text{string} + \text{string} \Rightarrow \text{hadrons} \\ \gamma \eta_c \Rightarrow gg \Rightarrow \text{string} + \text{string} \Rightarrow \text{hadrons} \\ \gamma + \text{exclusive radiative decay channels} \\ \pi\pi J/\psi, \eta J/\psi, \pi^0 J/\psi \end{cases}$$

$$e^+e^- \Rightarrow \gamma^* \Rightarrow \psi(3770) \Rightarrow \left\{ \begin{array}{l} \gamma^* \Rightarrow e^+e^-, \mu^+\mu^-, \tau^+\tau^- \\ \gamma^* \Rightarrow q\bar{q} \Rightarrow \text{string} \Rightarrow \text{hadrons} \\ ggg \Rightarrow \text{string} + \text{string} + \text{string} \Rightarrow \text{hadrons} \\ \gamma gg \Rightarrow \text{string} + \text{string} \Rightarrow \text{hadrons} \\ \pi\pi J/\psi, \eta J/\psi, \pi^0 J/\psi \\ \gamma + \text{exclusive radiative decay channels} \\ D\bar{D} \end{array} \right.$$

$$e^+e^- \Rightarrow \gamma^* \Rightarrow \left\{ \begin{array}{l} \psi(4040) \Rightarrow D\bar{D}, D^*\bar{D}^*, D\bar{D}^*, D^*\bar{D}, D_s\bar{D}_s, \text{other decay modes} \\ \psi(4160) \Rightarrow D\bar{D}, D^*\bar{D}^*, D\bar{D}^*, D^*\bar{D}, D_s\bar{D}_s, D_s\bar{D}_s^*, \text{other decay modes} \\ \psi(4415) \Rightarrow D\bar{D}, D^*\bar{D}^*, D\bar{D}^*, D^*\bar{D}, D_s\bar{D}_s, D_s\bar{D}_s^*, D_s^*\bar{D}_s^*, \text{other decay modes} \\ X(4260) \Rightarrow \text{possible decay modes} \\ X(4360) \Rightarrow \text{possible decay modes} \end{array} \right.$$

5 Parameters tuning

The generator LUARLW is used to determine the hadronic efficiency in R value measurement. In BEPCII energy region, R values are measured with inclusive scheme, so it is enough that LUARLW agree with data inclusively, no need request it agree with data exclusively. There are some phenomenological parameters in LUARLW. In the BEPC energy, the main parameters are those which determine the ratios of mesons and baryons with different quantum number (s, L, J) in the string fragmentation process. In LUARLW these parameters are stored in array $PARJ(1-20)$ as did in JETSET^[6], their default values were set by fitting the data collected at LEP^[7]. The values of $PARJ(1-20)$ are tuned to make the MC agree with the experimental data taken with BESIII. It is found that some parameters are not constants in BEPCII energy region, some experiential and smooth functions (no physics bases, just for fit data well) are adopted to give the values of parameters. Table 1 lists the tuned parameters and their tuned values.

In the processes of LUARLW's parameters tuning, one may find the earlier version of LUARLW code can not agree data well by tuning the parameters only. At this issue, one has to add some new mechanism and functions, i.e, modify the code. As the preliminary stage of parameters tuning, the parameters are tuned by hand, and compare the MC with data. After many many time manually tuning, MC agree with data for most of the important distributions generally, which mean the values of the related parameters of LUARLW are generally reasonable. When arriving at this stage, the parameters of LUARLW should be optimized by fitting all distributions again by the way described in reference ^[7].

6 Comparison of data and MC

The main principle of parameters tuning is requesting MC agree data well as many as possible for those final state distributions which related to the event selection criteria, especially for those distributions which are sensitive to the efficiency, such as charged multiplicity, polar angle $\cos\theta$, charged track momentum, deposit energy of charged and neutral tacks in EMC, angle between isolate photon and charged track for 2-prong events. The single particle (mesons and baryons) fraction in final states are also compared, which reflect whether the tuned values of parameter sets $PARJ(1-17)$ are reasonable or not. It is noticed that the generator LUARLW is not perfect yet, which certainly can not describe entire and intricate hadron production very well, it needs further modification when simply tuning of parameters can not improve the agreement between MC and data.

Table 1: The parameters tuned with BESIII data.

parameter	default	tuned	meaning
c_0	-	5.0	parameter in preliminary hadron multiplicity distribution $P_n(s)$
c_1	-	0.05	parameter in preliminary hadron multiplicity distribution $P_n(s)$
c_2	-	-0.25	parameter in preliminary hadron multiplicity distribution $P_n(s)$
α	-	0.80	parameter α in $\mu = \alpha + \beta \exp(\gamma \sqrt{s})$
β	-	0.25	parameter β in $\mu = \alpha + \beta \exp(\gamma \sqrt{s})$
γ	-	1.10	parameter γ in $\mu = \alpha + \beta \exp(\gamma \sqrt{s})$
$\sigma_{\perp}$	-	Ecm-dependent	effective transverse momentum width in like-Gaussian form
PARJ(01)	0.10	Ecm-dependent	diquark/quark production ratio, baryon suppression (B/M)
PARJ(02)	0.30	Ecm-dependent	$s/(u,d)$ production ratio, strange meson suppression (K/π)
PARJ(03)	0.40	Ecm-dependent	strange diquark suppression, strange baryon suppression (Λ/p)
PARJ(04)	0.05	0.05	suppression of spin 1 diquark compared to spin 0 ones
PARJ(05)	0.50	0.50	relative occurrence of baryon produced by $BM\bar{B}$ and by $B\bar{B}$
PARJ(06)	0.50	0.50	suppression for having $s\bar{s}$ shared by B and $\bar{B}$ of $BM\bar{B}$ situation
PARJ(07)	0.50	0.50	suppression for having strange meson M in $BM\bar{B}$ configuration
PARJ(11)	0.50	Ecm-dependent	suppression of light meson has spin 1 compared to spin 0 (ρ/π)
PARJ(12)	0.60	0.60	suppression of strange meson of spin 1 compared to spin 0 (K^*/K)
PARJ(13)	0.75	0.75	suppression of charm meson of spin 1 compared to spin 0 (D^*/D)
PARJ(14)	0.00	0.09	probability that spin $s=0$ and orbital $L=1$ with total $J=1$ meson
PARJ(15)	0.00	0.07	probability that spin $s=1$ and orbital $L=1$ with total $J=0$ meson
PARJ(16)	0.00	0.08	probability that spin $s=1$ and orbital $L=1$ with total $J=1$ meson
PARJ(17)	0.00	0.10	probability that spin $s=1$ and orbital $L=1$ with total $J=2$ meson

A objective standard that the values of the tuned parameters are optimized are to see the quantities χ^2 about the differences between MC data for above mentioned distributions and $\Delta(N_{had}/\bar{\epsilon})$ are acceptable small.

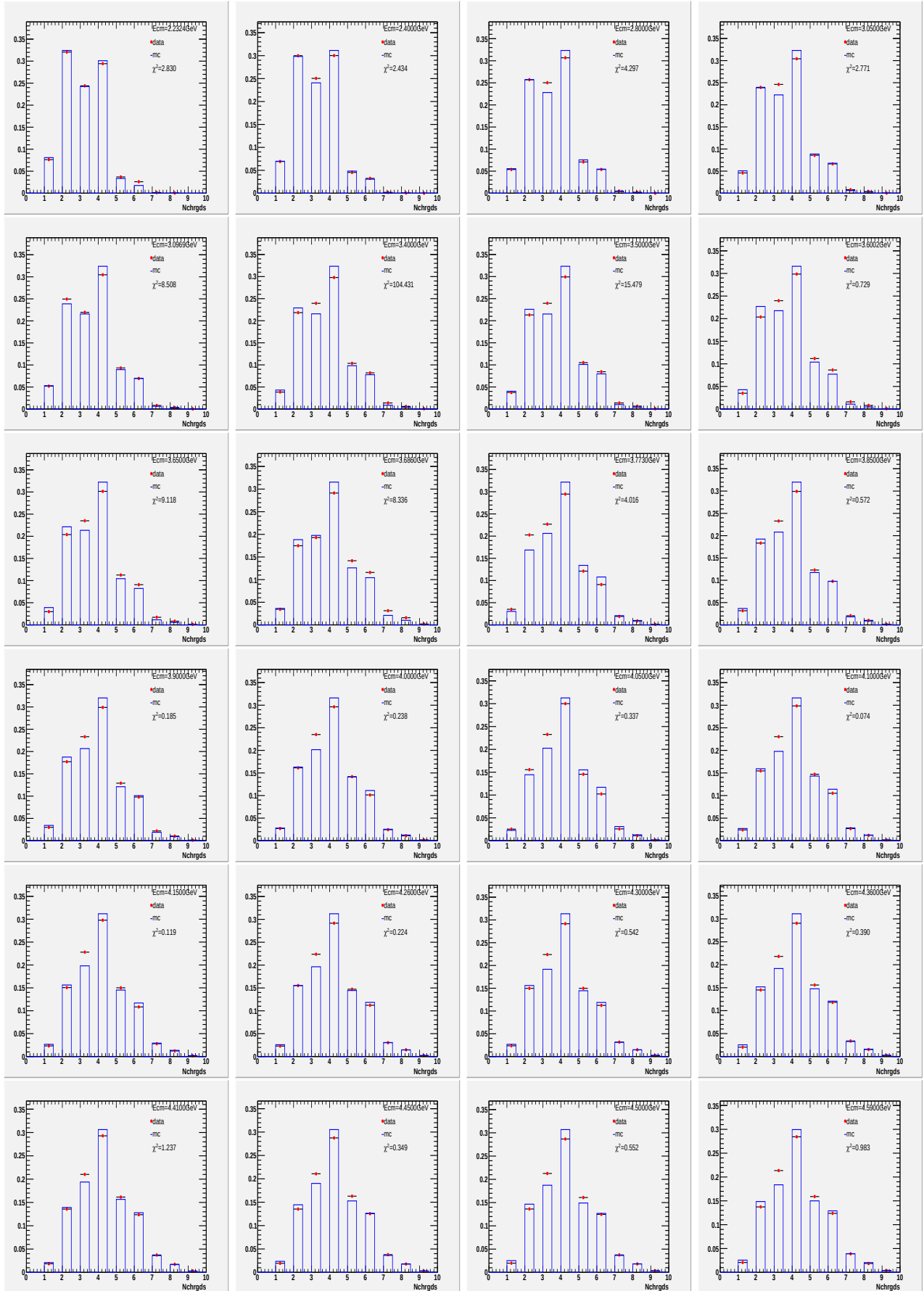


Fig. The multiplicity of charged track distribution.

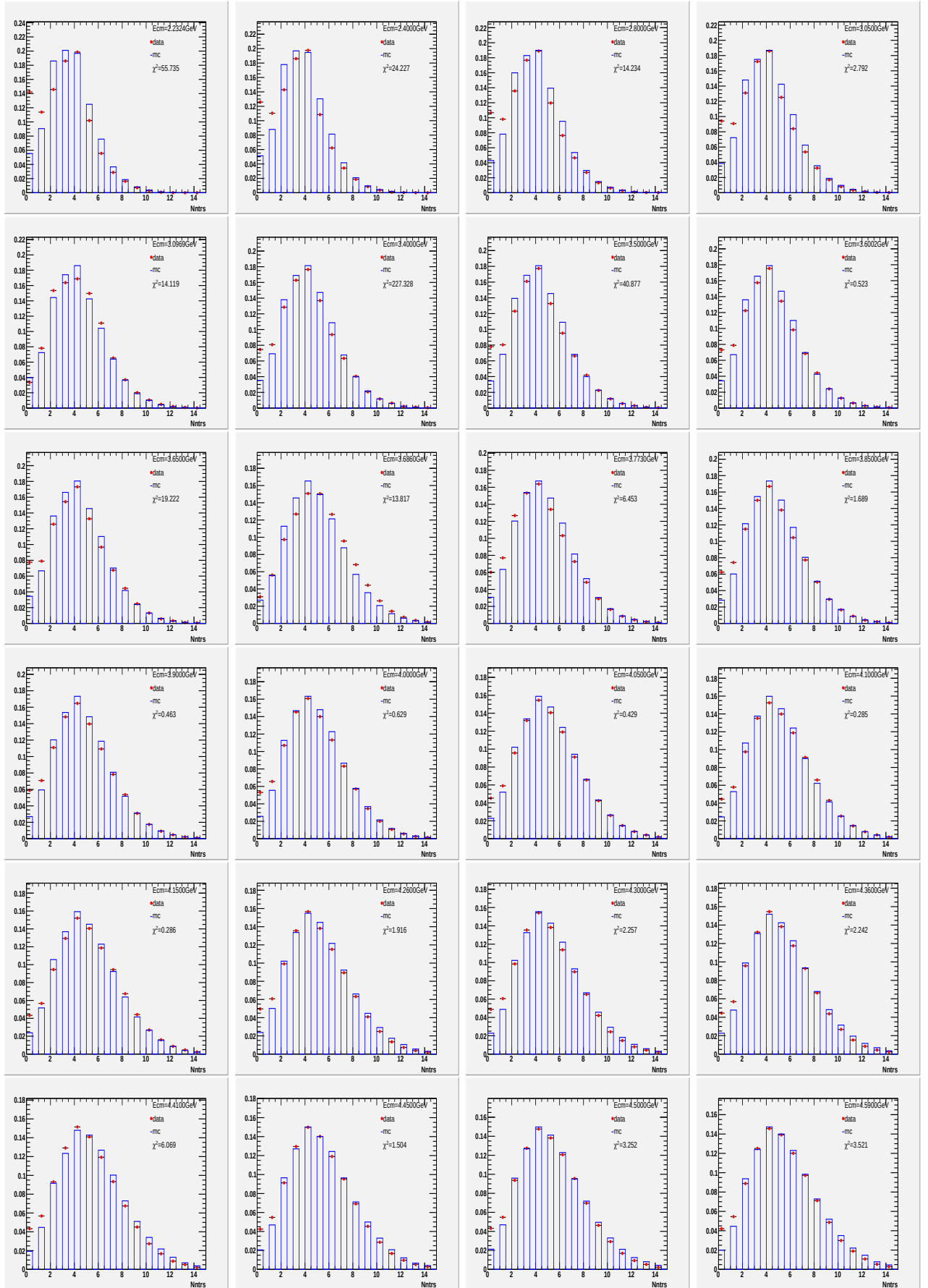


Fig. The multiplicity of neutral track distribution.

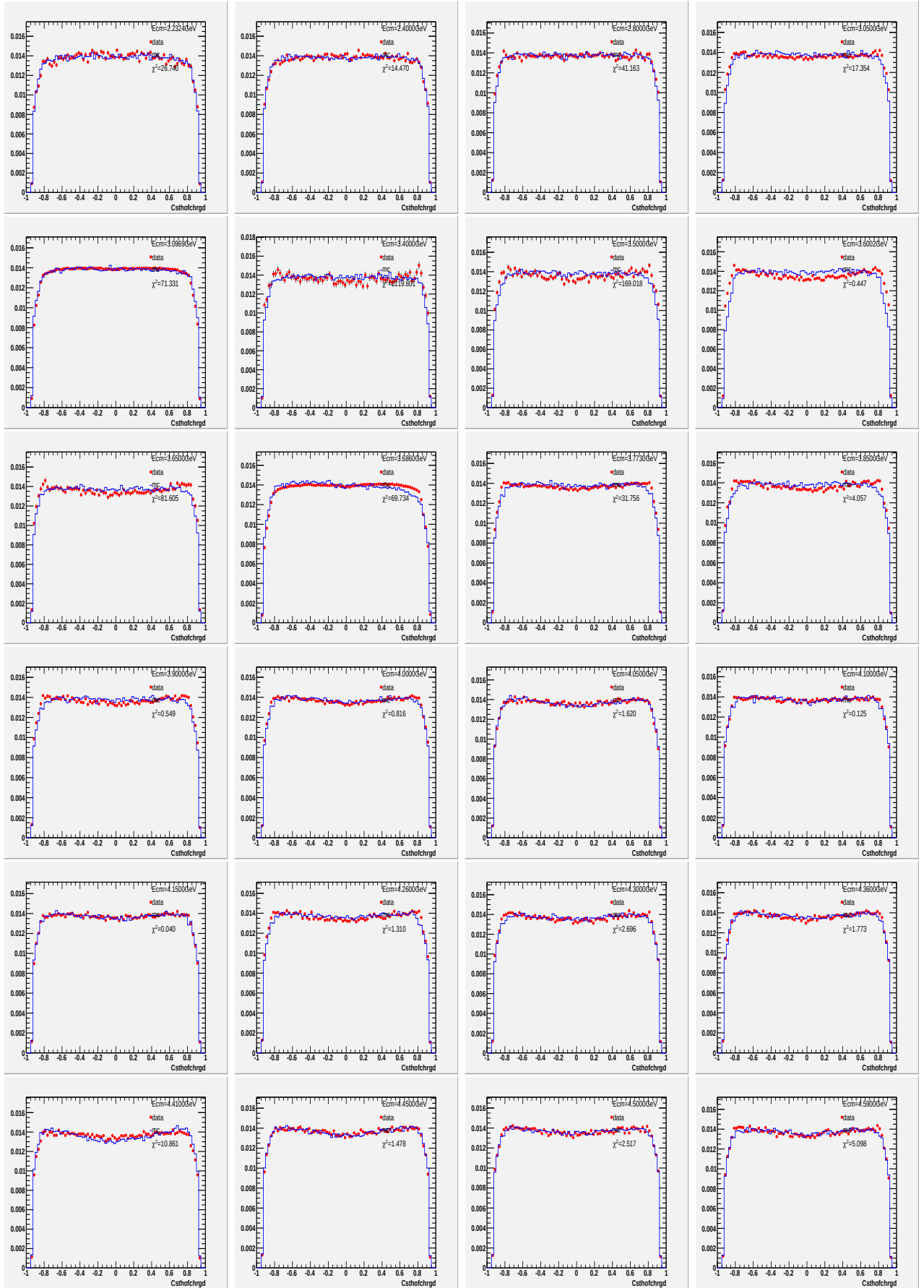


Fig. The polar angle $\cos\theta$ of charged track distribution.

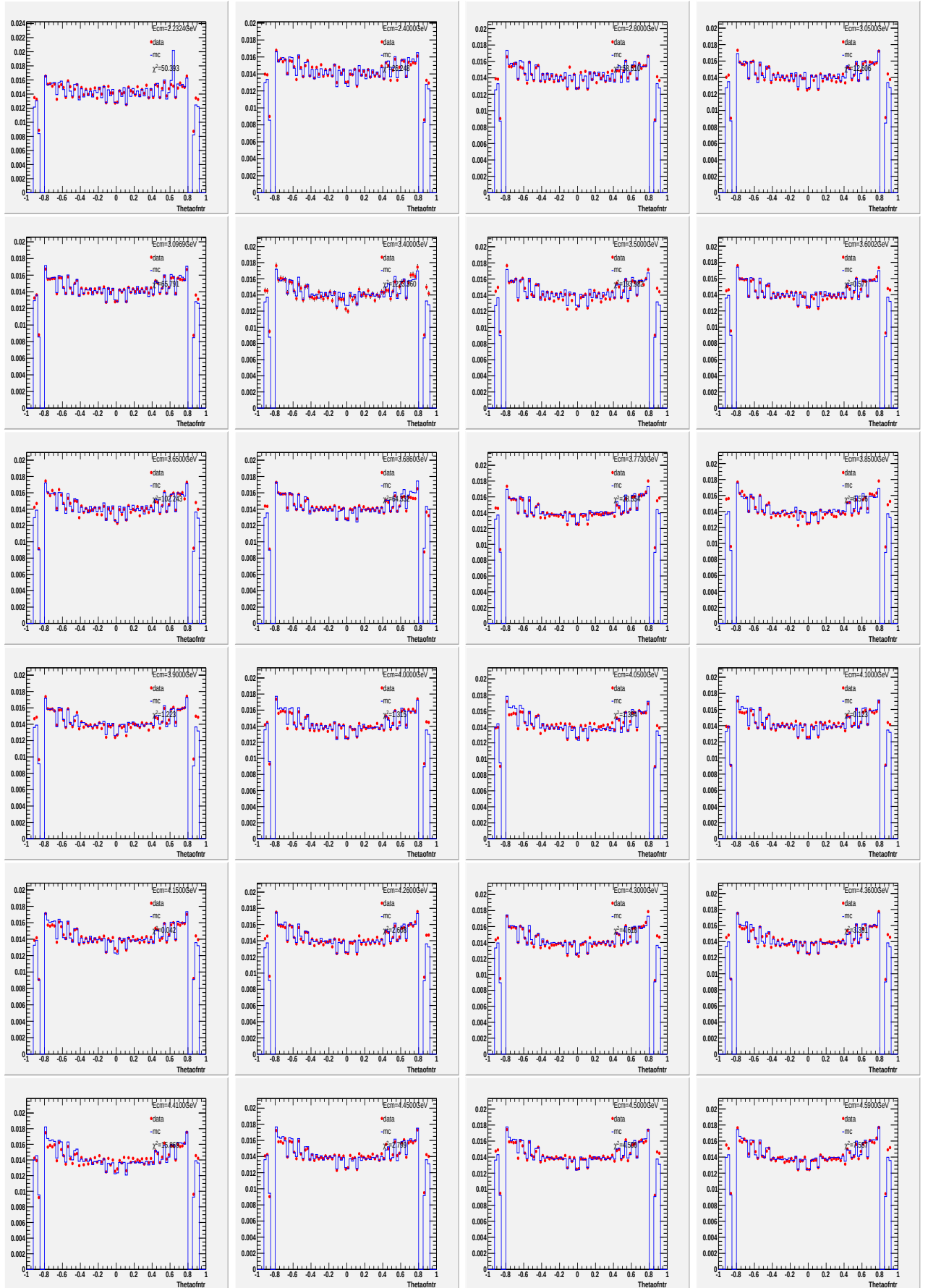


Fig. $\cos\theta$ of neutral tracks.

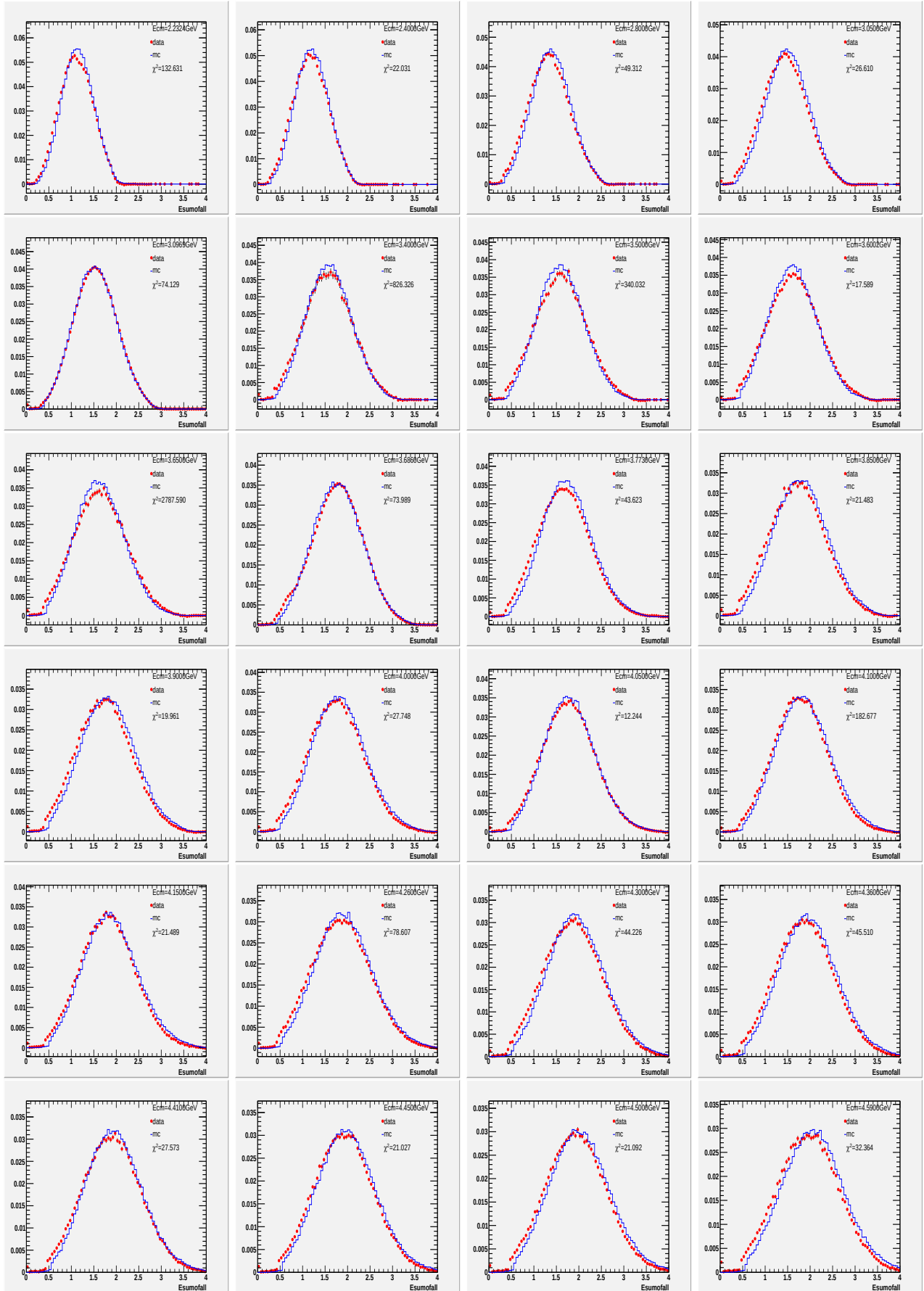


Fig. The deposit energy in ECM of charged and neutral tracks.

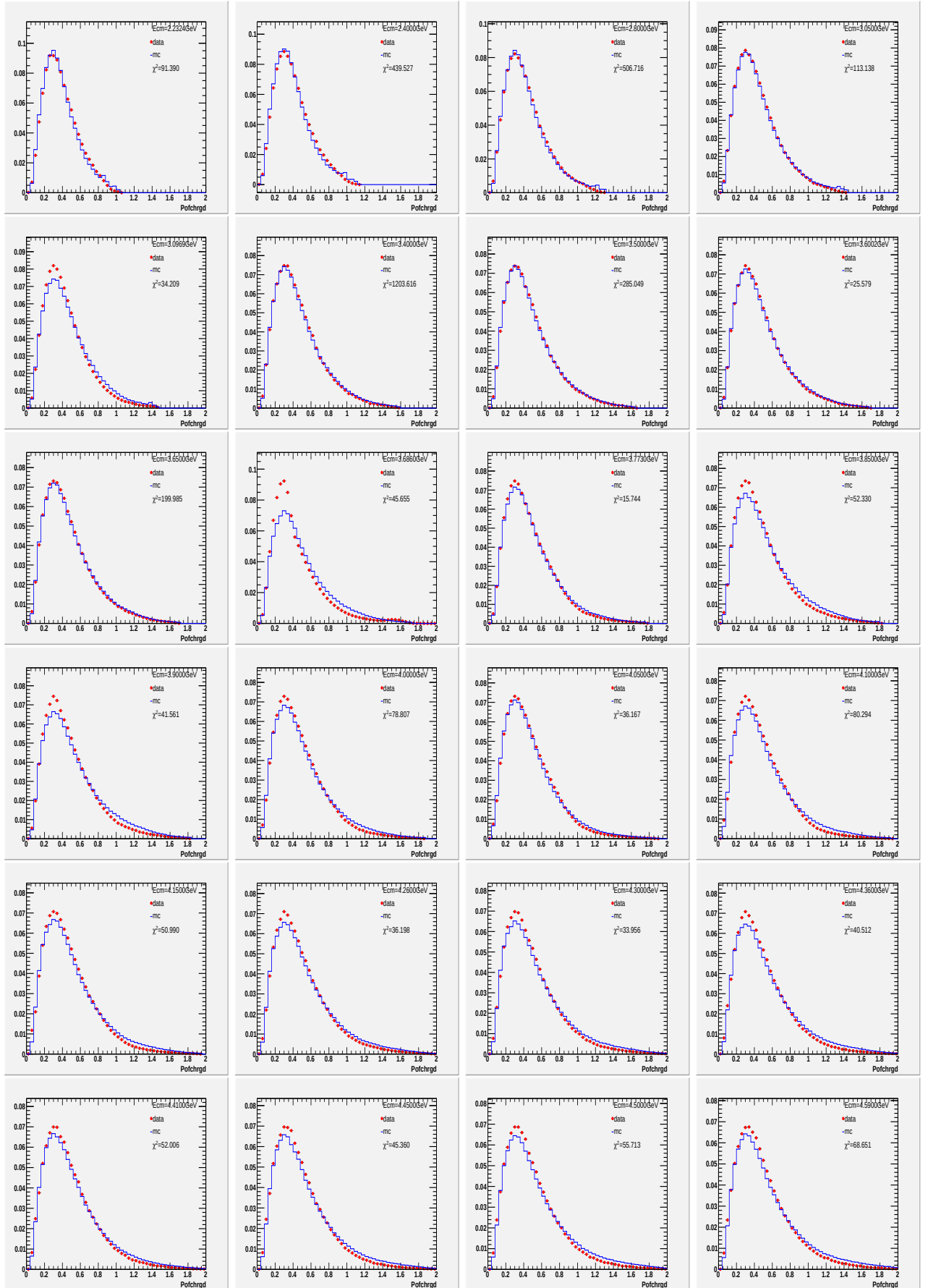


Fig. The momentum of charged tracks.

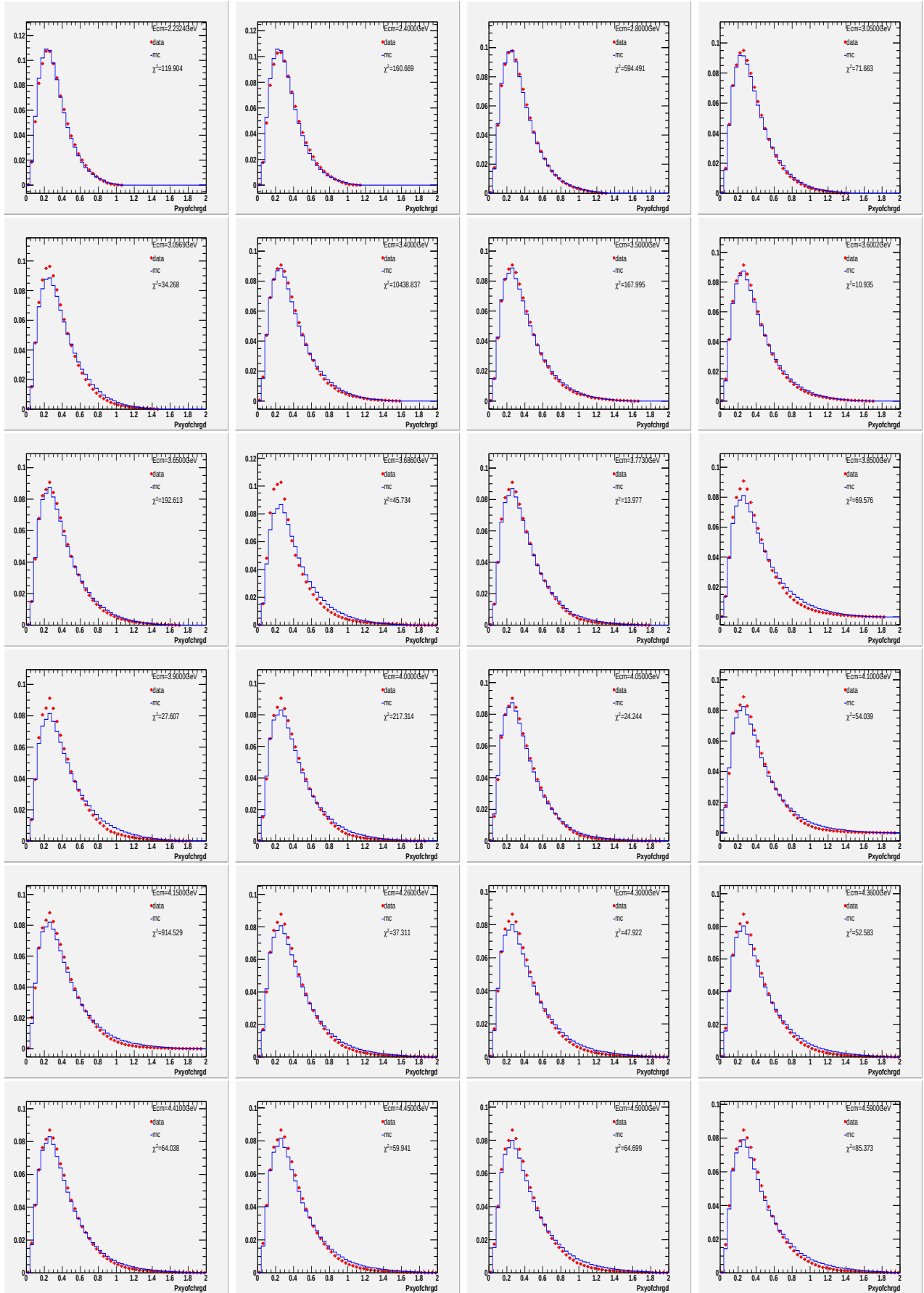


Fig. The transverse momentum P_{xy} of charged tracks.

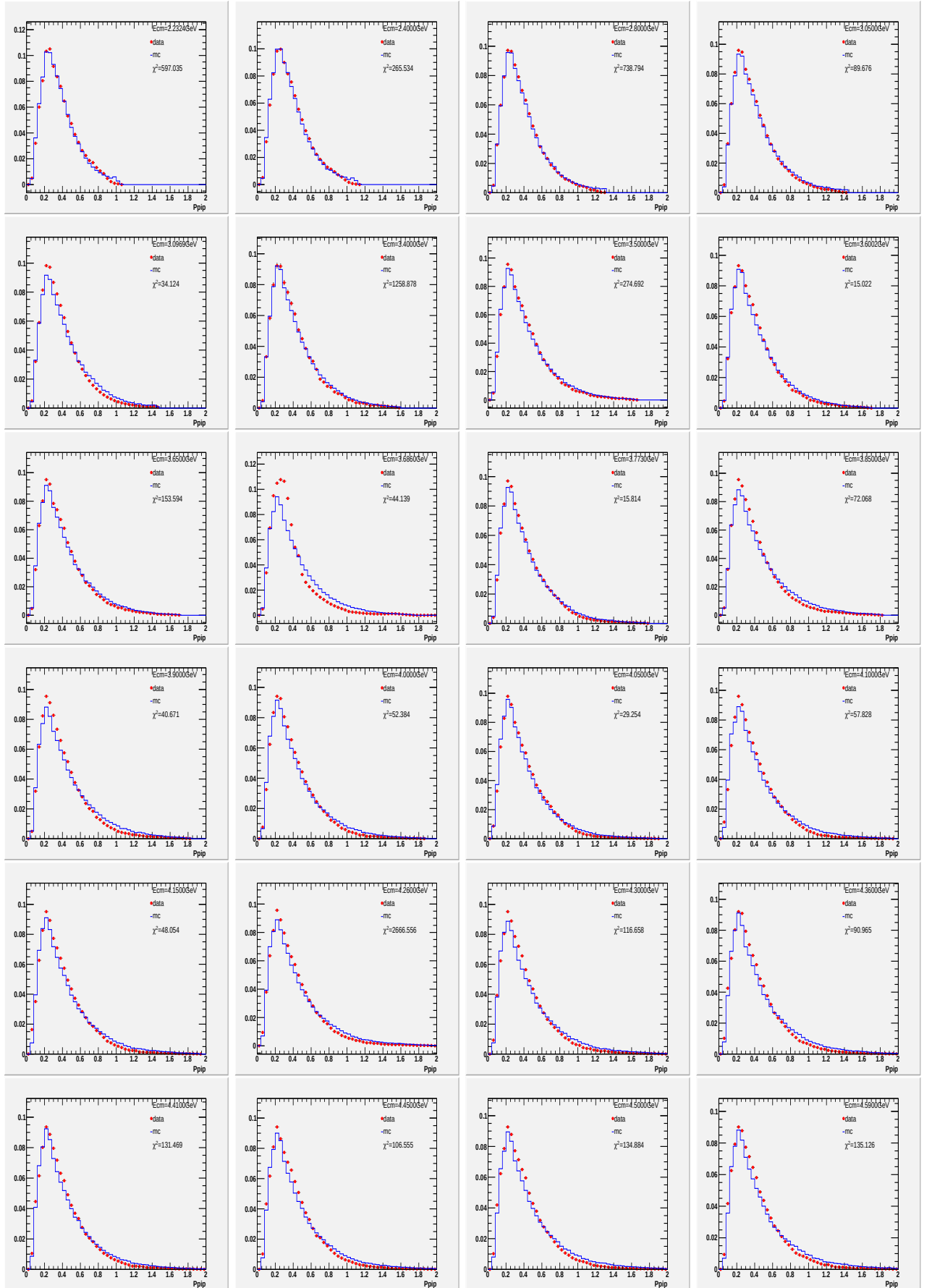


Fig. P of π^+ .

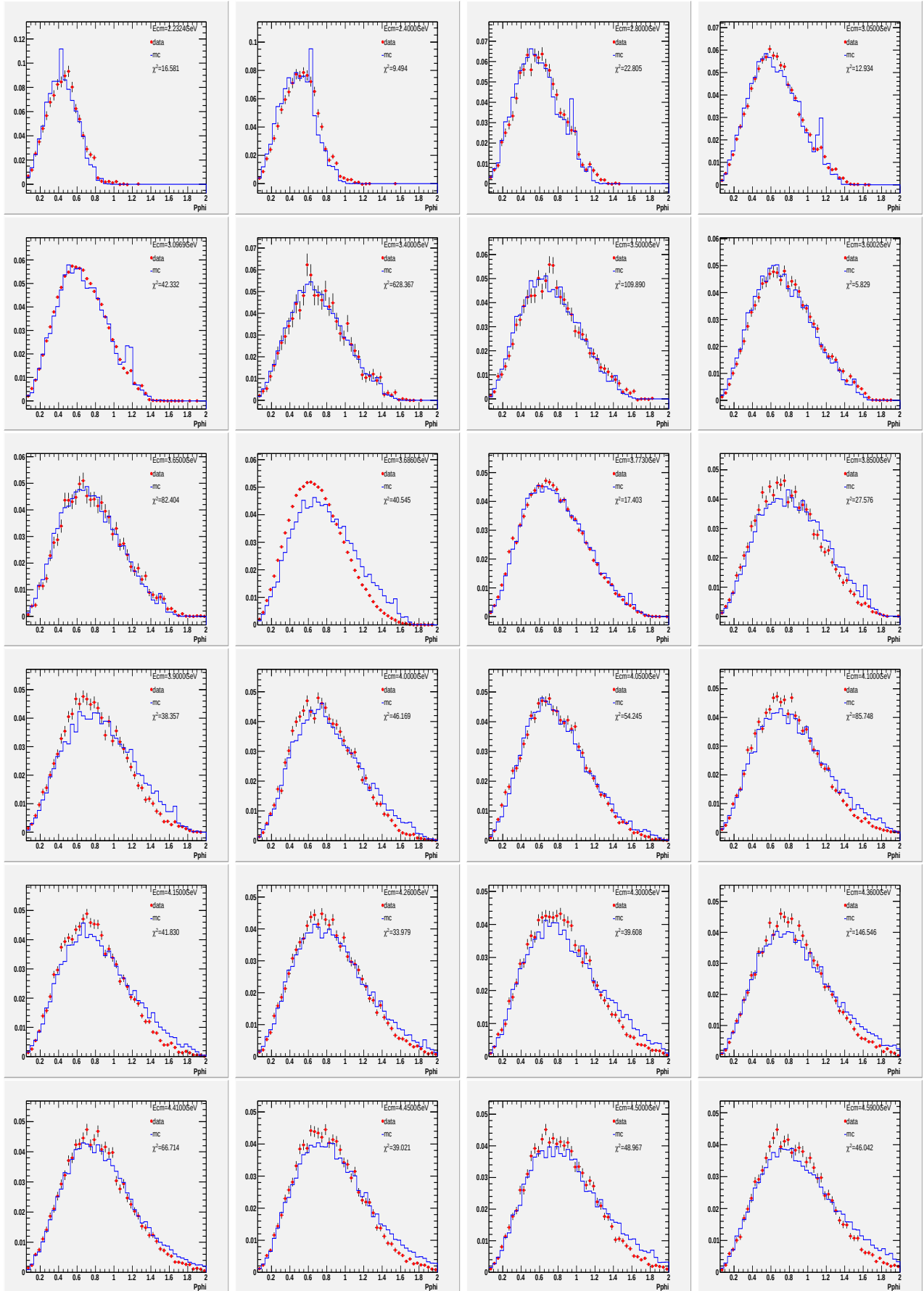


Fig. P of ϕ .

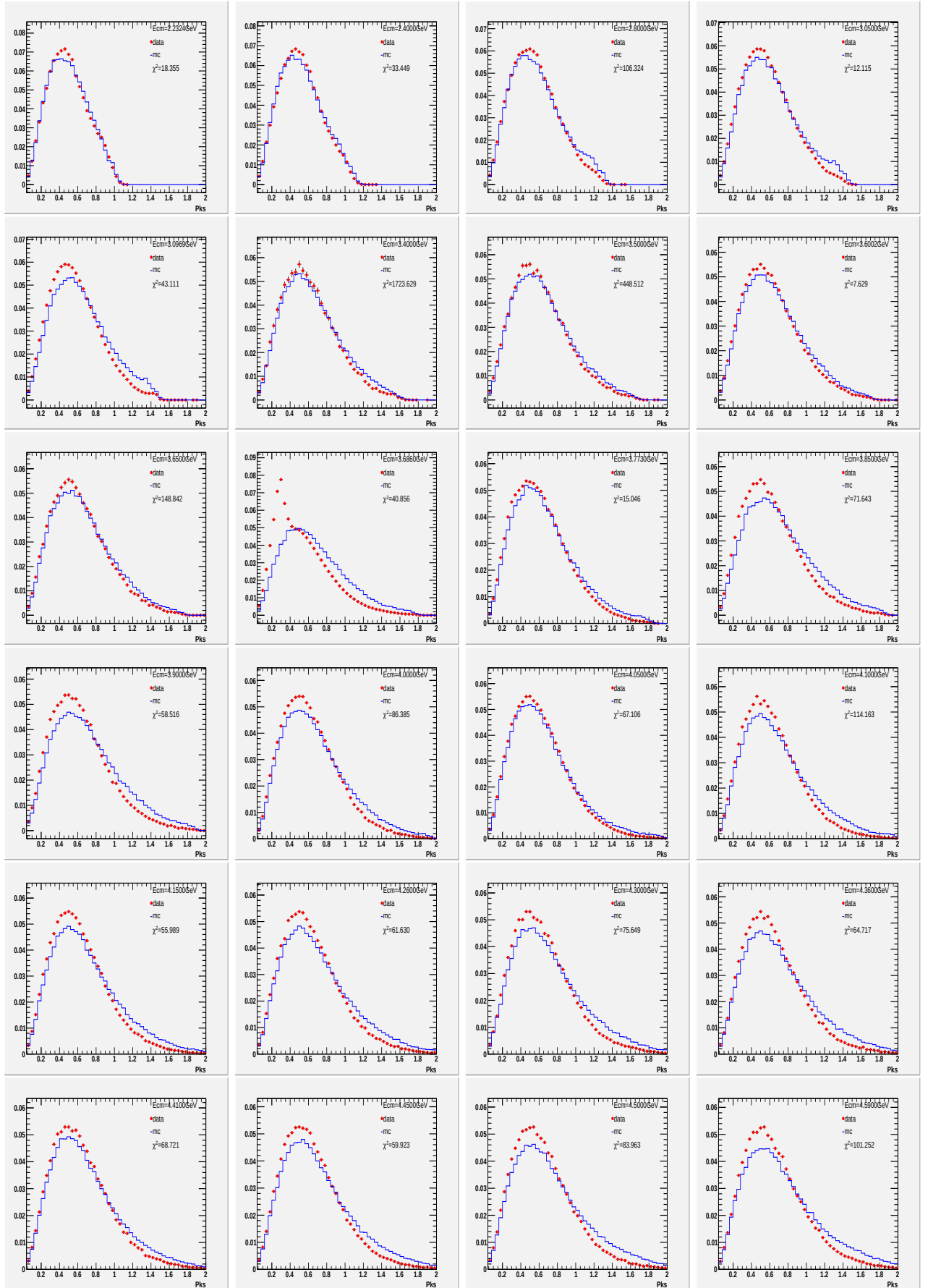


Fig. P of K_s .

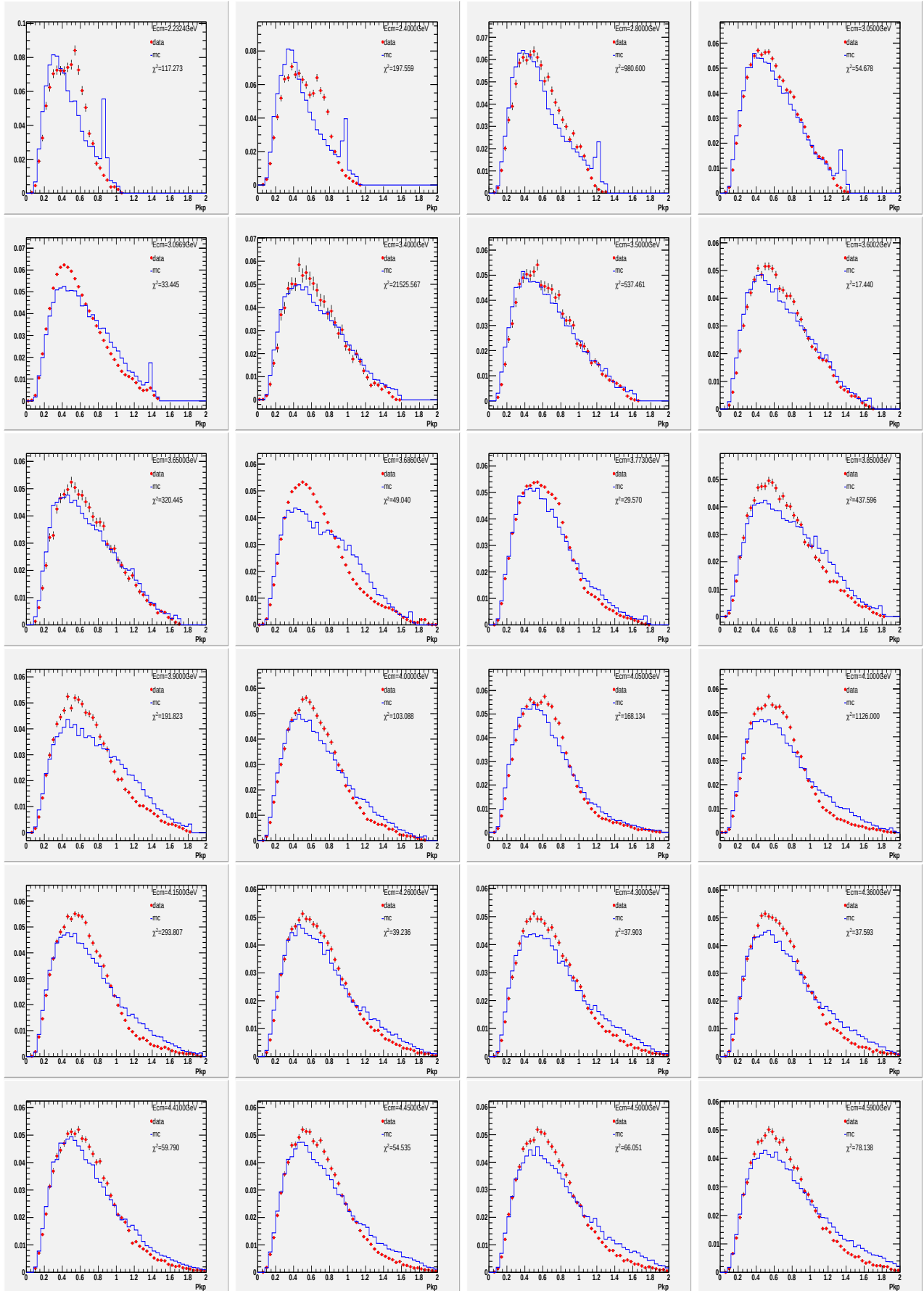


Fig. P of K^+ .

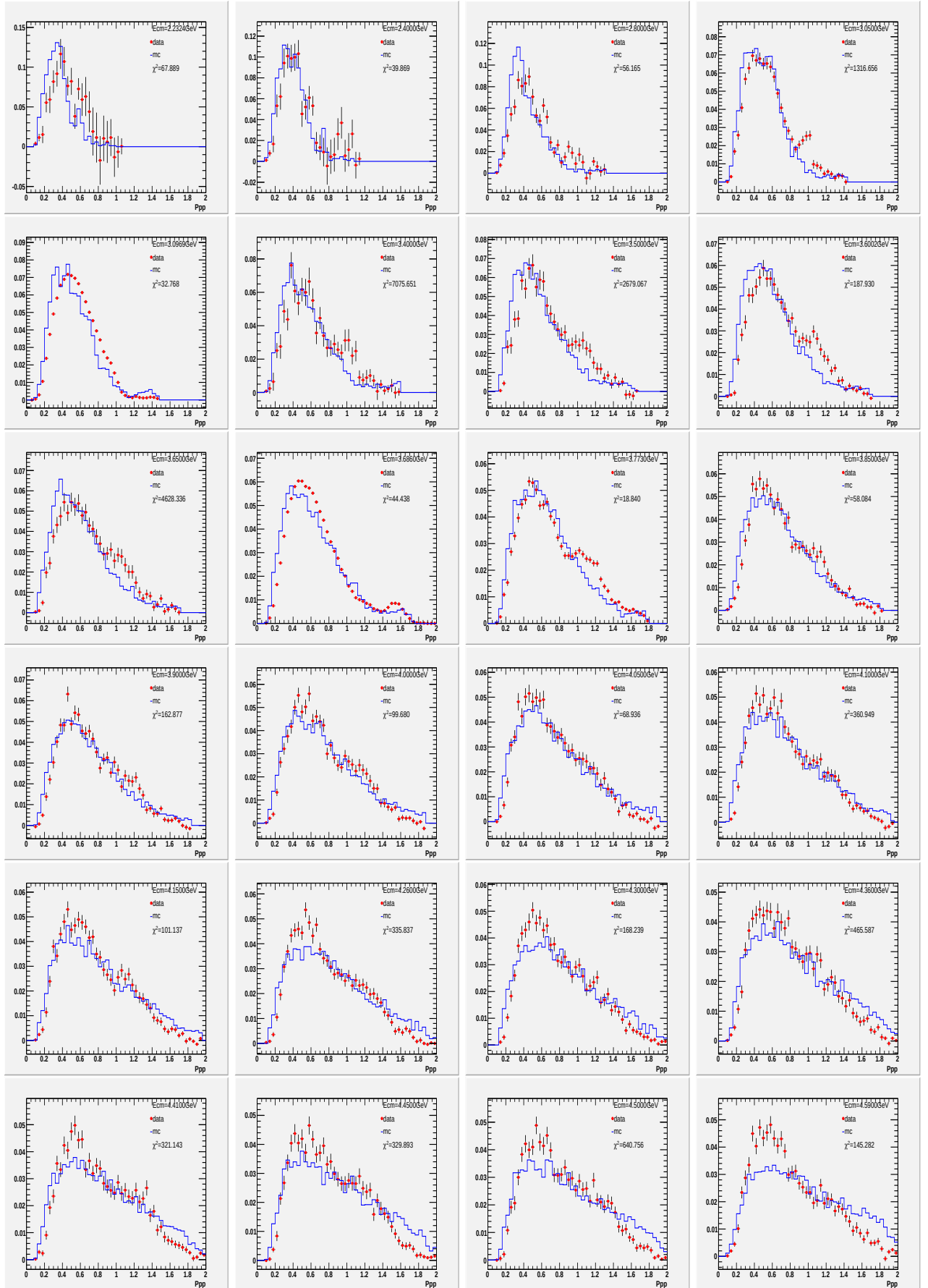


Fig. P of p^+ .

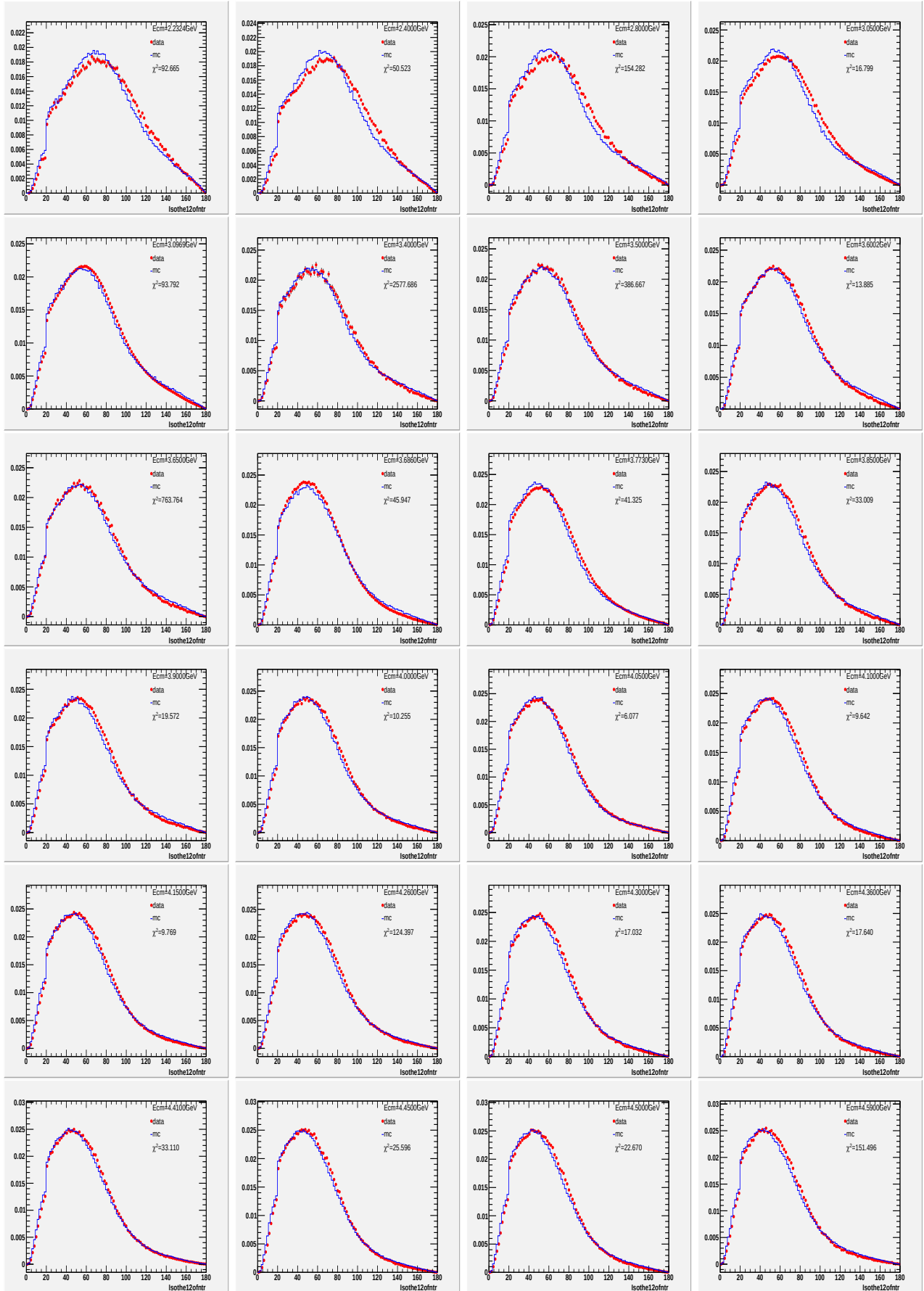


Fig.2 The angle between photon and charged track in 2-prong event.

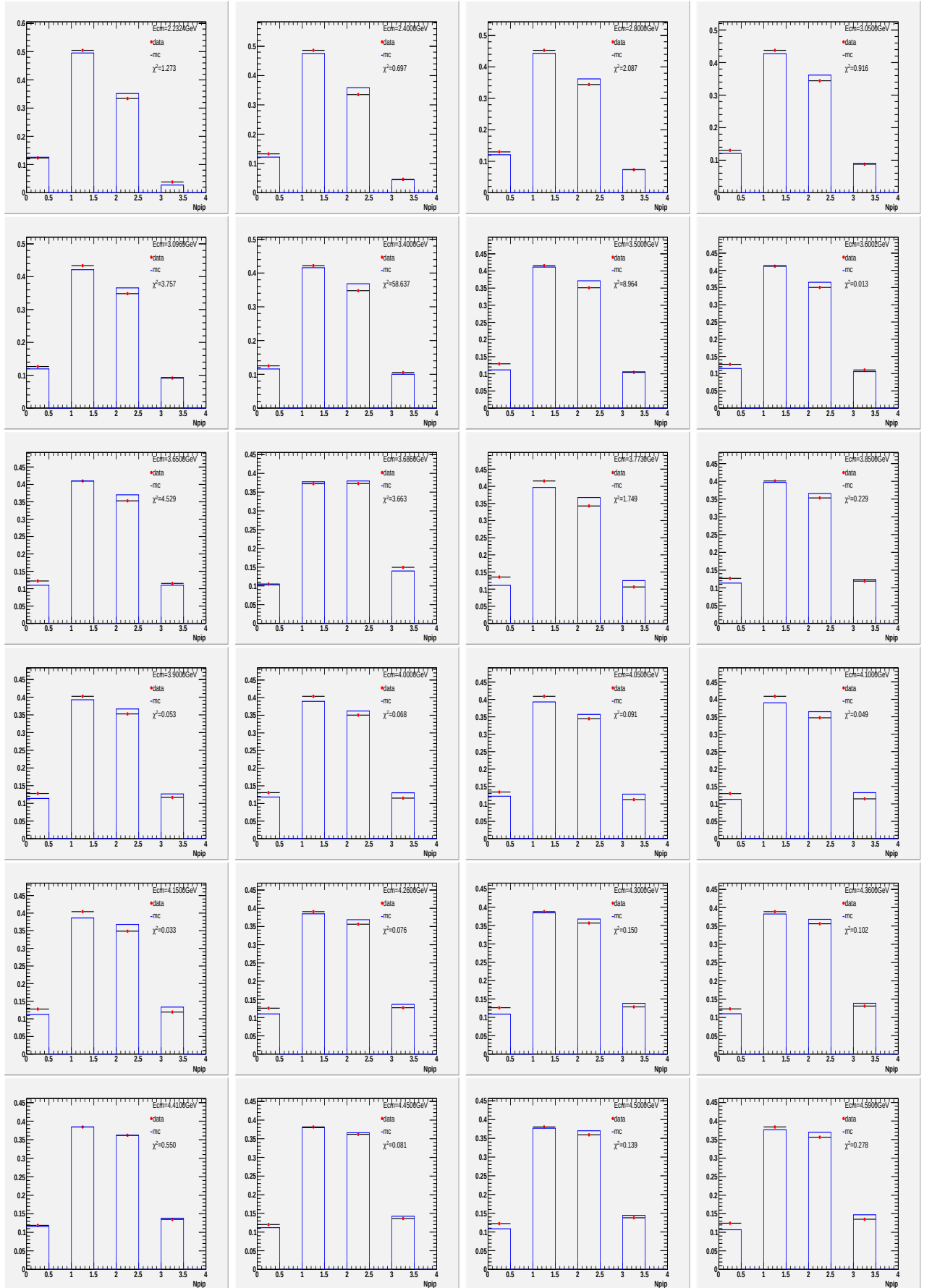


Fig.2 N_{π^+} .

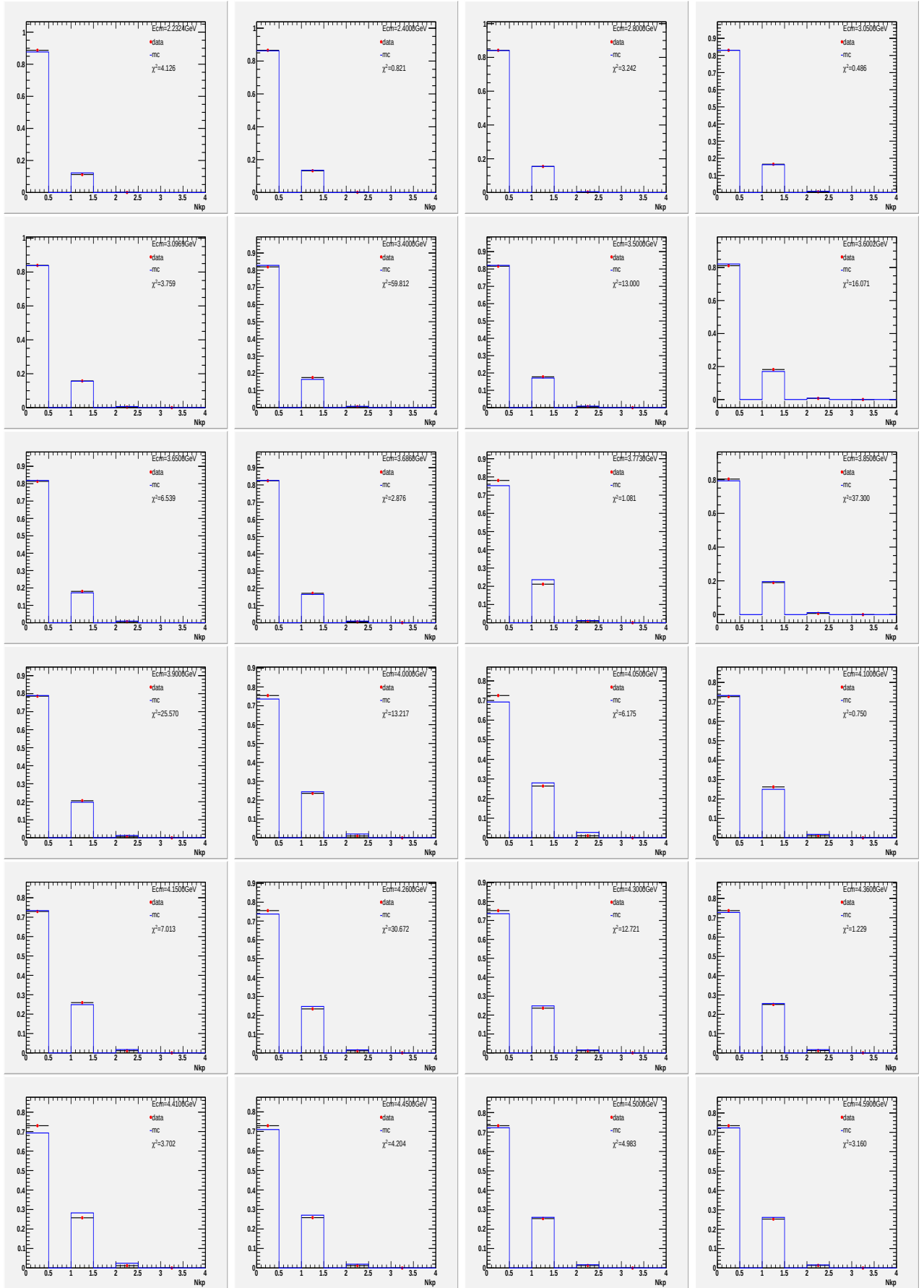


Fig.2 N_{K^+} .

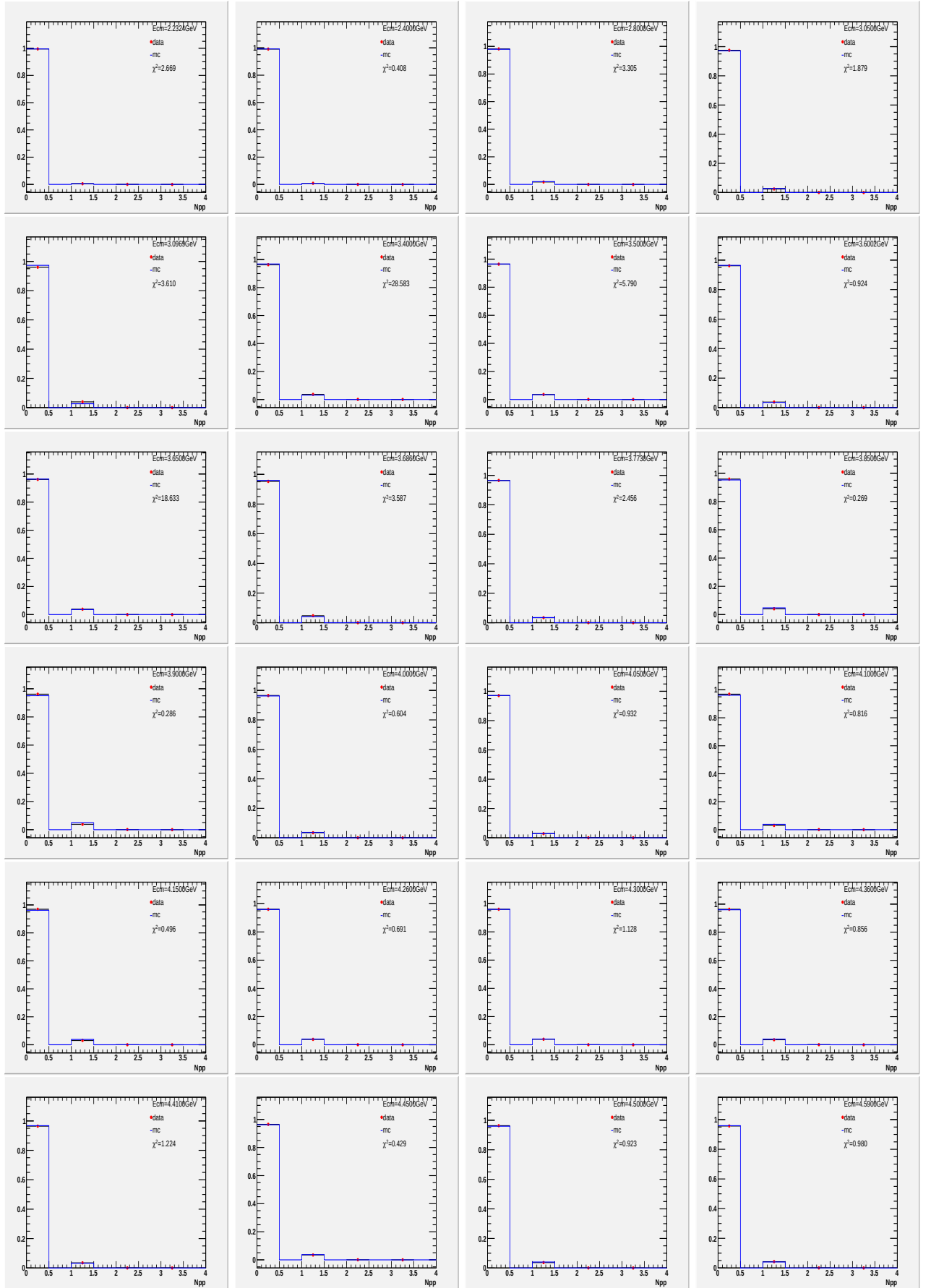


Fig.2 N_{p+} .

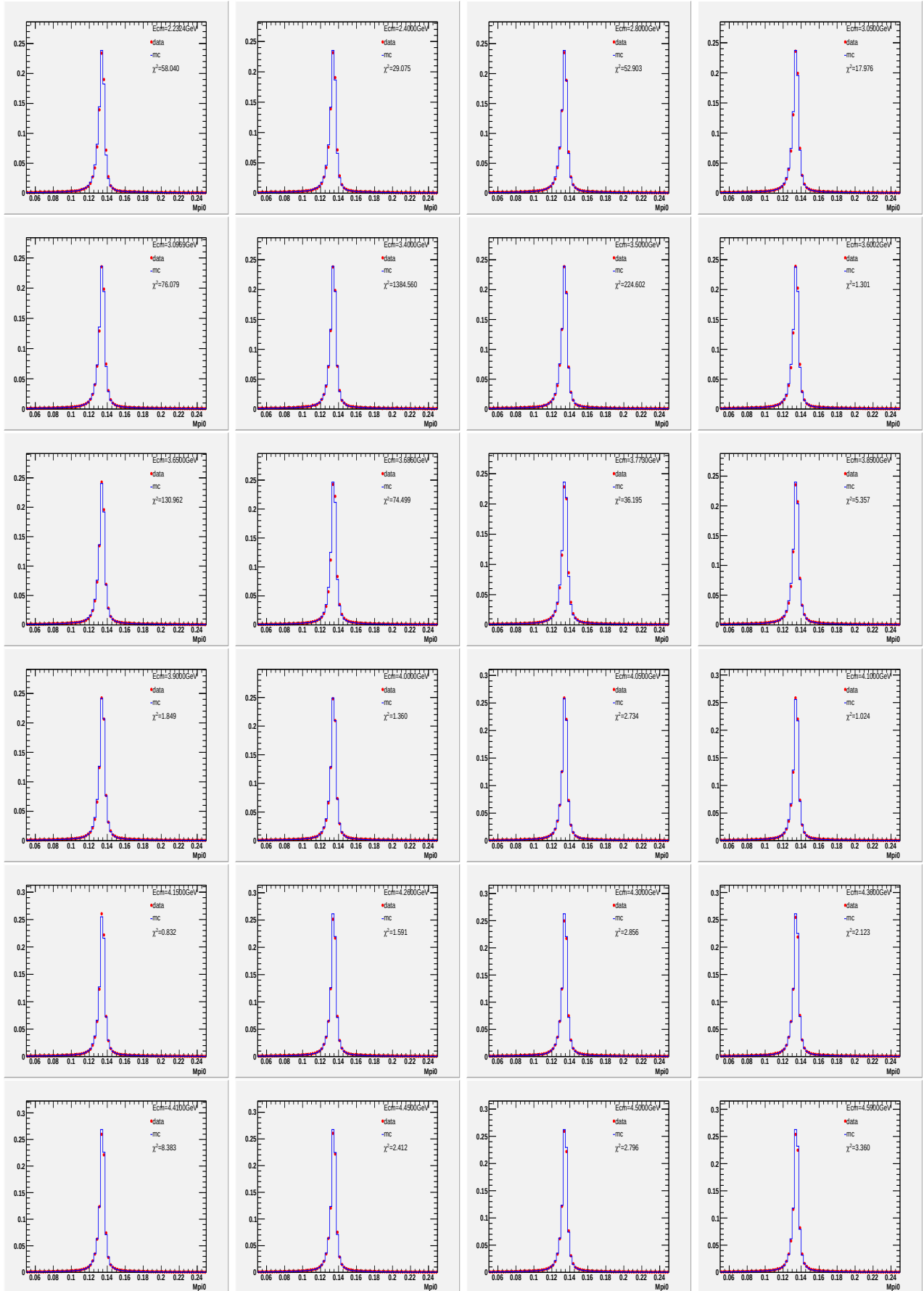


Fig.2 M_{π^0} .

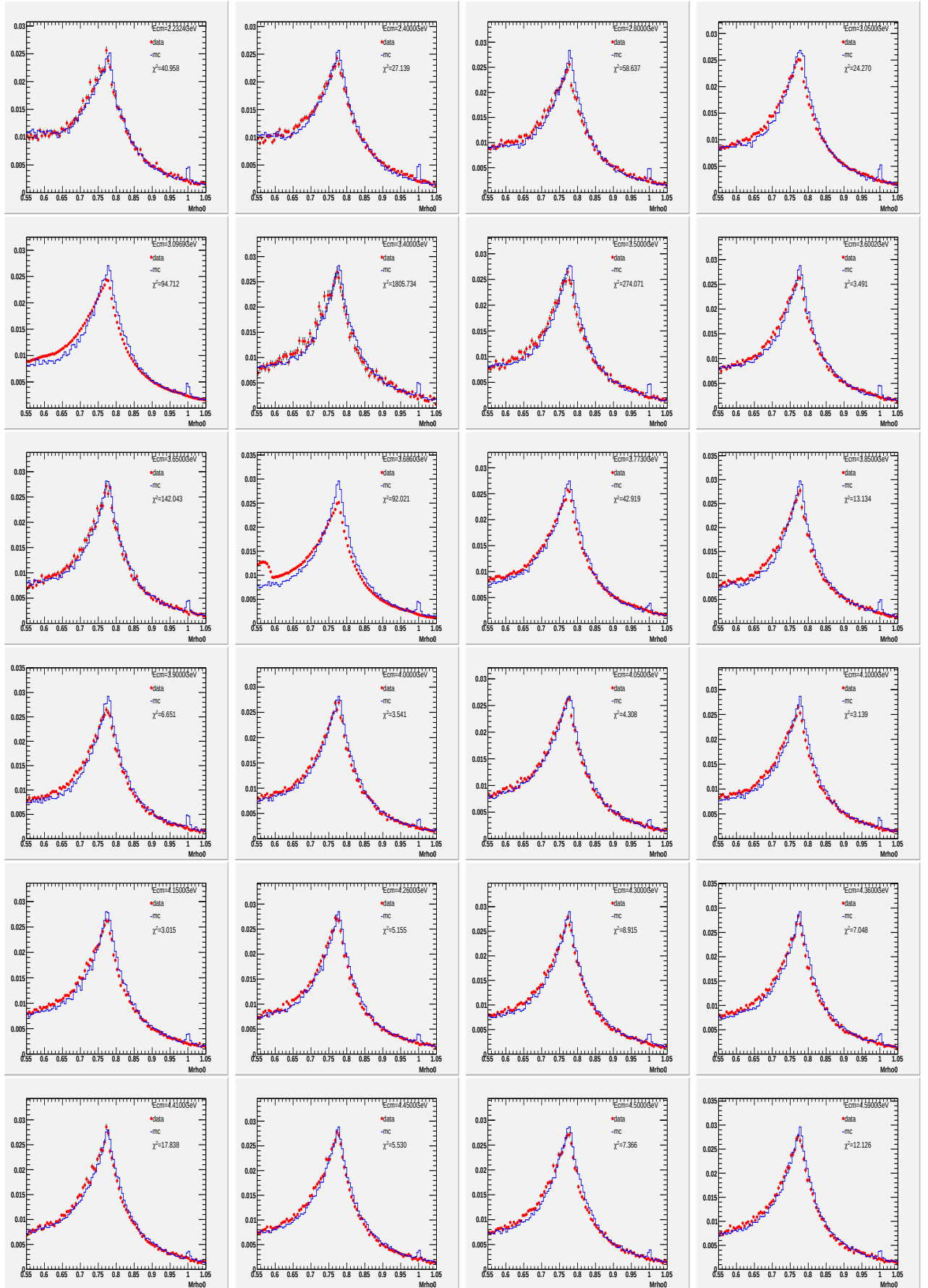


Fig. $M_{\pi^0\pi^0}$.

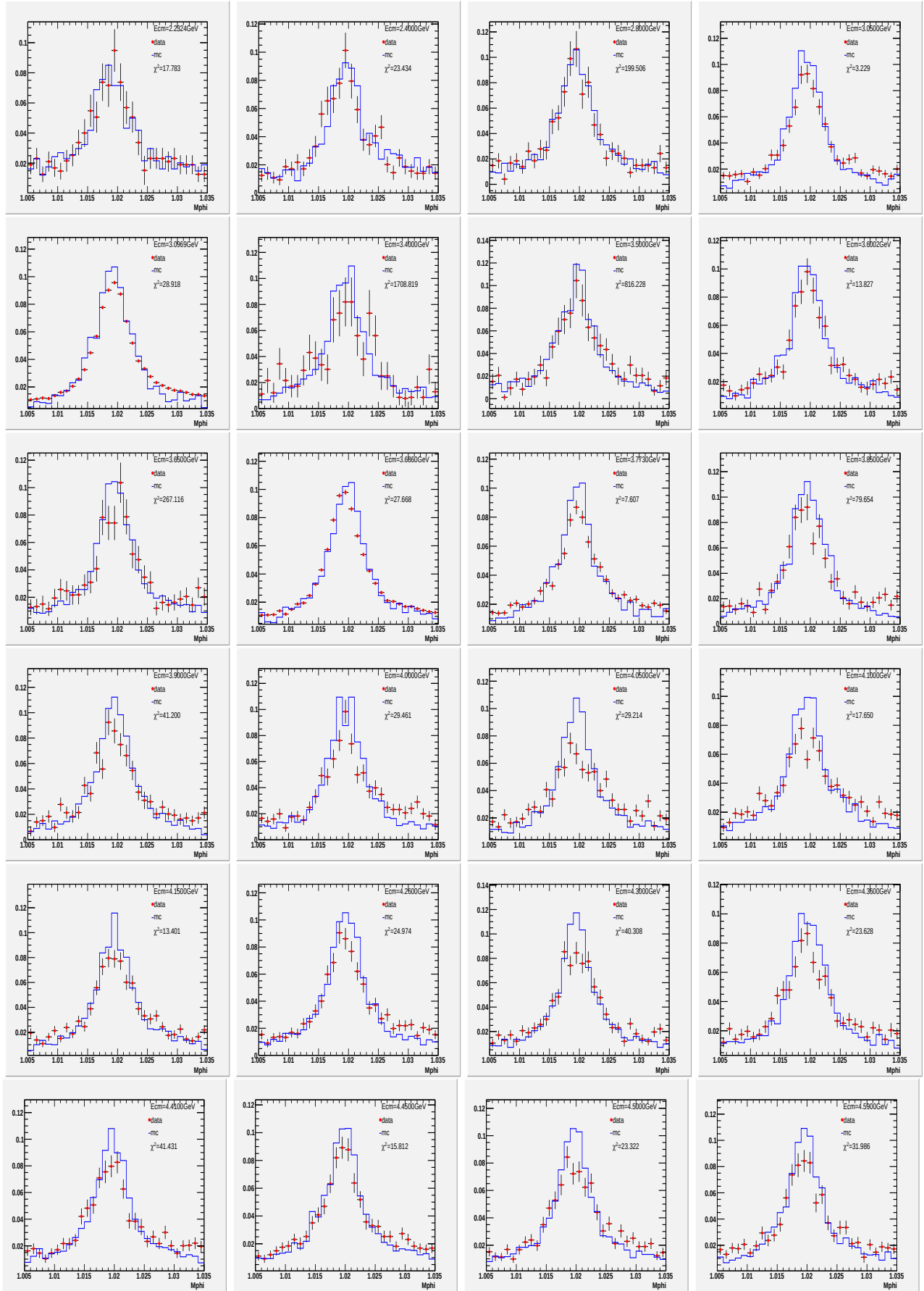


Fig. M_ϕ .

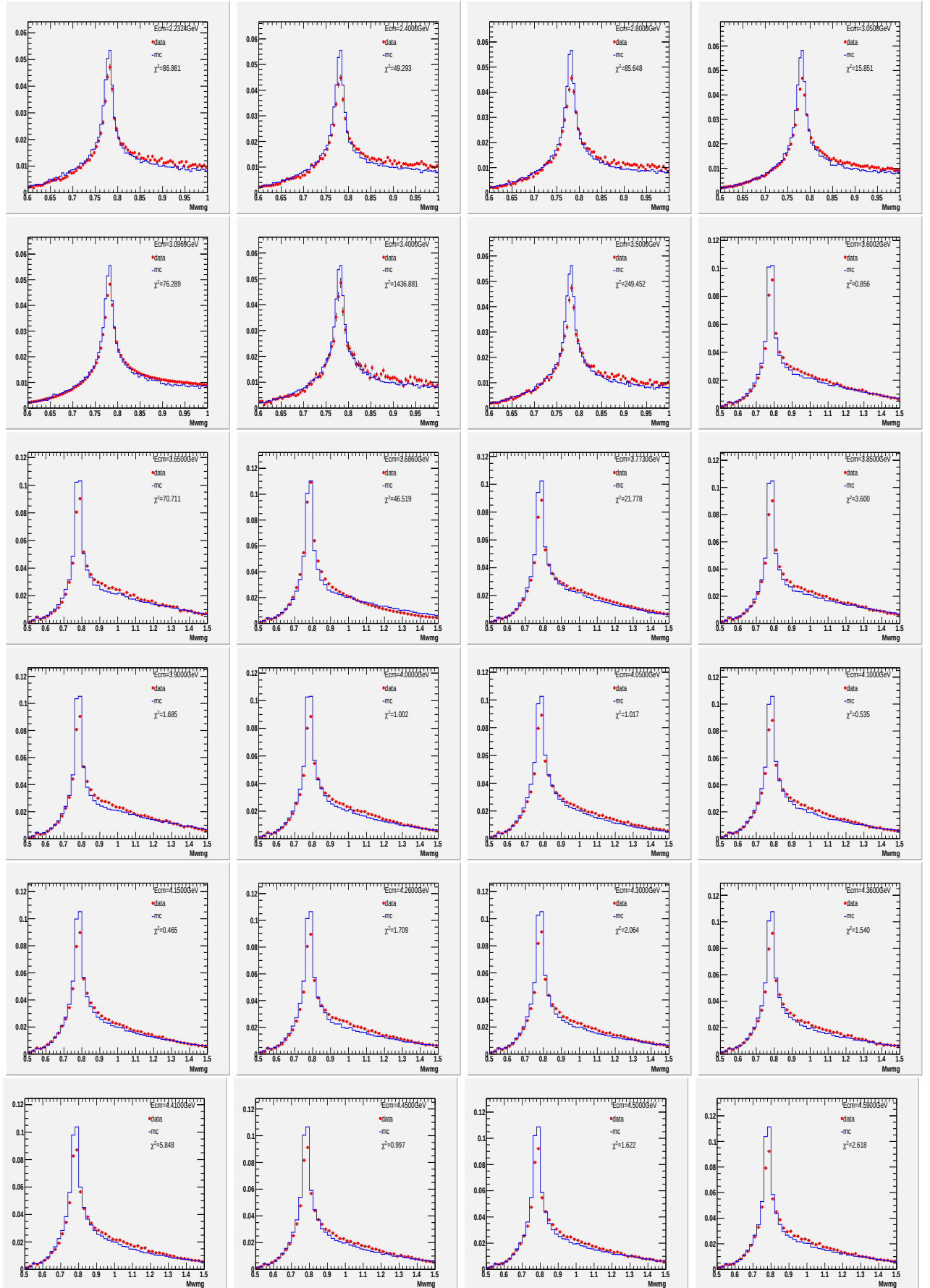


Fig. M_{ω} .

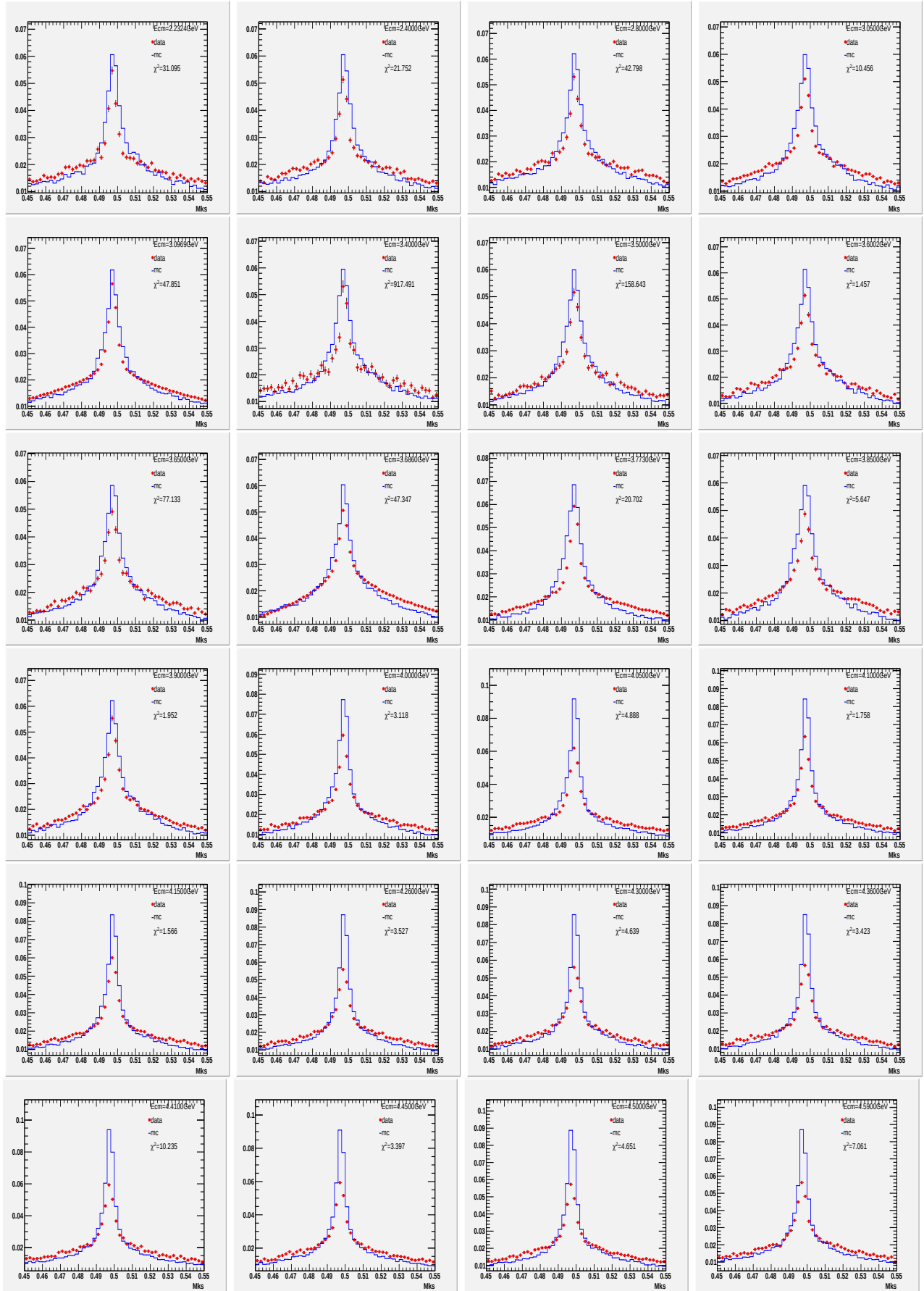


Fig. M_{K_s} .

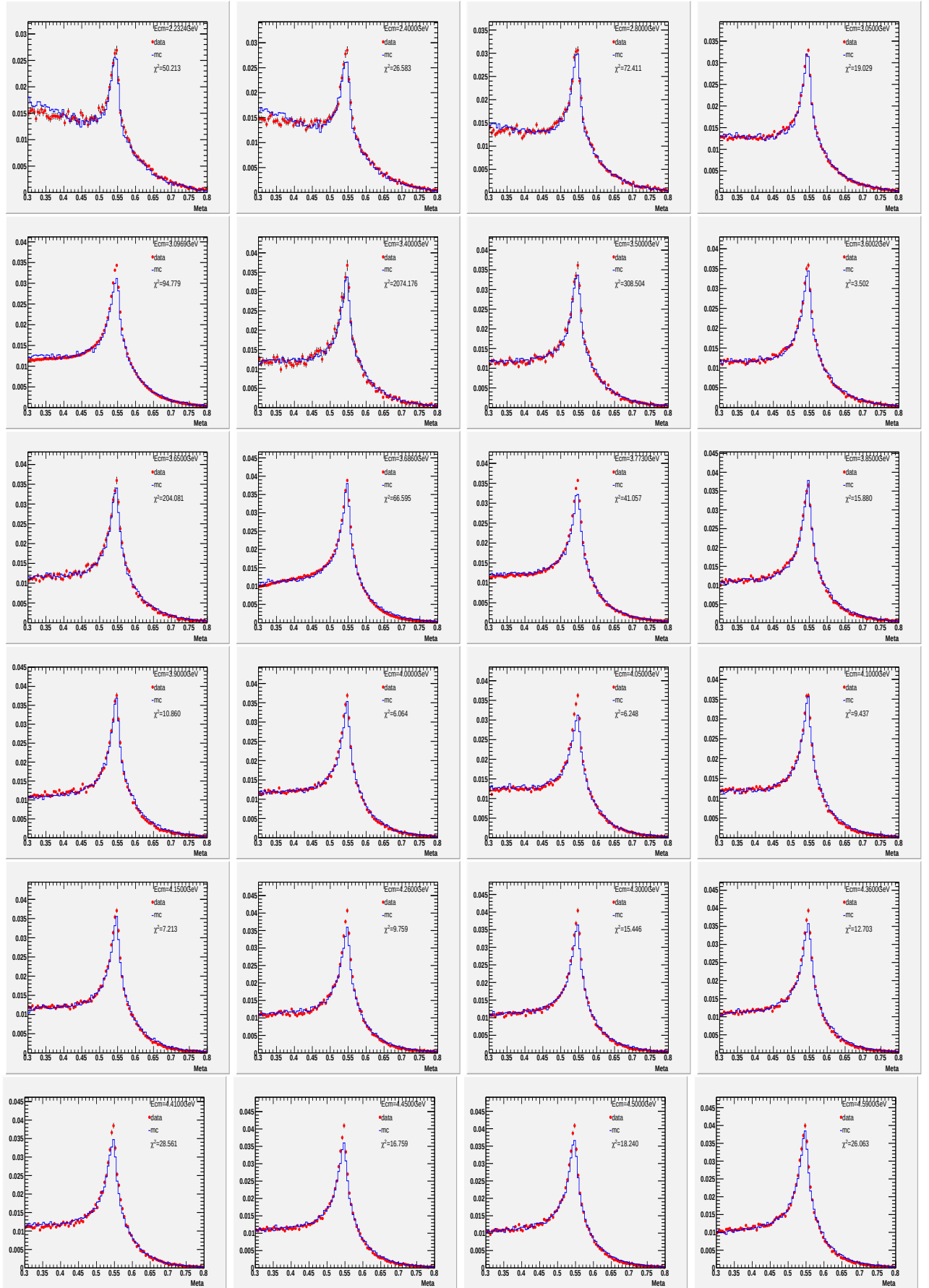


Fig. M_{η} .

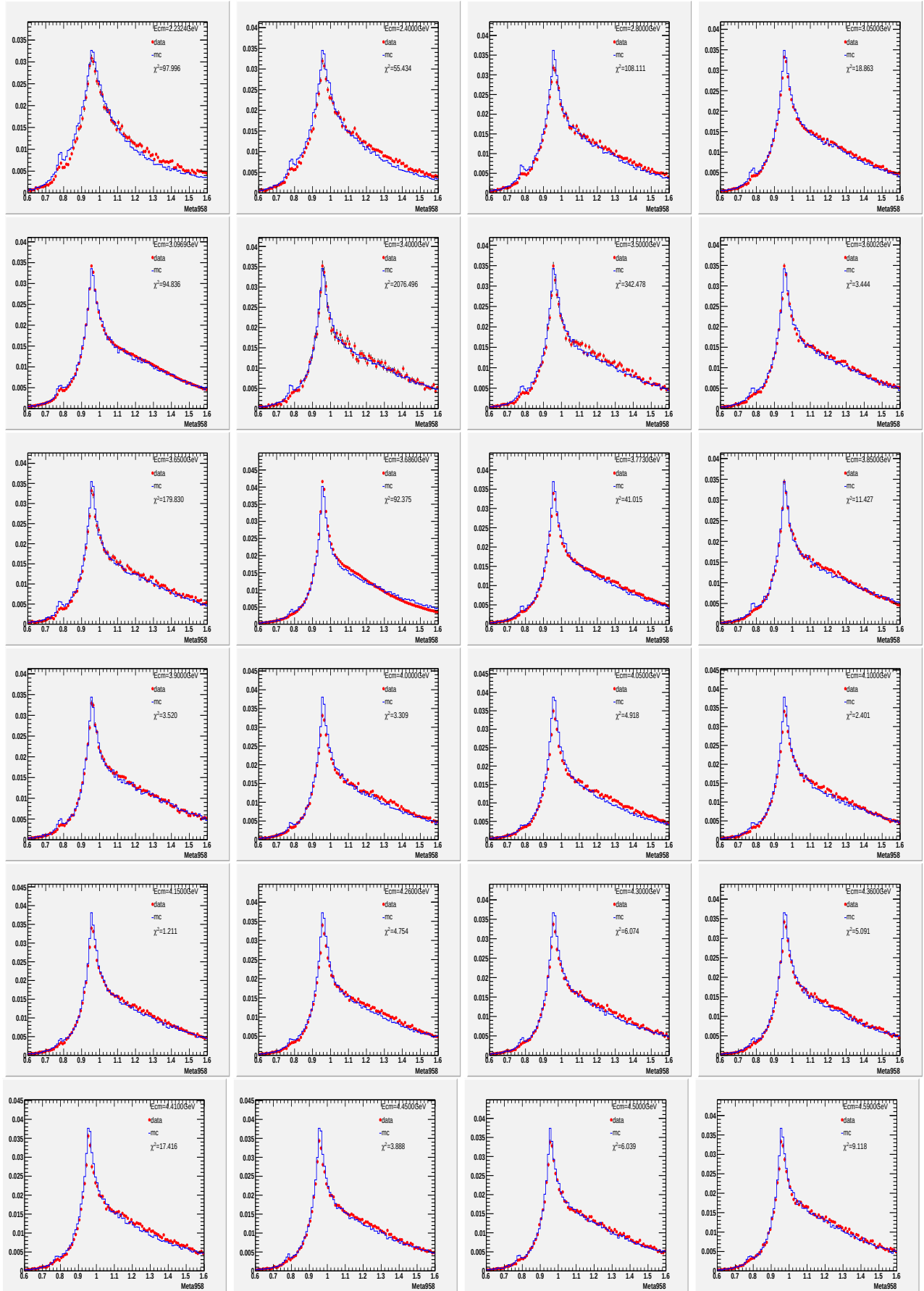


Fig. $M_{\eta_{958}}$.

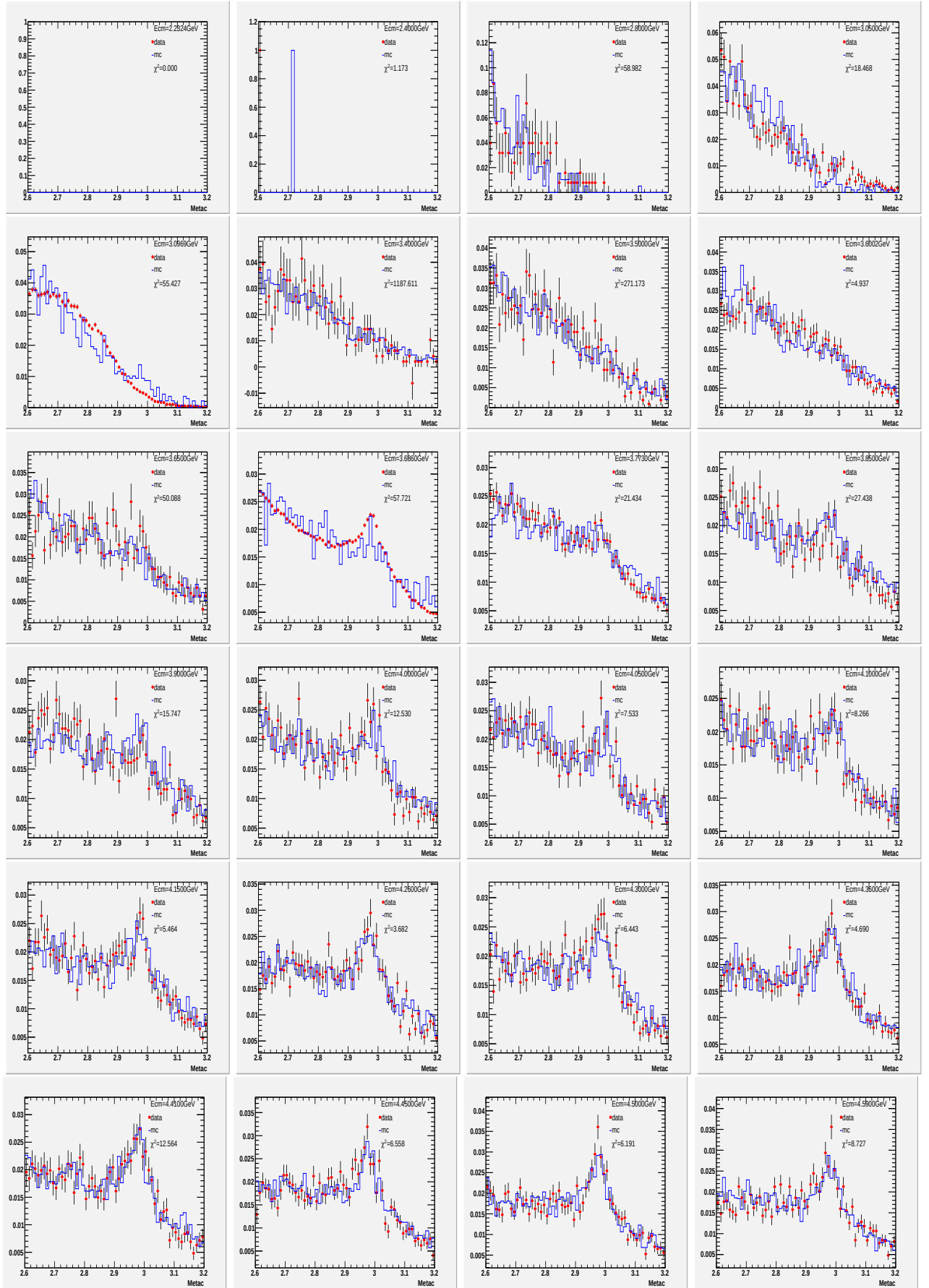


Fig. M_{η_c} .

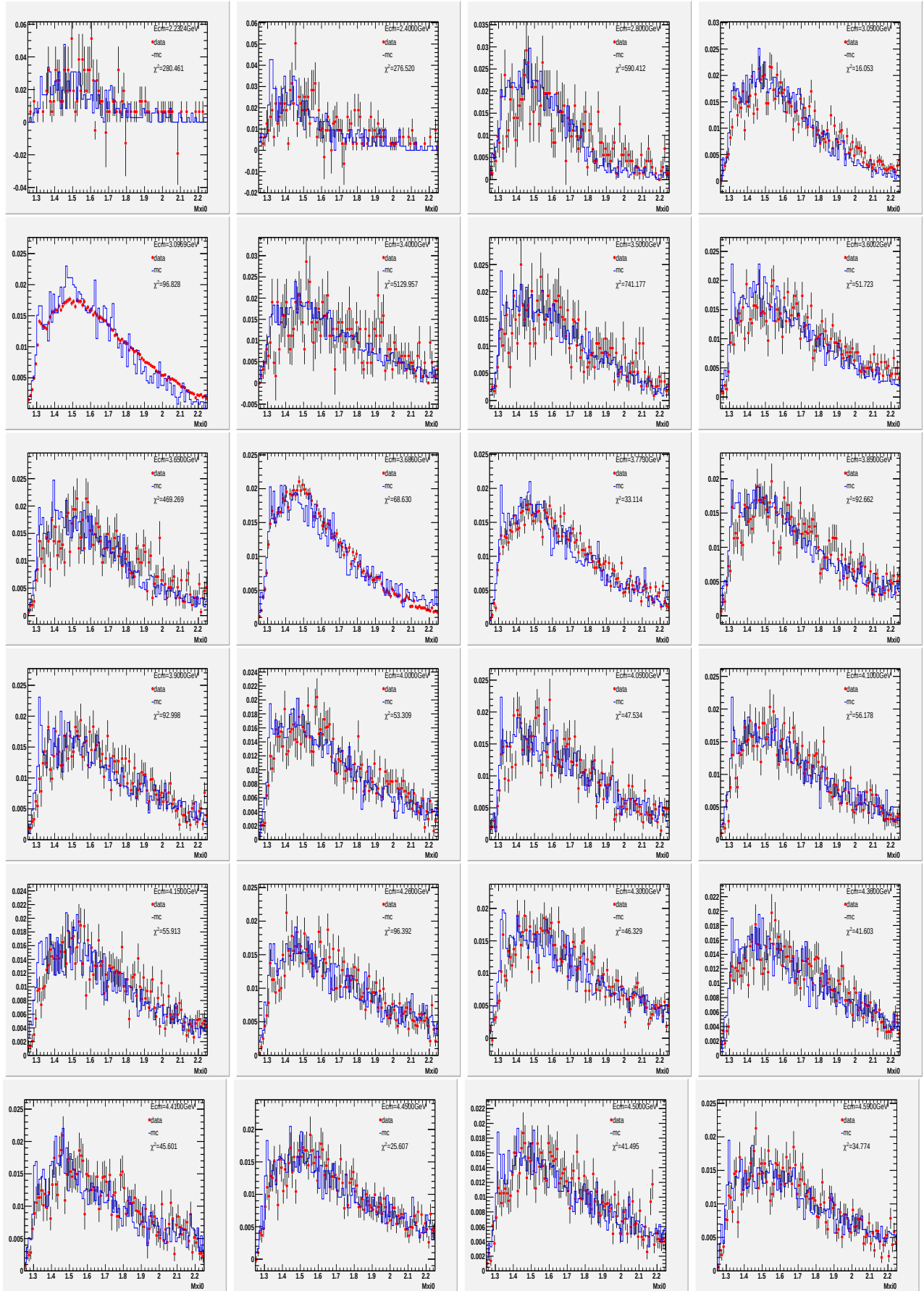


Fig. M_{χ_0} .

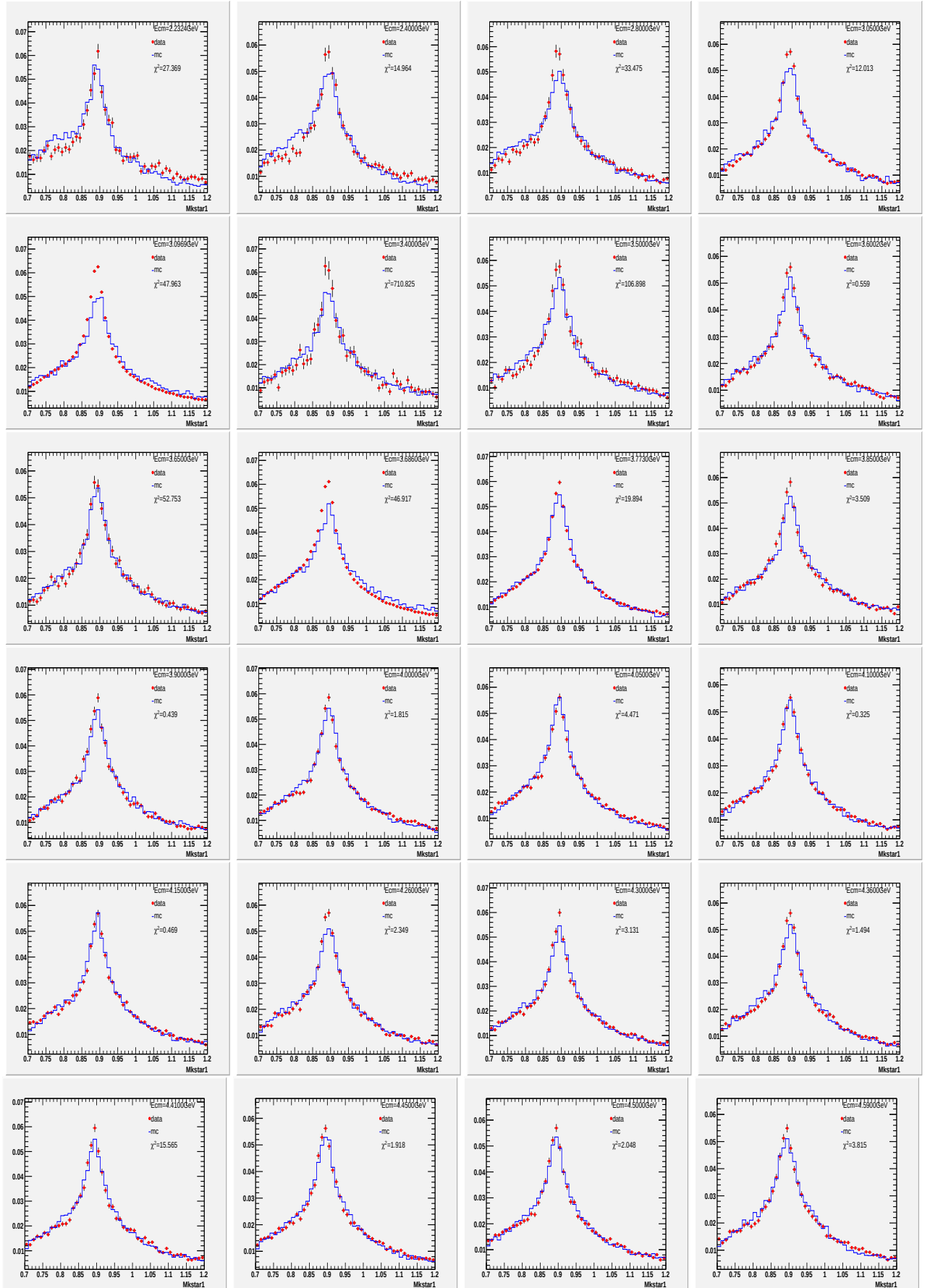


Fig. $N_{K^*_1}$.

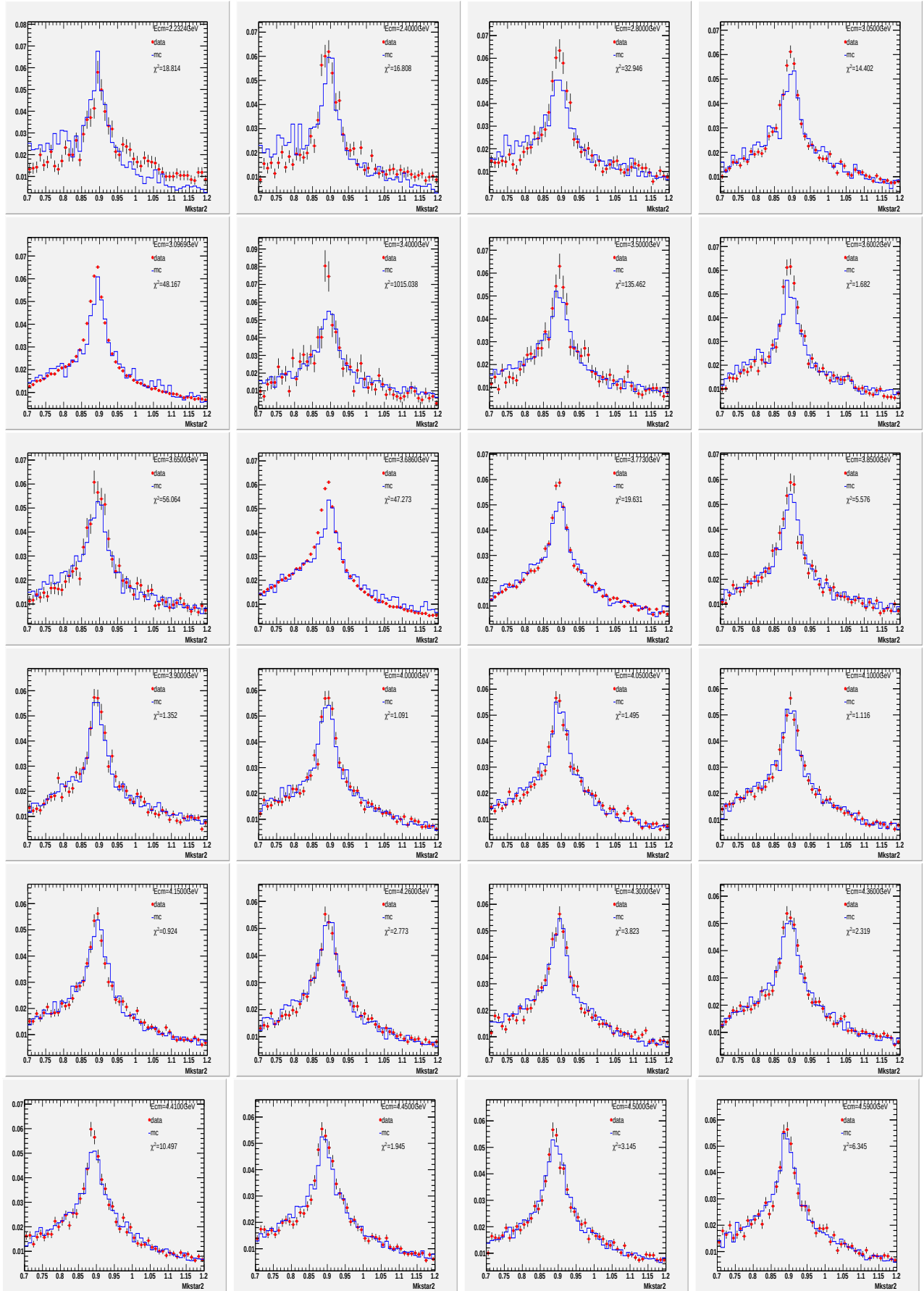


Fig. N_{K^*2} .

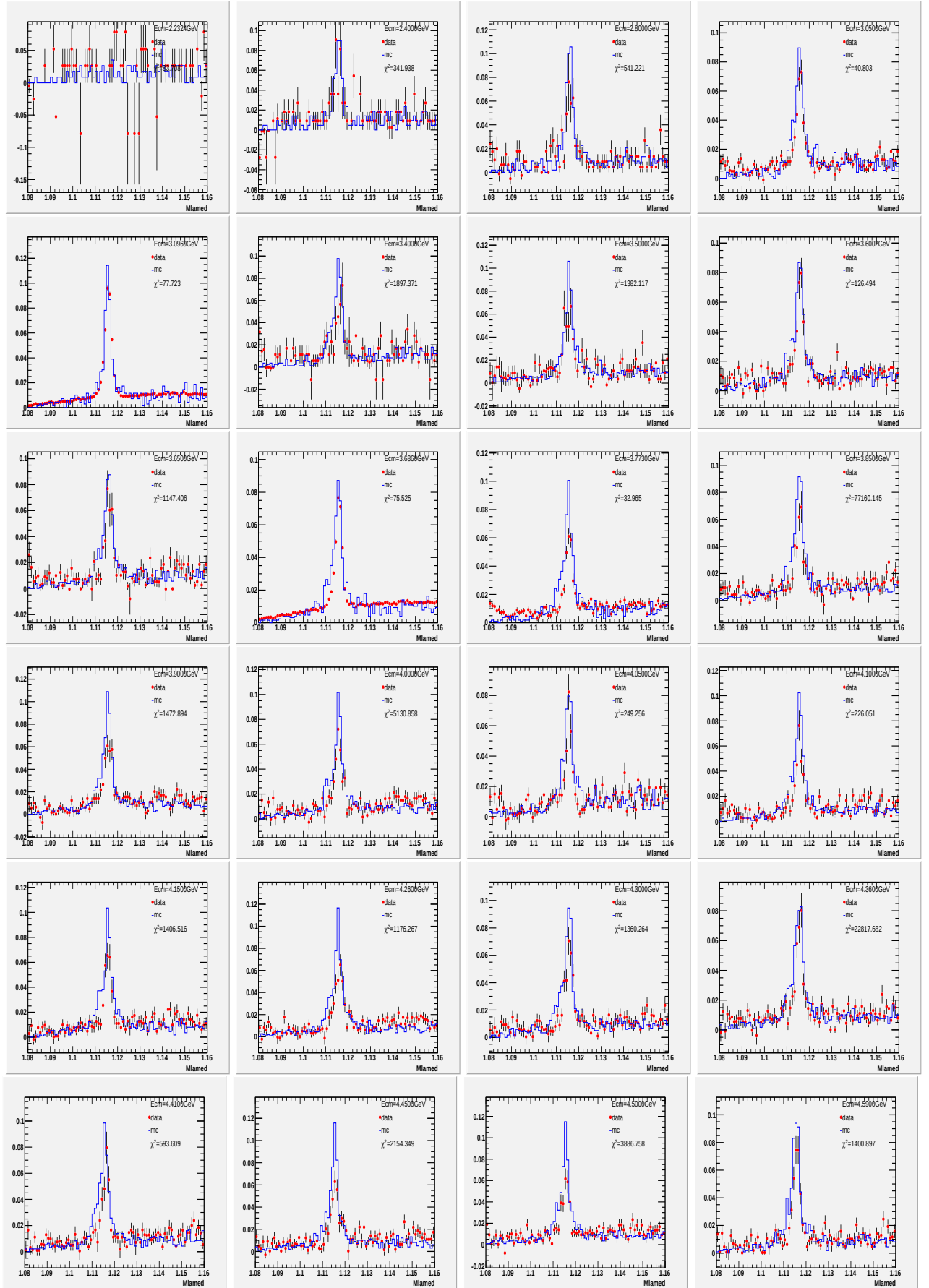


Fig. M_{λ} .

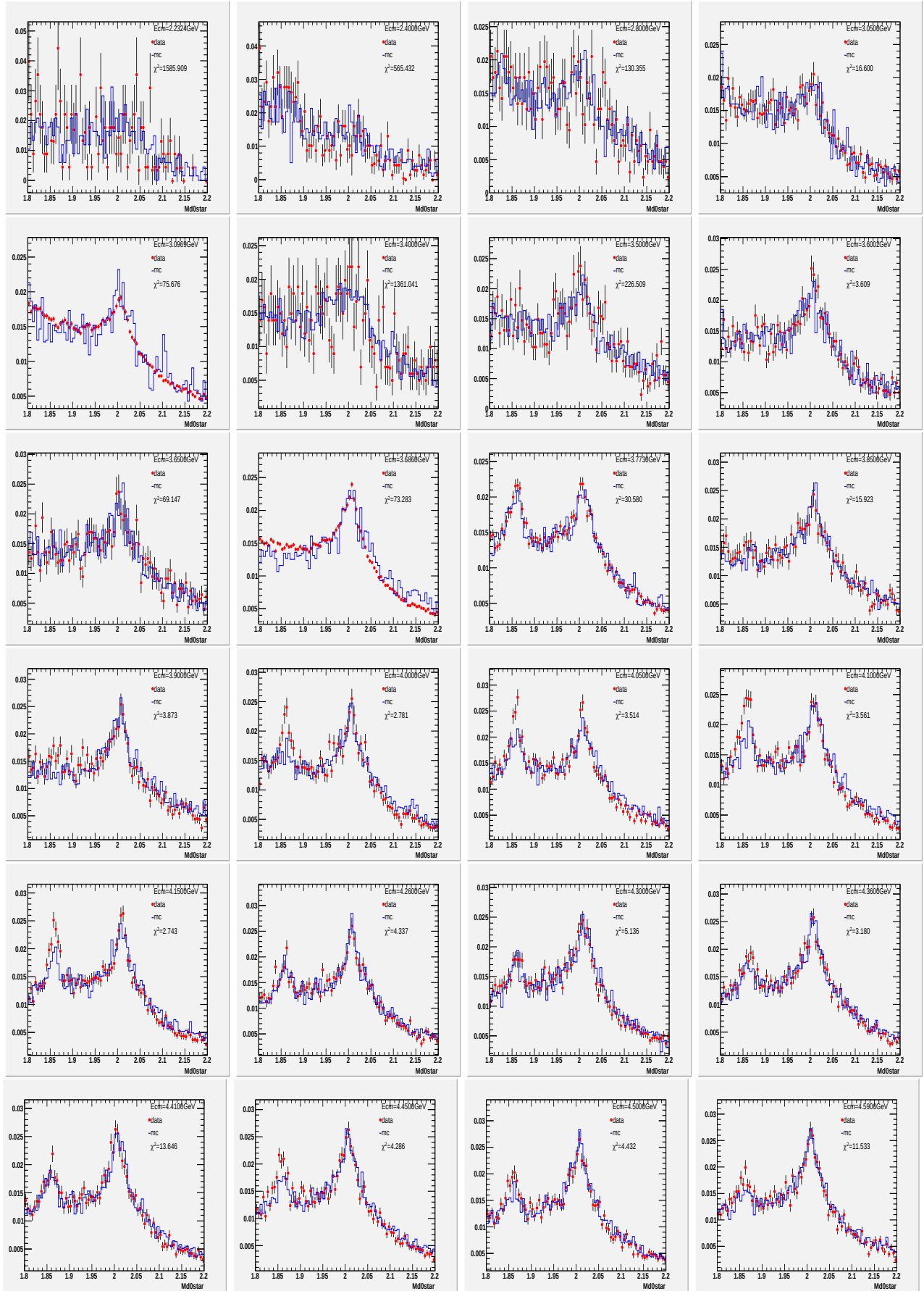


Fig. $M_{D_0^*}$.

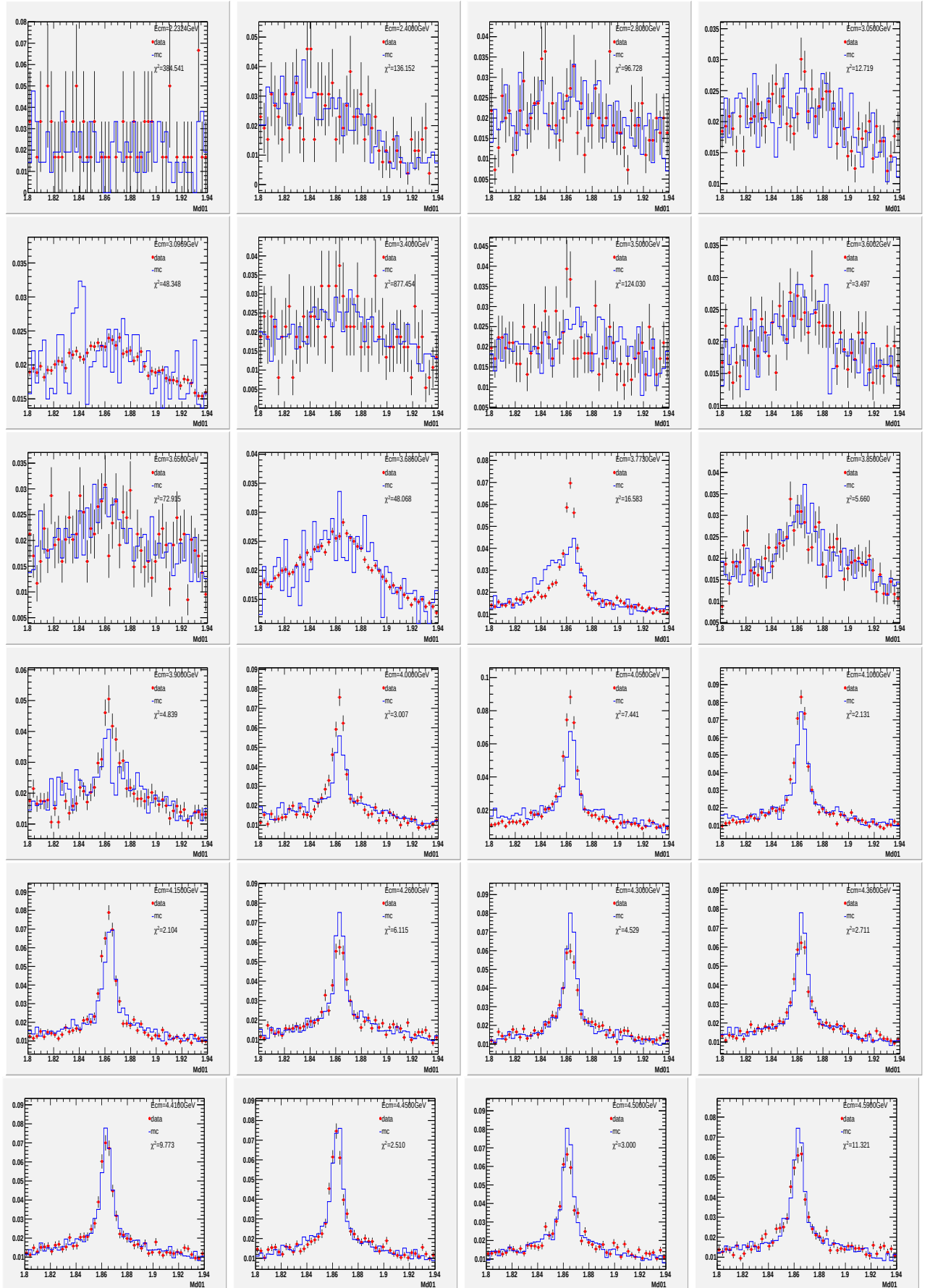


Fig. $M_{D_0^1}$.

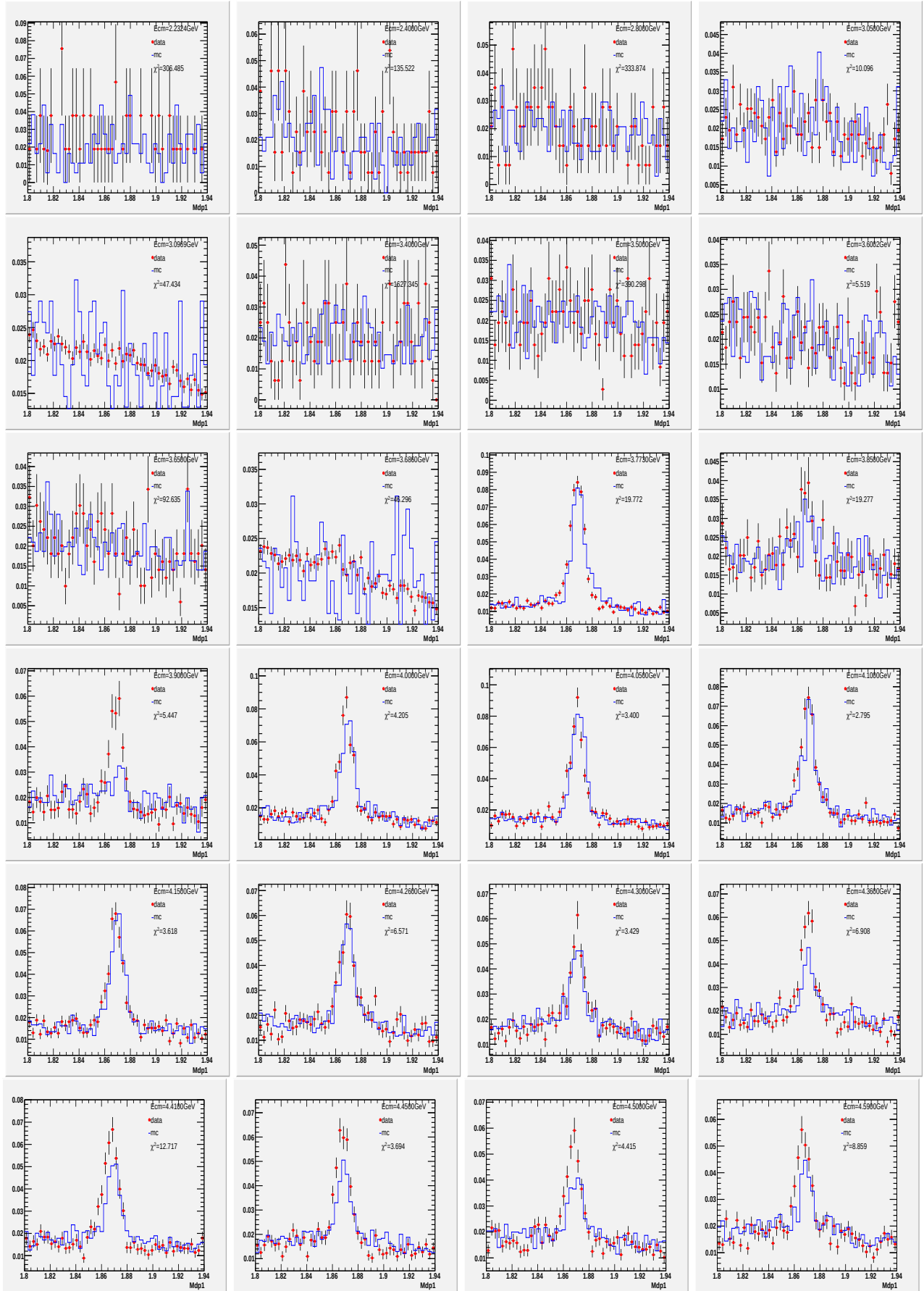


Fig. $M_{D_1^+}$.

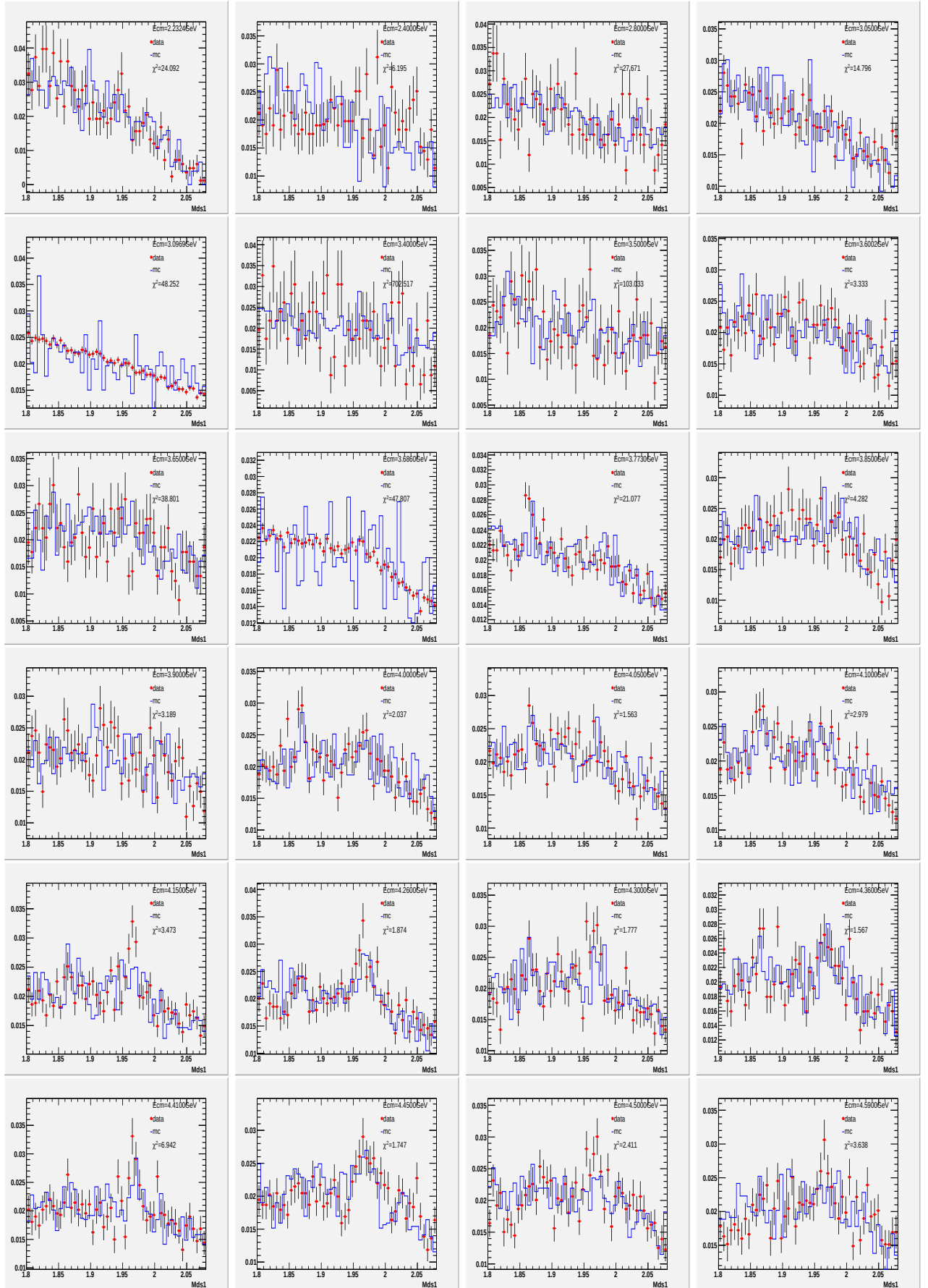


Fig. $M_{D_s^1}$.

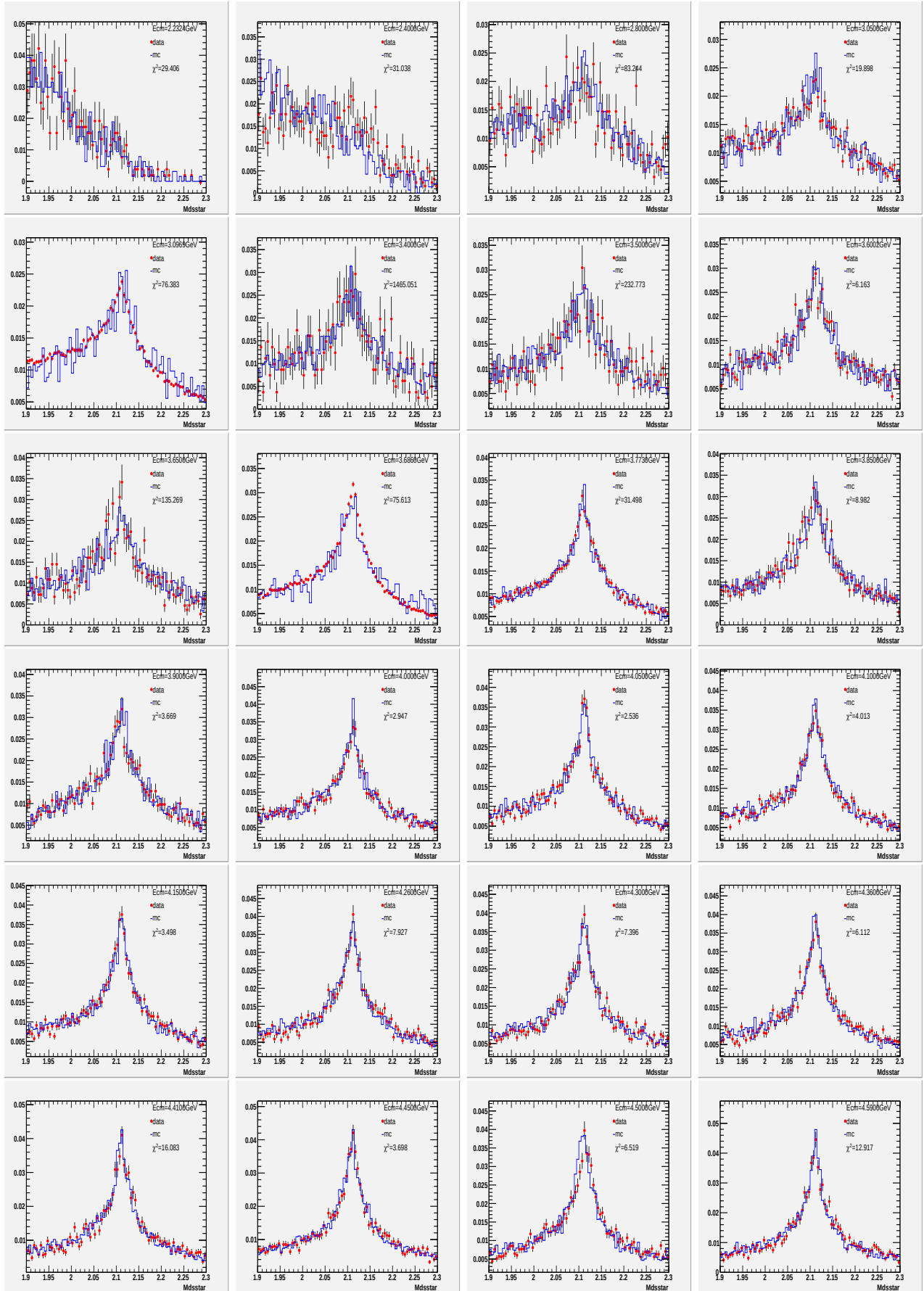


Fig. $M_{D_s^+}$.

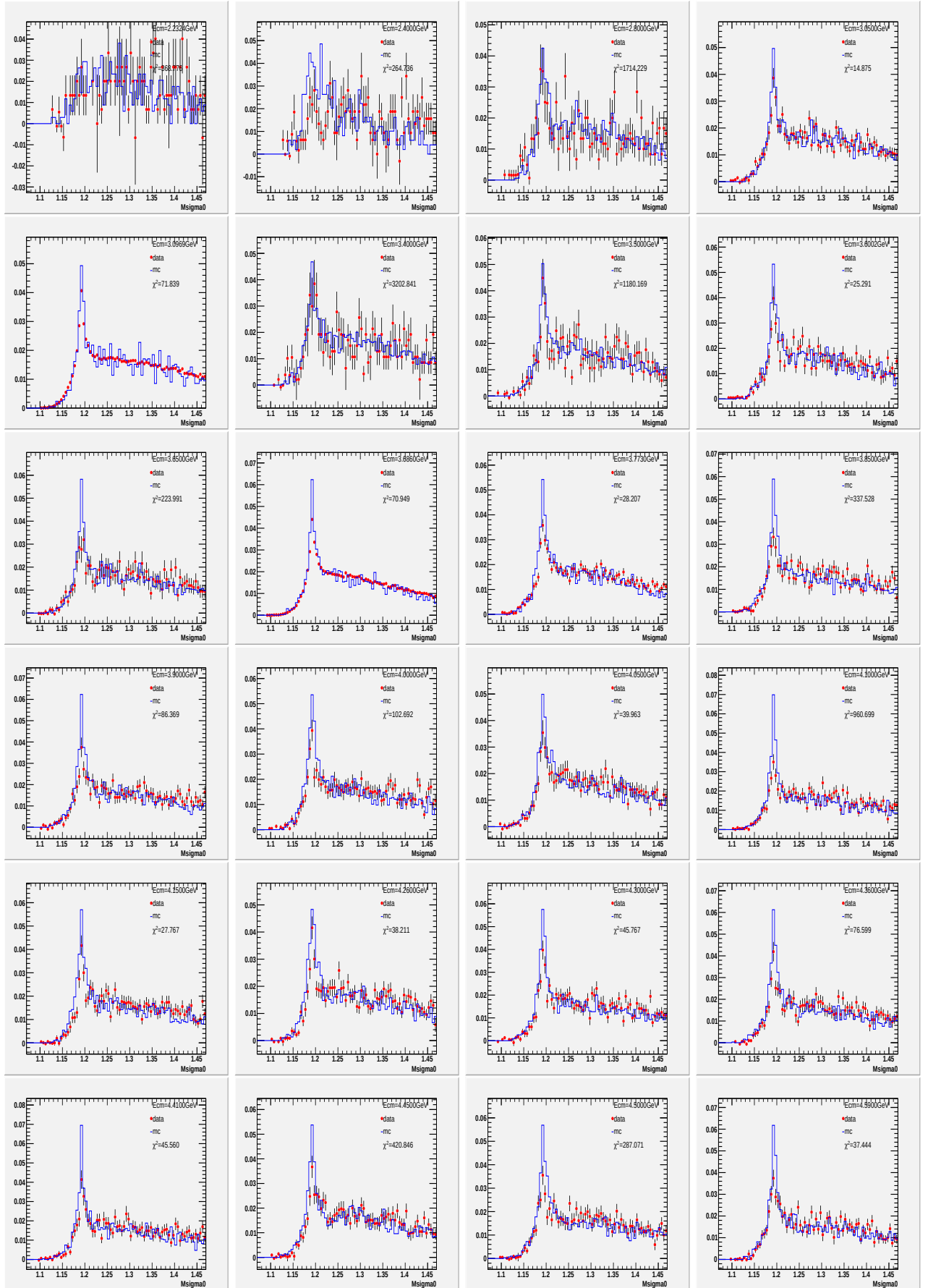


Fig. M_{Σ^0} .

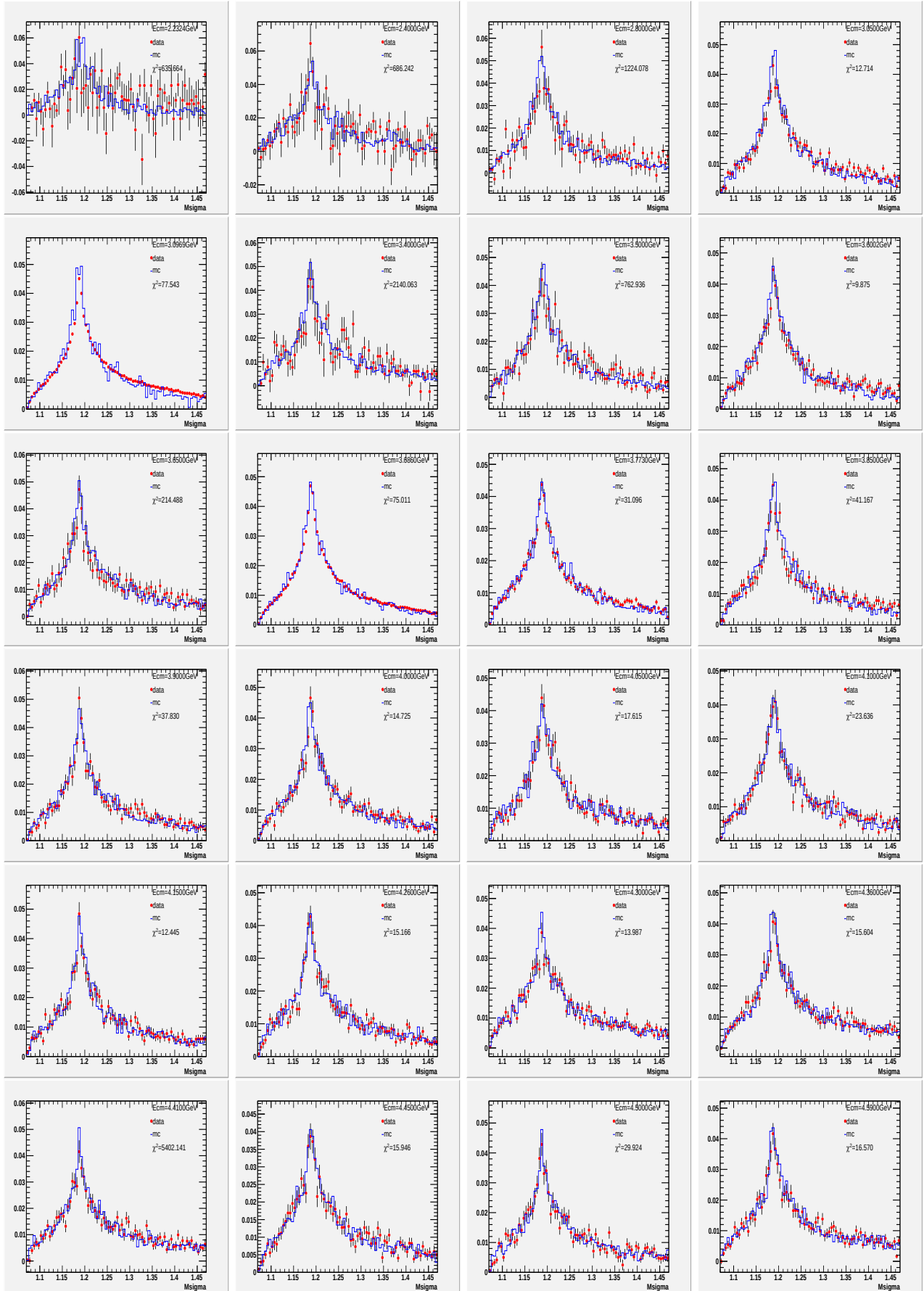


Fig. M_Σ .

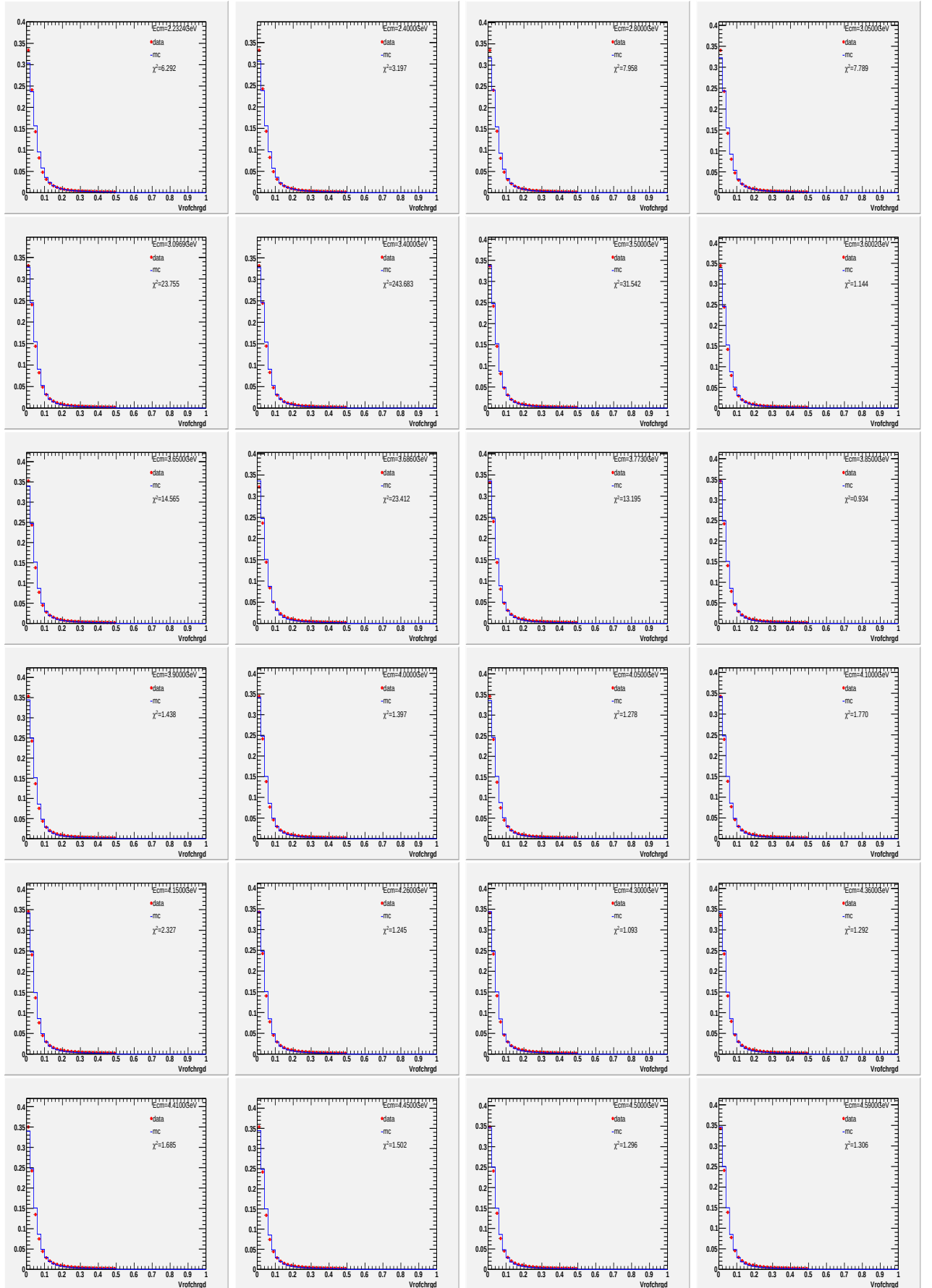


Fig. V_r

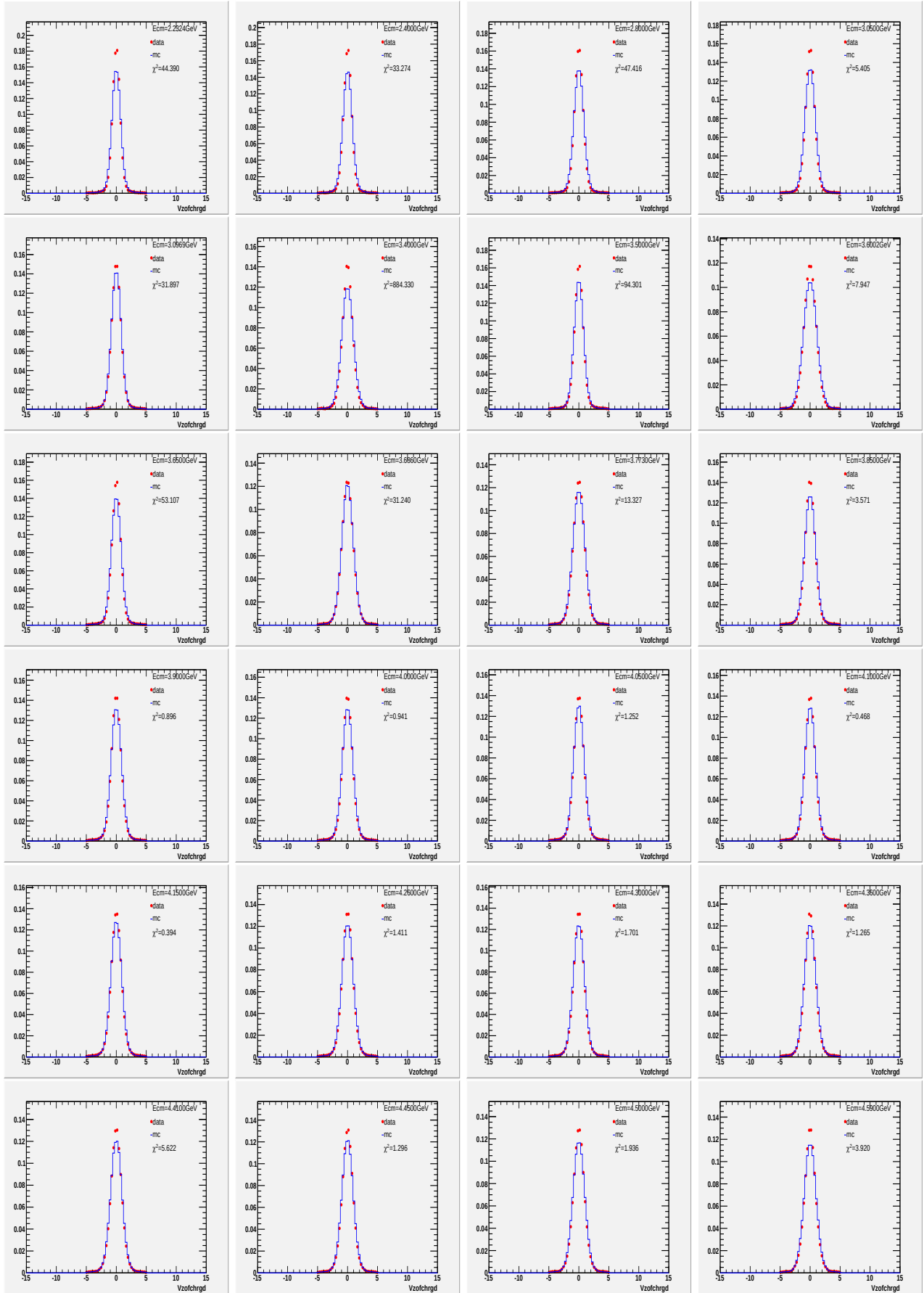


Fig. V_r

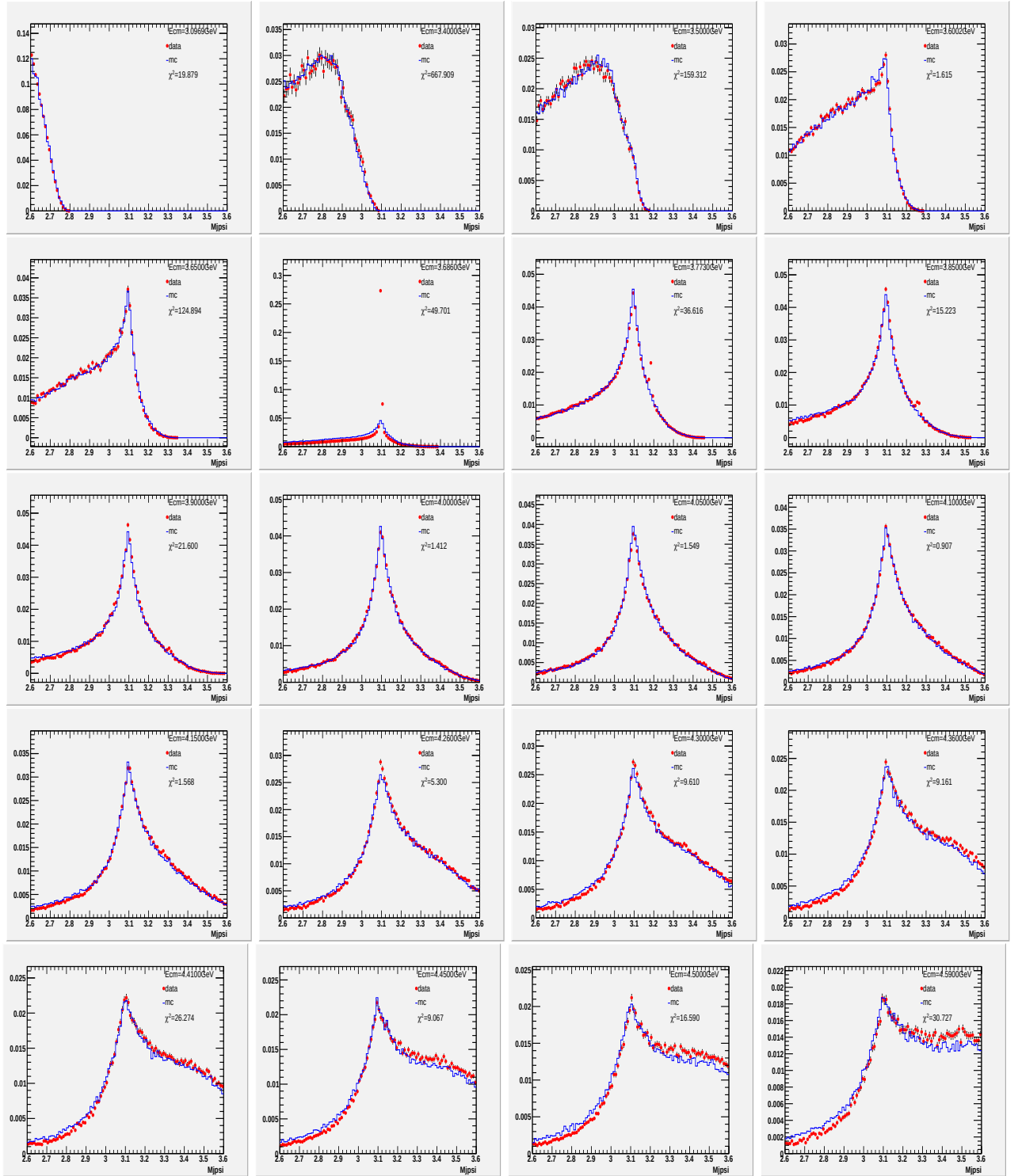


Fig. $M_{J\psi}$.

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