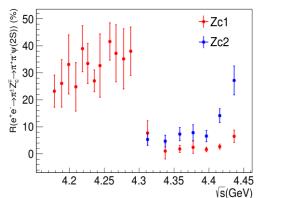


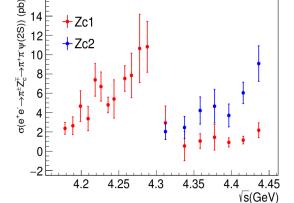


Partial Wave Analysis Amplitude Check

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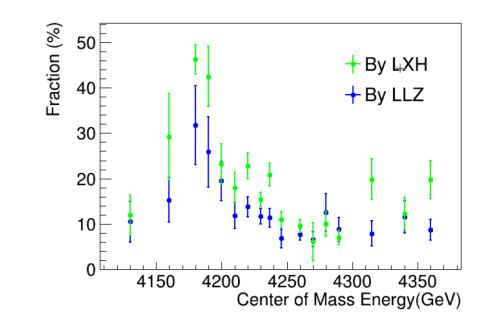


Covariant tensor formalism

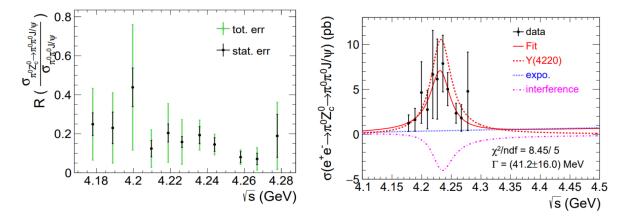


• The $Z_c(3900)$ cross section line shapes from $e^+e^- \rightarrow \pi^{+/0}\pi^{-/0}J/\psi$ and $e^+e^- \rightarrow \pi^+\pi^-\psi(3686)$ are different

• To validate the analysis of $e^+e^- \rightarrow \pi^+\pi^-\psi(3686)$, checks on amplitude and program are needed



Helicity formalism



Helicity amplitude construction

• Two body decay

 $a(J_a, \eta_a) \rightarrow b(J_b, \eta_b) + c(J_c, \eta_c)$

$$A_{\lambda_b,\lambda_c}^{J_a}(\theta,\phi;M) = N_{J_a} F_{\lambda_b,\lambda_c}^{J_a} D_{M,\lambda}^{J_a*}(\phi,\theta,0), (\lambda = \lambda_b - \lambda c)$$

• $F_{\lambda_b,\lambda_c}^J$ is helicity decay amplitude

$$F_{\lambda_b,\lambda_c}^{J_a} = \sum_{ls} \left(\frac{2l+1}{2J_a+1}\right)^{1/2} < l0s\lambda |J_a\lambda\rangle < s_b\lambda_b s_c - \lambda_c |s\lambda\rangle > G_{ls}^{J_a} r^l B_l(r)$$
$$G_{ls}^J = 4\pi \left(\frac{w}{p}\right)^{\frac{1}{2}} \langle JMls |\mathcal{M}| JM \rangle$$

- \blacktriangleright *G*_{LS} is LS coupling partial wave amplitude
- > With a definite set of helicity of (b,c), G_{LS} should be same
- > In fit, G_{LS} is a float complex parameter

• Helicity coupling amplitudes depend on the Lorentz factor for particles with spin 1 or higher

$$\xi_{s}(\lambda) \equiv f_{\lambda}^{(1)}(\gamma_{s}) = \begin{cases} [\chi^{(1)*}(\lambda) \cdot \omega(\lambda)] \\ 1 \quad \text{for } \lambda = \pm 1 \\ \gamma_{s} \quad \text{for } \lambda = 0, \end{cases} \qquad f_{\lambda}^{(2)}(\gamma_{s}) = \begin{cases} 1 \quad \text{for } \lambda = \pm 2 \\ \gamma_{s} \quad \text{for } \lambda = \pm 1 \\ \frac{2}{3}\gamma_{s}^{2} + \frac{1}{3} \quad \text{for } \lambda = 0. \end{cases}$$

Covariant helicity
coupling amplitude
$$F^{J}_{\lambda\nu} = \sum_{\ell S} g_{\ell S} A_{\ell S}(\lambda \nu),$$
where

$$A_{\ell S}(\lambda \nu) = \left(\frac{2\ell+1}{2J+1}\right)^{1/2} (\ell 0 S \delta | J\delta) (s\lambda \sigma - \nu | S\delta)$$
$$\times W^n r^{\ell} f^s_{\lambda}(\gamma_s) f^{\sigma}_{\nu}(\gamma_{\sigma}),$$

PhysRevD.57.431 (1998) by S. U. Chung

Helicity amplitude construction

Sequential decays: $Y \to \pi^+\pi^- J/\psi, J/\psi \to l^+l^-$

• Helicity formalism

$$\begin{split} \frac{d\sigma}{d\phi} &= \sum_{\lambda_Y,\Delta\lambda_l} \left| \sum_{\lambda_{R_i}} A_{R_i} \right|^2 \\ &= \sum_{\lambda_Y,\Delta\lambda_l} \left| \sum_{\lambda_{R_i},\lambda_{\psi}} F_1^J D_1^J \cdot F_2^J D_2^J \cdot F_3^J D_3^J \right|^2 \\ &= \sum_{\lambda_Y,\Delta\lambda_l} \left| \sum_{\lambda_{R_i},\lambda_{\psi}} (\sum_{LS} G_{LS} \cdot CG_1 \cdot B_l) D_1^J \cdot (\sum_{ls} G_{ls} \cdot CG_2 \cdot B_l) D_2^J \cdot F_3^J D_3^J \right|^2 \end{split}$$

• $J/\psi \rightarrow l^+l^-$ is included in helicity formalism

For the last step $J/\psi \to \ell^+ \ell^-$, at the relativistic limit, by QED calculation, $F_{1/2,1/2}^{J_{J/\psi}} = F_{-1/2,-1/2}^{J_{J/\psi}} \approx 0$. Here we define $\Delta \lambda_{\ell} = \lambda_{\ell^+} - \lambda_{\ell^-}$, we can see only $\Delta \lambda_{\ell} = \pm 1$ is allowed.

 $A_{Z_c}(\lambda_Y, \lambda_{Z_c}, \lambda_{\ell^+}, \lambda_{\ell^-}) = F^{J_Y}_{\lambda_{Z_c}, 0} D^{J_Y}_{\lambda_Y, \lambda_{Z_c}}(\theta_{Z_c}, \phi_{Z_c}) \cdot BW(Z_c) \cdot F^{J_{Z_c}}_{\lambda_{J/\psi}, 0} D^{J_{Z_c}}_{\lambda_{Z_c}, \lambda_{J/\psi}}(\theta_{J/\psi}, \phi_{J/\psi}) \\ \cdot F^{J_{J/\psi}}_{\lambda_{\ell^+}, \lambda_{\ell^-}} D^{J_{J/\psi}}_{\lambda_{J/\psi}, \lambda_{\ell^+} - \lambda_{\ell^-}}(\theta_{\ell^+}, \phi_{\ell^+}),$

$$egin{array}{rcl} Decay:Y& o\psi&f_0& ext{ } f_0& o\pi^+&\pi^-\ J^{PC}:1^{--}& o1^{--}&0^+& ext{ } 0^+& o0^-&0^- \end{array}$$

 $A_{R_f}(\lambda_Y, \lambda_{R_f}, \lambda_{\ell^+}, \lambda_{\ell^-}) = F_{\lambda_{R_f}, \lambda_{J/\psi}}^{J_Y} D_{\lambda_Y, \lambda_{R_f} - \lambda_{J/\psi}}^{J_Y}(\theta_{R_f}, \phi_{R_f}) \cdot BW(R_f) \cdot F_{0,0}^{J_{R_f}} D_{\lambda_{R_f},0}^{J_{R_f}}(\theta_{\pi^+}, \phi_{\pi^+}) \\ \cdot F_{\lambda_{\ell^+}, \lambda_{\ell^-}}^{J_{J/\psi}} D_{\lambda_{J/\psi}, \lambda_{\ell^+} - \lambda_{\ell^-}}^{J_{J/\psi}}(\theta_{\ell^+}, \phi_{\ell^+}),$

$$\frac{d\sigma}{d\phi} = \sum_{\lambda_Y, \Delta\lambda_l} |\sum_{\lambda_{Z_c}, \lambda_{R_f}} (A_{R_f} + e^{i\Delta\lambda_\ell \alpha_\ell (Z_c^+)} A_{Z_c^+} + e^{i\Delta\lambda_\ell \alpha_\ell (Z_c^-)} A_{Z_c^-})|^2$$

Helicity amplitude construction

• Sequential decay: $Y \to \pi^+\pi^- J/\psi$, $J/\psi \to l^+l^-$

$$\begin{aligned} Decay: Y & \to Z_c \quad \pi & L = 0(S - wave) \\ J^{PC}: 1^{--} & \to 1^+ & 0^- & L = 2(D - wave) \\ F_{1,0}^1 &= +g_{01}\sqrt{\frac{1}{3}}r^0 + g_{21}\sqrt{\frac{1}{6}}r^2 \\ &+ g_{01}\sqrt{\frac{1}{3}}r^0\gamma_s - g_{21}\sqrt{\frac{2}{3}}r^2\gamma_s & L = 0(S - wave) \\ 1^+ &\to 1^{--}0^- & L = 0(S - wave) \\ &1^+ \to 1^{--}0^- & L = 2(D - wave) \\ &F_{1,0}^1 &= +g_{01}\sqrt{\frac{1}{3}}r^0 + g_{21}\sqrt{\frac{1}{6}}r^2 \\ &+ g_{01}\sqrt{\frac{1}{3}}r^0\gamma_s - g_{21}\sqrt{\frac{2}{3}}r^2\gamma_s & \end{bmatrix} \end{aligned}$$

• Four components: SS, SD, DS and DD

Amplitude construction

Sequential decays: $Y \rightarrow \pi^+\pi^- J/\psi$

• Covariant tensor formalism

$$egin{aligned} &A=\phi_{\mu}(m_{1})\omega_{
u}^{*}(m_{2})A^{\mu
u}=\phi_{\mu}(m_{1})\omega_{
u}^{*}(m_{2})\sum_{i}\Lambda_{i}U_{i}^{\mu
u}\ &\sum_{m_{1}}^{2}\phi_{\mu}(m_{1})\phi_{\mu'}^{*}(m_{1})=\delta_{\mu\mu'}(\delta_{\mu1}+\delta_{\mu2})\ &\sum_{m_{2}=1}^{3}\omega_{
u}(m_{2})\omega_{
u'}^{*}(m_{2})=-g_{
u
u'}+rac{p_{(\psi)
u}p_{(\psi)
u'}}{p_{\psi}^{2}}\equiv- ilde{g}_{
u
u'}(p_{(\psi)}) \end{aligned}$$

$$\begin{split} \frac{d\sigma}{d\Phi_n} \propto \frac{1}{2} \sum_{m_1}^2 \sum_{m_2}^3 \phi_\mu(m_1) \omega_\nu^*(m_2) A^{\mu\nu} \phi_{\mu'}^*(m_1) \omega_{\nu'}(m_2) A^{*\mu'\nu'} \\ \frac{d\sigma}{d\Phi_n} \propto -\frac{1}{2} \sum_{\mu=1}^2 \tilde{g}_{\nu\nu'}(p_{(\psi)}) A^{\mu\nu} A^{*\mu\nu'} \\ = -\frac{1}{2} \sum_{i,j} \Lambda_i \Lambda_j^* \sum_{\mu=1}^2 U_i^{\mu\nu} \tilde{g}_{\nu\nu'}(p_{(\psi)}) U_j^{*\mu\nu'} \end{split}$$

• $U^{\mu\nu}$ is the partial wave amplitude constructed according to LS coupling $U^{\mu\nu} = (A_{LS})(A_{ls})$

$$U^{\mu
u}_{(Y o Z_c^\pm\pi^\mp)SS}= ilde{g}^{\mu
u}_{(Z_c^+)}f^{(Z_c^+)}_{(01)}+ ilde{g}^{\mu
u}_{(Z_c^-)}f^{(Z_c^-)}_{(02)}$$

$$U^{\mu
u}_{(Y o Z^{\pm}_c\pi^{\mp})SD} = ilde{t}^{(2)\mu
u}_{(\psi\pi^+)}f^{(Z^+_c)}_{(01)} + ilde{t}^{(2)\mu
u}_{(\psi\pi^-)}f^{(Z^-_c)}_{(02)}$$

$$U^{\mu
u}_{(Y o Z_c^{\pm}\pi^{\mp})DS} = ilde{T}^{(2)\mu\lambda}_{(Z_c^{+}\pi^{-})} ilde{g}_{(Z_c^{+})\lambda\sigma} g^{\sigma
u} f^{(Z_c^{+})}_{(01)} + ilde{T}^{(2)\mu\lambda}_{(Z_c^{-}\pi^{+})} ilde{g}_{(Z_c^{-})\lambda\sigma} g^{\sigma
u} f^{(Z_c^{-})}_{(02)}$$

$$U^{\mu\nu}_{(Y \to Z_c^{\pm} \pi^{\mp})DD} = \tilde{T}^{(2)\mu\lambda}_{(Z_c^{+} \pi^{-})} \tilde{t}^{(2)}_{(\psi\pi^{+})\lambda\sigma} g^{\sigma\nu} f^{(Z_c^{-})}_{(01)} + \tilde{T}^{(2)\mu\lambda}_{(Z_c^{-} \pi^{+})} \tilde{t}^{(2)}_{(\psi\pi^{-})\lambda\sigma} g^{\sigma\nu} f^{(Z_c^{-})}_{(02)}$$

$$egin{array}{rcl} Decay: Y & o\psi & f_0 & f_0 & o\pi^+ & \pi^- \ J^{PC}: 1^{--} & o1^{--} & 0^+ & 0^+ & o0^- & 0^- \end{array}$$

$$egin{aligned} U^{\mu
u}_{(Y
ightarrow\psi(2S)f_0)SS} &= \langle\psi f_0|01
angle = g^{\mu
u} f^{(f_0)}_{(12)} \ U^{\mu
u}_{(Y
ightarrow\psi(2S)f_0)DS} &= \langle\psi f_0|21
angle = ilde{T}^{(2)\mu
u}_{(\psi f_0)} f^{(f_0)}_{(12)} \end{aligned}$$

Two points:

- $J/\psi \rightarrow l^+l^-$ is included in the helicity formalism and not in covariant tensor formalism
- Lorentz factor is considered in covariant helicity formalism

$$\begin{split} F^{J_a}_{\lambda_b,\lambda_c} &= \sum_{r} (\frac{2l+1}{2J_a+1})^{1/2} < l0s\lambda | J_a\lambda > < s_b\lambda_b s_c - \lambda_c | s\lambda > G^{J_a}_{ls} r^l B_l(r) \\ F^J_{\lambda\nu} &= \sum_{\ell S} g_{\ell S} A_{\ell S}(\lambda \nu), \end{split}$$

where

$$A_{\ell S}(\lambda \nu) = \left(\frac{2\ell+1}{2J+1}\right)^{1/2} (\ell 0 S \delta | J\delta) (s\lambda \sigma - \nu | S\delta)$$
$$\times W^n r^{\ell} f^s_{\lambda}(\gamma_s) f^{\sigma}_{\nu}(\gamma_{\sigma}),$$

$$\frac{d\sigma}{d\phi} = \sum_{\lambda_Y, \Delta\lambda_l} |\sum_{\lambda_{Z_c}, \lambda_{R_f}} (A_{R_f} + e^{i\Delta\lambda_\ell \alpha_\ell (Z_c^+)} A_{Z_c^+} + e^{i\Delta\lambda_\ell \alpha_\ell (Z_c^-)} A_{Z_c^-})|^2$$

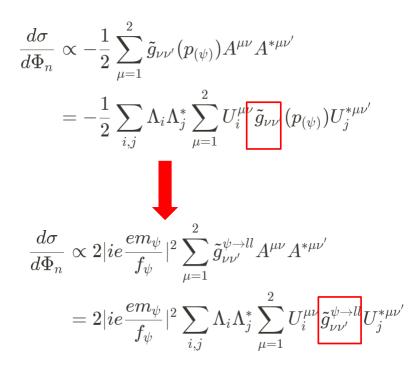
$$egin{aligned} rac{d\sigma}{d\Phi_n} \propto -rac{1}{2} \sum_{\mu=1}^2 ilde{g}_{
u
u'}(p_{(\psi)}) A^{\mu
u} A^{*\mu
u'} \ &= -rac{1}{2} \sum_{i,j} \Lambda_i \Lambda_j^* \sum_{\mu=1}^2 U_i^{\mu
u} ilde{g}_{
u
u'}(p_{(\psi)}) U_j^{*\mu
u'} \end{aligned}$$

Amplitude construction

$$A = \phi_{\mu}(m_1)\omega_{
u}^*(m_2)A^{\mu
u} = \phi_{\mu}(m_1)\omega_{
u}^*(m_2)\sum_i \Lambda_i U_i^{\mu
u}$$

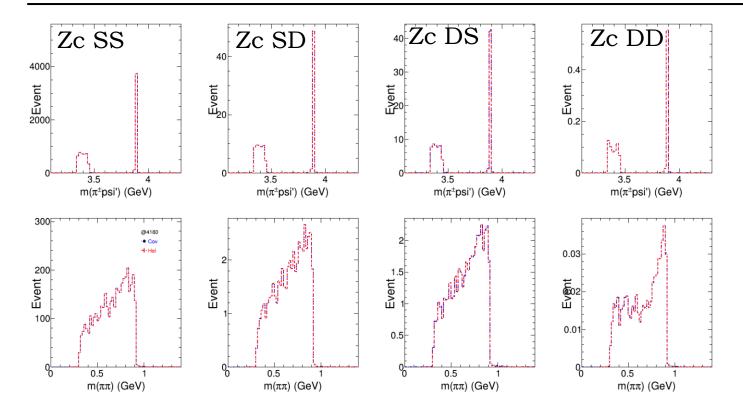
 ${\cal B}=ie\omega_eta(m_2)ar u_{e^-}\gamma^eta
u_{e^+}rac{em_\psi}{f_\psi}$

- $\omega_{\beta}(m_2)$ 是 ψ 的极化矢量
- $f_{\psi} = 11.2$ 是常数
- $\bar{u}_{e^-}(\nu_{e^+})$ 是电子的极化矢量



$$\begin{split} & \int l\psi \rightarrow e^{+}e^{-} \\ & \mathcal{M} = ie \ \bar{u}_{e} \ \gamma^{\nu} \mathcal{V}e^{+} \cdot \frac{e^{m_{n_{\nu}}}}{f_{n_{\nu}}} \ \xi_{n_{\nu}} \\ & e^{+}e^{-} \\$$

Test with MC



- The simplest MC sample:
 - Only Zc SS component
 - No BW width
 - Generated by helicity formalism
- Fit in two formalisms

With $J/\psi \rightarrow l^+l^-$ Cov: SS wave fraction = 0.976046 Hel: SS wave fraction = 0.975296 Hel: SD wave fraction = 0.0130789 Hel: DS wave fraction = 0.0130789 Hel: DS wave fraction = 0.0114711 Hel: DD wave fraction = 0.000154264

Without		fraction fraction	
${\rm J}/\psi \to l^+ l^-$		fraction fraction	0.0834154 0.272534

The test result supports the necessity of $J/\psi \rightarrow l^+l^-$

Test with MC

- The fractions are consistent between covariant tensor formalism and helicity formalism when $J/\psi \rightarrow l^+l^-$ is included
- Different (LS) wave components for two formalisms are in match
- $J/\psi \rightarrow l^+l^-$ part is necessary

With $J/\psi \rightarrow l^+ l^-$ SS wave: Cov = 0.11943Hel = 0.117659Cov = 0.108338Zc mass=3880 MeV Hel = 0.0220939SD wave: Cov = 0.0223932Cov = 0.143104Width = 1 MeVDS wave: Cov = 0.723714Hel = 0.725169Cov = 0.321651DD wave: Cov = 0.134462Hel = 0.135078Cov = 0.426907Cov = 0.0721274Zc mass=3885 MeV SS wave: Cov = 0.140743Hel = 0.137783SD wave: Cov = 0.0218467Hel = 0.0217328Cov = 0.283996Width = 1 MeVCov = 0.130078DS wave: Cov = 0.724729Hel = 0.725696Cov = 0.513799DD wave: Cov = 0.112681Hel = 0.114788SS wave: Cov = 0.126908Hel = 0.122224Cov = 0.0780388Zc mass=3890 MeV Hel = 0.0222862SD wave: Cov = 0.0227674Cov = 0.164744Width = 1 MeVDS wave: Cov = 0.722514Hel = 0.724973Cov = 0.243955DD wave: Cov = 0.127811Hel = 0.130517Cov = 0.513263

MC sample: Four components: SS+SD+DS+DD BW: different Zc mass and 1 MeV width

Without $J/\psi \rightarrow l^+l^-$

Zc3900 MC sample: only SS component, BW function has width

Without $J/\psi \rightarrow l^+l^-$ in covariant tensor formalism

Cov: SS wave fraction = 0.292975	Hel: SS wave fraction = 0.905202
Cov: SD wave fraction = 0.500984	Hel: SD wave fraction = 0.00800091
Cov: DS wave fraction = 0.0744949	Hel: DS wave fraction = 0.0859795
Cov: DD wave fraction = 0.131546	Hel: DD wave fraction = 0.000817607

With $J/\psi \rightarrow l^+l^-$ in covariant tensor formalism

Cov: SS wave fraction = 0.930268	Hel: SS wave fraction = 0.905202
Cov: SD wave fraction = 0.00704404	Hel: SD wave fraction = 0.00800091
Cov: DS wave fraction = 0.062201	Hel: DS wave fraction = 0.0859795
Cov: DD wave fraction = 0.000487272	Hel: DD wave fraction = 0.000817607

- The test result supports the necessity of $J/\psi \rightarrow l^+ l^-$
- Differences appear in fractions of two formalisms
- Invariant scattering amplitude has mass-dependent term? This term is treated as a constant?

$$G^J_{ls} = 4\pi \left(rac{w}{p}
ight)^{rac{1}{2}} \langle JMls | \mathcal{M} | JM
angle$$

5.3.1 $\psi \rightarrow \rho^0 \pi^0 \rightarrow \pi^+ \pi^- \pi^0$

第一个例子,我们采用相空间产生子,产生了 $\psi \to \rho^0 \pi^0 \to \pi^+ \pi^- \pi^0$ 过程的相空间 事例,这时Breit-Wigner分布已经包含在内了,但是不包含势垒因子和角分布。

5.3.1.0.1 截面 我们先把每个事例的截面计算出来(去掉Breit-Wigner),看看两种 理论描述的计算结果,并求它们的比值,下面是打印出的数值信息: Helicity计算的截面、张量方法计算的截面和它们的比值。它们在计算精度内完全一致。

Helicity: 0.	.353986	L-S: 0.353986	Ratio: 1.000000
Helicity: 0.	.795525	L-S: 0.795525	Ratio: 1.000000
Helicity: 0.	.410571	L-S: 0.410571	Ratio: 1.000000
Helicity: 0.	.004812	L-S: 0.004812	Ratio: 1.000000
Helicity: 0.	.704383	L-S: 0.704383	Ratio: 1.000000
Helicity: 0.	.480020	L-S: 0.480020	Ratio: 1.000000

不考虑能量依赖, $\psi \to \rho^0 \pi^0$ 道的完整角分布表达式为

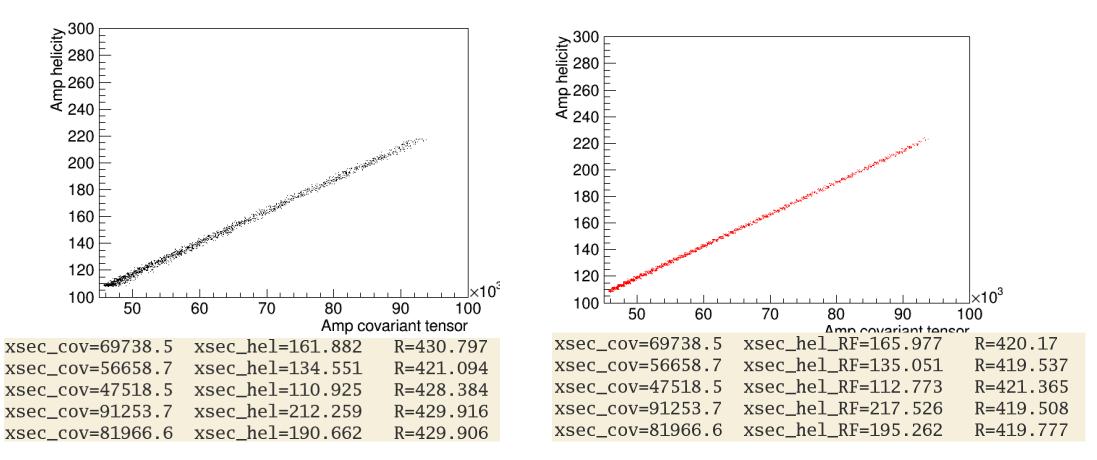
$$(\frac{d\sigma}{d\Phi})_H \propto |\mathcal{C}_1|^2 \sum_M \left| D_{M1}^{1\star}(\theta\phi) D_{10}^{1\star}(\theta'\phi') - D_{M-1}^{1\star}(\theta\phi) D_{-10}^{1\star}(\theta'\phi') \right|^2$$
$$\propto \frac{1}{2} [(1+\cos^2\theta)\sin^2\theta' + \sin^2\theta\sin^2\theta'\cos 2\phi']$$

$$\frac{d\sigma}{d\Phi} \propto \frac{|\mathcal{C}_2|^2}{M_{\psi}^2} \sum_{\mu=1}^2 U_{\rho}^{\mu} U_{\rho}^{\star \mu}
= |\mathcal{C}_2|^2 \epsilon_{ijk} p_1^j p_2^k \epsilon^{imn} p_{1m} p_{2n}
= |\mathcal{C}_2|^2 (\vec{p}_1 \times \vec{p}_2)_{xy}^2.$$

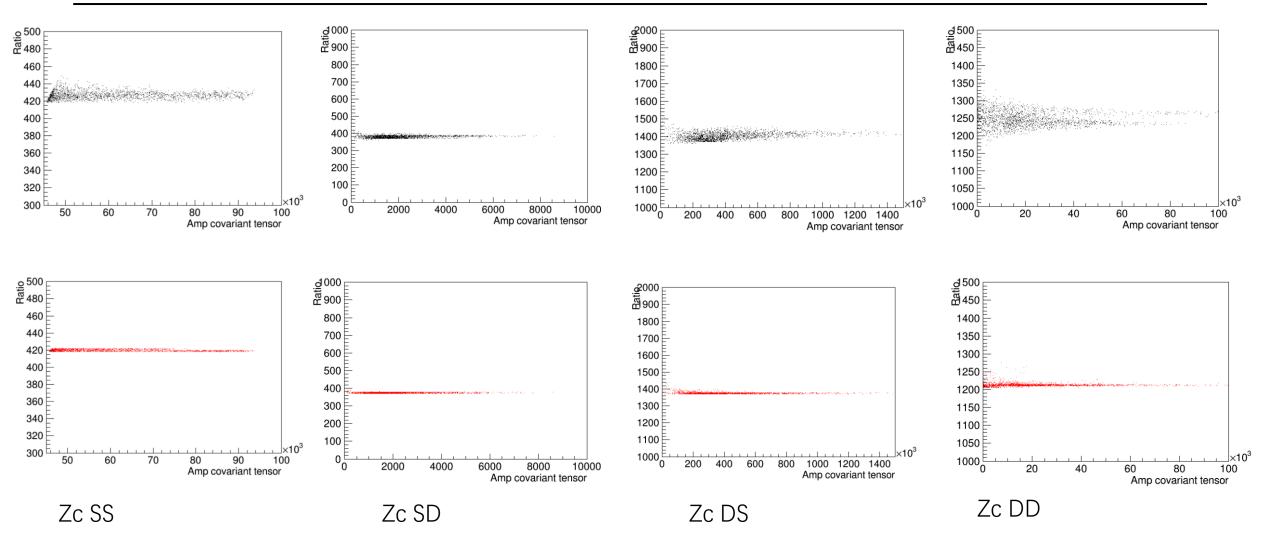
分类号 密级 UDC 编号	$D \to \rho^0 \pi^0 \to \pi^+ \pi^- \pi^0, \ \phi \pi^0 \to K^+ K^- \pi^0$ Both P wave for two steps	$D \rightarrow f_2 \pi^0 \rightarrow \pi^+ \pi^- \pi^0$ Both D wave for two steps
中国高等科技中心 博士后研究报告 <u>分波分析方法</u>	Helicity formali $Z = D_{0\lambda}^{0}(0, \theta_{0}\phi_{0})f_{\lambda}(\gamma)D_{\lambda0}^{*1}(0\theta\phi)$ $= \gamma d_{00}^{1}(\theta) \qquad \lambda = 0$ $= \gamma \cos \theta ,$	Solution $Z = D_{0\lambda}^{0}(0, \theta_{0}\phi_{0})f_{\lambda}(\gamma)D_{\lambda0}^{*2}(0\theta\phi)$ $= (\frac{2}{3}\gamma^{2} + \frac{1}{3})d_{00}^{2}(\theta)$ $= (\frac{2}{3}\gamma^{2} + \frac{1}{3})(\frac{3\cos^{2}\theta - 1}{2}),$
	Covariant forma	alism:
合作导师 <u>叶铭汉 院士,朱永生 研究员</u> 工作完成日期 <u>2004年7月— 2006年8月</u> 报告提交日期 <u>2006年8月</u> 中国高等科技中心(北京) 2006年8月	$Z = -2p(p_a - p_b)^3$ = $-2p[(\gamma\beta E^* - \gamma q_3) - (\gamma\beta E^* + \gamma q_3)]$ = $4p\gamma(q\cos\theta).$ = $4pq\gamma\cos\theta$,	$Z = \tilde{T}^{(2)}_{\mu\nu}(p_a + p_b + p_c)\tilde{t}^{\mu\nu(2)}(p_a + p_b)$ $= \left[(p_a + p_b - p_c)_i(p_a + p_b - p_c)_j - \frac{1}{3}\delta_{ij}(p_a + p_b - p_c)^2 \right]$ $\left[(p_a - p_b)^i(p_a - p_b)^j - \frac{1}{3}\delta^{ij}(p_a - p_b)^2 \right]$ $= \left[(p_a + p_b - p_c) \cdot (p_a - p_b) \right]^2 + \frac{1}{3}(p_a + p_b - p_c)^2(p_a - p_b)^2$ $- \frac{1}{3}(p_a + p_b - p_c)^2(p_a - p_b)^2 - \frac{1}{3}(p_a + p_b - p_c)^2(p_a - p_b)^2$ $= \left[(p_a + p_b - p_c) \cdot (p_a - p_b) \right]^2 - \frac{1}{3}(p_a + p_b - p_c)^2(p_a - p_b)^2$ $= \left[(p_a + p_b - p_c) \cdot (p_a - p_b) \right]^2 - \frac{1}{3}(p_a + p_b - p_c)^2(p_a - p_b)^2$ $= \left[(p_a + p_b - p_c) \cdot (p_a - p_b) \right]^2 - \frac{1}{3}(p_a + p_b - p_c)^2(p_a - p_b)^2$

 $= \left[\frac{64}{3} \times p^2 q^2 \left[(\frac{2}{3} \gamma^2 + \frac{1}{3}) (\frac{3 \cos^2 \theta}{2}) - \frac{1}{2} \right] \right]$

MC: $Y \rightarrow \pi^+\pi^- J/\psi$, $J/\psi \rightarrow l^+l^-$ phsp MC Amplitude: only Zc SS component



- The ratio is close to a constant as well
- The relativistic factor also makes the ratio more stable



其中 $z = [qd]^2$, $z_0 = [q_0d]^2$, d 是相互作用的典型力程。

L	$B_L(q)$
0	1
1	$\sqrt{rac{2z}{1+z}}$
2	$\sqrt{\frac{13z^2}{(z-3)^2+9z}}$
3	$\sqrt{\frac{277z^3}{z^3+6z^2+45z+225}}$
4	$\sqrt{\frac{12746z^4}{z^4 + 10z^3 + 135z^2 + 1575z + 11025}}$

);

$$p^{\alpha} = (W;0,0,0),$$

$$p^{\alpha} = (W;0,0,0),$$

$$q^{\alpha} = (q_{0};0,0,q) = (\gamma_{s}m;0,0,\gamma_{s}\beta_{s}m),$$

$$A_{\ell S}(\lambda \nu) = \left(\frac{2\ell+1}{2J+1}\right)^{1/2} (\ell 0S\delta|J\delta)(s\lambda \sigma - \nu|S\delta)$$

$$\times W r^{2} f^{S}_{\lambda}(\gamma_{s}) f^{\sigma}_{\nu}(\gamma_{\sigma}),$$
where
$$W = q_{0} + k_{0}, \quad q_{0} = \sqrt{m^{2}+q^{2}}, \quad k_{0} = \sqrt{m^{2}+q^{2}},$$

$$W = q_{0} + k_{0}, \quad q_{0} = \sqrt{m^{2}+q^{2}}, \quad k_{0} = \sqrt{m^{2}+q^{2}},$$

$$W = q_{0} + k_{0}, \quad q_{0} = \sqrt{m^{2}+q^{2}}, \quad k_{0} = \sqrt{m^{2}+q^{2}},$$

$$W = q_{0} + k_{0}, \quad q_{0} = \sqrt{m^{2}+q^{2}}, \quad k_{0} = \sqrt{m^{2}+q^{2}},$$

$$W = q_{0} + k_{0}, \quad q_{0} = \sqrt{m^{2}+q^{2}}, \quad k_{0} = \sqrt{m^{2}+q^{2}},$$

$$W = q_{0} + k_{0}, \quad q_{0} = \sqrt{m^{2}+q^{2}}, \quad k_{0} = \sqrt{m^{2}+q^{2}}, \quad k_{0} = \sqrt{m^{2}+q^{2}},$$

 $p^{\alpha} = (W; 0, 0, 0),$

 $(0,q) = (\gamma_s m; 0, 0, \gamma_s \beta_s m),$

(17)

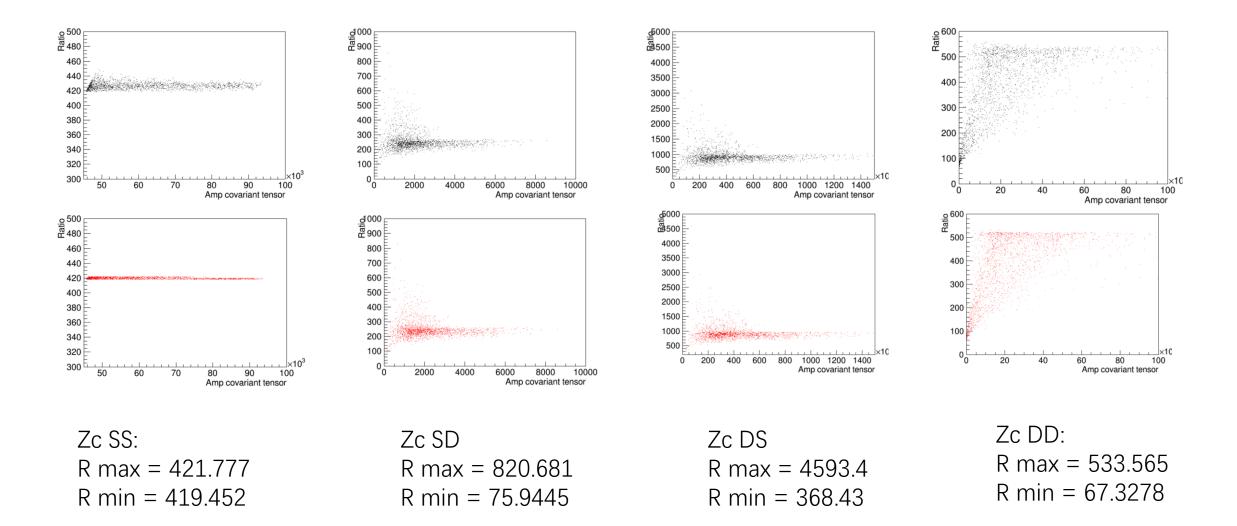
 $=(q_0-k_0;0,0,2q),$

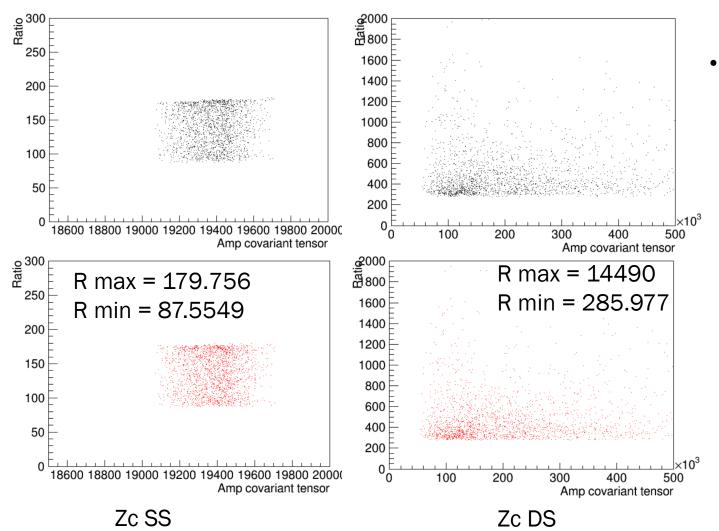
 $k_0, \quad q_0 = \sqrt{m^2 + q^2}, \quad k_0 = \sqrt{\mu^2 + q^2}, \quad \text{and}$ q_0^{-1} r = q - k, the wave functions in the J rest frame are given by

$$Q_{abc}^2 = \frac{(s_a + s_b - s_c)^2}{4s_a} - s_b$$

Here Q_0 is a hadron "scale" parameter $Q_0 =$ 0.197321/R GeV/c, where R is the radius of the centrifugal barrier in fm. We remark that in these Blatt-Weisskopf factors, the approximation is made that the centrifugal barrier may be replaced by a square well of radius R.

$$\begin{split} B_1(Q_{abc}) &= \sqrt{\frac{2}{Q_{abc}^2 + Q_0^2}}, \\ B_2(Q_{abc}) &= \sqrt{\frac{13}{Q_{abc}^4 + 3Q_{abc}^2Q_0^2 + 9Q_0^4}}, \\ B_3(Q_{abc}) &= \sqrt{\frac{277}{Q_{abc}^6 + 6Q_{abc}^4Q_0^2 + 45Q_{abc}^2Q_0^4 + 225Q_0^6}}, \\ B_4(Q_{abc}) &= \sqrt{\frac{12746}{Q_{abc}^8 + 10Q_{abc}^6Q_0^2 + 135Q_{abc}^4Q_0^4 + 1575Q_{abc}^2Q_0^6 + 11025Q_0^8}}. \end{split}$$





Without $J/\psi \rightarrow l^+l^-$, even the amplitude ratio of Zc SS component is not a constant

Summary

- In this report, tests are performed to demonstrate the consistency between two commonly used formalisms: helicity formalism and covariant tensor formalism
- For example process $Y \to \pi^+ \pi^- J/\psi \to \pi^+ \pi^- l^+ l^-$, test results support the necessicity of $J/\psi \to l^+ l^-$
 - → How about other case? $\phi \to KK$? $\omega \to \pi^+\pi^-\pi^0$?

For *Y* → $\pi^+\pi^-\psi(3686)$, only use $\psi(3686) \rightarrow \pi^+\pi^- J/\psi \rightarrow \pi^+\pi^- l^+ l^-$?

- Even for L ≥ 2(D wave or higher), amplitudes constructed from helicity formalism and covariant tensor formalism are consistent
 - Which one describes physical process better? Or Is it possible to search a better amplitude construction method?

BACKUP

