• Two body decay

 $\left| rac{d\sigma}{d\phi} = \sum_{\lambda_V = \Lambda \lambda_I} \left| \sum_{\lambda_{R_i}} A_{R_i} \right|$

 $a(J_a, \eta_a) \rightarrow b(J_b, \eta_b) + c(J_c, \eta_c)$

 $= \sum \left| \sum F_1^J D_1^J \cdot F_2^J D_2^J \cdot F_3^J D_3^J \right|$

$$A_{\lambda_b,\lambda_c}^{J_a}(\theta,\phi;M) = N_{J_a} F_{\lambda_b,\lambda_c}^{J_a} D_{M,\lambda}^{J_a*}(\phi,\theta,0), (\lambda = \lambda_b - \lambda c)$$

• Sequential decays: $Y \to \pi^+\pi^- J/\psi$, $J/\psi \to l^+l^-$

$$\begin{split} & F_{\lambda_b,\lambda_a}^{J_a} \text{ is including decay dimplitude} \\ & F_{\lambda_b,\lambda_c}^{J_a} = \sum_{ls} (\frac{2l+1}{2J_a+1})^{1/2} < l0s\lambda | J_a\lambda > < s_b\lambda_b s_c - \lambda_c | s\lambda > G_{ls}^{J_a} r^l B_l(r) \\ & F_{\lambda_1\lambda_2}^J = \eta \eta_1 \eta_2(-)^{J-s_1-s_2} F_{-\lambda_1-\lambda_2}^J \\ & G_{ls}^J = 4\pi \left(\frac{w}{p}\right)^{\frac{1}{2}} \langle JMls | \mathcal{M} | JM \rangle \end{split}$$

is helicity decay amplitude

- G_{LS} is LS coupling partial wave amplitude
- With a definite set of helicity of (b,c), G_{LS} should be same
- In fit, *G*_{LS} is float parameter
- To obtain the contribution of a LS wave component

 $= \sum_{\lambda_Y, \Delta\lambda_l} \left| \sum_{\lambda_{R_i}, \lambda_{\psi}} (\sum_{LS} G_{LS} \cdot CG_1 \cdot B_l) D_1^J \cdot (\sum_{ls} G_{ls} \cdot CG_2 \cdot B_l) D_2^J \cdot F_3^J D_3^J \right|^2 \longrightarrow \sum_{\lambda_Y, \Delta\lambda_l} \left| \sum_{\lambda_{R_i}, \lambda_{\psi}} (G_{LS} \cdot CG_1 \cdot B_l) D_1^J \cdot (G_{ls} \cdot CG_2 \cdot B_l) D_2^J \cdot F_3^J D_3^J \right|^2$

 F^{J}

For the last step $J/\psi \to \ell^+ \ell^-$, at the relativistic limit, by QED calculation, $F_{1/2,1/2}^{J_{J/\psi}} = F_{-1/2,-1/2}^{J_{J/\psi}} \approx 0$. Here we define $\Delta \lambda_{\ell} = \lambda_{\ell^+} - \lambda_{\ell^-}$, we can see only $\Delta \lambda_{\ell} = \pm 1$ is allowed.

$$\sum_{\lambda_Y,\Delta\lambda_l} |\sum_{\lambda_{R_i},\lambda_{\psi}} (G_{LS} \cdot CG_1 \cdot B_l) D_1^J \cdot (G_{ls} \cdot CG_2 \cdot B_l) D_2^J \cdot F_3^J D_3^J |^2$$

$$egin{array}{ll} Z_c
ightarrow \psi & \pi \ 1^+
ightarrow 1^{--} 0^- \end{array}$$

$$F_{1,0}^{1} = +g_{01}\sqrt{\frac{1}{3}}r^{0} + g_{21}\sqrt{\frac{1}{6}}r^{2}$$
$$F_{0,0}^{1} = +g_{01}\sqrt{\frac{1}{3}}r^{0}\gamma_{s} - g_{21}\sqrt{\frac{2}{3}}r^{2}\gamma_{s}$$

$$A = \phi_\mu(m_1) \omega^*_
u(m_2) A^{\mu
u} = \phi_\mu(m_1) \omega^*_
u(m_2) \sum_i \Lambda_i U_i^{\mu
u}$$

$$egin{aligned} rac{d\sigma}{d\Phi_n} \propto -rac{1}{2} \sum_{\mu=1}^2 ilde{g}_{
u
u'}(p_{(\psi)}) A^{\mu
u} A^{*\mu
u'} \ &= -rac{1}{2} \sum_{i,j} \Lambda_i \Lambda_j^* \sum_{\mu=1}^2 U_i^{\mu
u} ilde{g}_{
u
u'}(p_{(\psi)}) U_j^{*\mu
u'} \end{aligned}$$

$$U^{\mu\nu} = (A_{LS})(A_{ls})$$

• $U^{\mu\nu}$ is the partial wave amplitude constructed according to LS coupling

$$egin{array}{rcl} Decay: Y & o\psi & f_0 & f_0 & o\pi^+ & \pi^- \ J^{PC}: 1^{--} & o1^{--} & 0^+ & 0^+ & o1^- & 0^- \end{array}$$

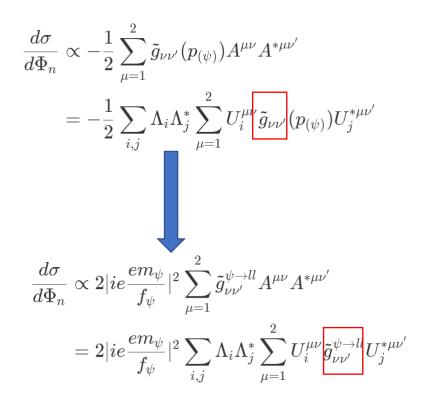
$$egin{aligned} U^{\mu
u}_{(Y o\psi(2S)f_0)SS} &= \langle\psi f_0|01
angle &= g^{\mu
u}f^{(f_0)}_{(12)} \ U^{\mu
u}_{(Y o\psi(2S)f_0)DS} &= \langle\psi f_0|21
angle &= ilde{T}^{(2)\mu
u}_{(\psi f_0)}f^{(f_0)}_{(12)} \end{aligned}$$

$$\begin{split} Decay &: Y \to Z_c \quad \pi \qquad Z_c \to \psi \quad \pi \\ J^{PC} &: 1^{--} \to 1^+ \quad 0^- \qquad 1^+ \to 1^{--} 0^- \\ U^{\mu\nu}_{(Y \to Z_c^{\pm} \pi^{\mp})SS} &= \tilde{g}^{\mu\nu}_{(Z_c^+)} f^{(Z_c^+)}_{(01)} + \tilde{g}^{\mu\nu}_{(Z_c^-)} f^{(Z_c^-)}_{(02)} \\ U^{\mu\nu}_{(Y \to Z_c^{\pm} \pi^{\mp})SD} &= \tilde{t}^{(2)\mu\nu}_{(\psi\pi^+)} f^{(Z_c^+)}_{(01)} + \tilde{t}^{(2)\mu\nu}_{(\psi\pi^-)} f^{(Z_c^-)}_{(02)} \\ U^{\mu\nu}_{(Y \to Z_c^{\pm} \pi^{\mp})DS} &= \tilde{T}^{(2)\mu\lambda}_{(Z_c^+ \pi^-)} \tilde{g}_{(Z_c^+)\lambda\sigma} g^{\sigma\nu} f^{(Z_c^+)}_{(01)} + \tilde{T}^{(2)\mu\lambda}_{(Z_c^- \pi^+)} \tilde{g}_{(Z_c^-)\lambda\sigma} g^{\sigma\nu} f^{(Z_c^-)}_{(02)} \\ U^{\mu\nu}_{(Y \to Z_c^{\pm} \pi^{\mp})DD} &= \tilde{T}^{(2)\mu\lambda}_{(Z_c^+ \pi^-)} \tilde{t}^{(2)}_{(\psi\pi^+)\lambda\sigma} g^{\sigma\nu} f^{(Z_c^+)}_{(01)} + \tilde{T}^{(2)\mu\lambda}_{(Z_c^- \pi^+)} \tilde{t}^{(2)}_{(\psi\pi^-)\lambda\sigma} g^{\sigma\nu} f^{(Z_c^-)}_{(02)} \end{split}$$

$$A=\phi_{\mu}(m_{1})\omega_{
u}^{*}(m_{2})A^{\mu
u}=\phi_{\mu}(m_{1})\omega_{
u}^{*}(m_{2})\sum_{i}\Lambda_{i}U_{i}^{\mu
u}$$

 ${\cal B}=ie\omega_eta(m_2)ar u_{e^-}\gamma^eta
u_{e^+}rac{em_\psi}{f_\psi}$

- $\omega_{\beta}(m_2)$ 是 ψ 的极化矢量
- $f_{\psi} = 11.2$ 是常数
- $\bar{u}_{e^-}(\nu_{e^+})$ 是电子的极化矢量



• F0(500): MC sample contains only DS component

Without $J/\psi \rightarrow l^+l^-$ in covariant tensor formalism

Cov: SS wave fraction = 0.588728 Hel: SS wave fraction = 0.00091914 Cov: DS wave fraction = 0.411272 Hel: DS wave fraction = 0.999081

With $J/\psi \rightarrow l^+l^-$ in covariant tensor formalism

Cov:	SS wave fraction = 0.0535918	Hel: SS wave fraction = 0.00091914
Cov:	DS wave fraction = 0.946408	Hel: DS wave fraction = 0.999081

• Zc3900: MC sample contains only SS component

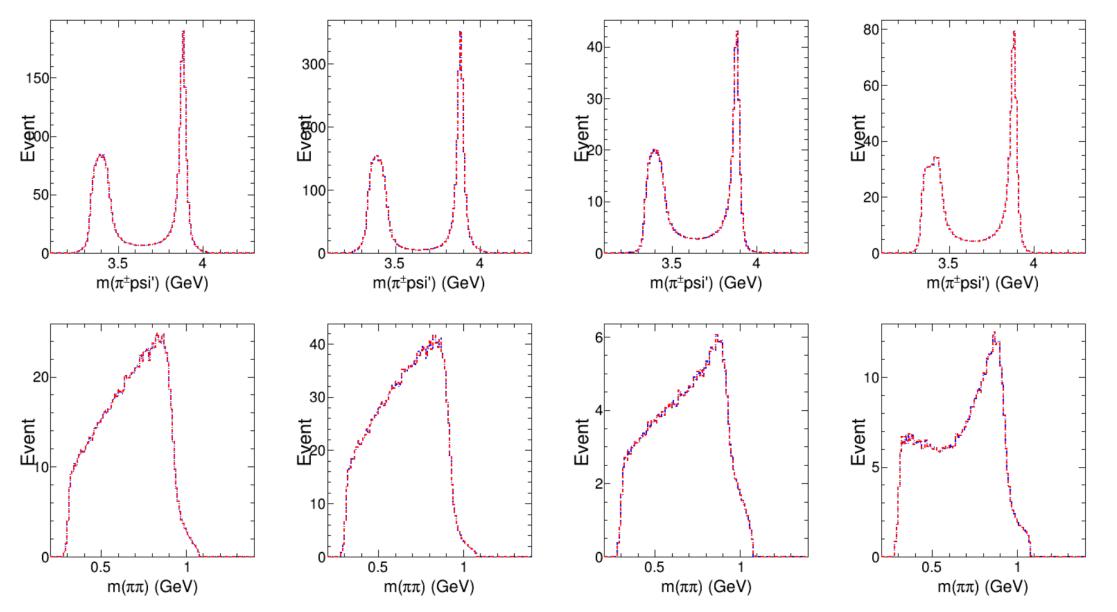
Without $J/\psi \rightarrow l^+l^-$ in covariant tensor formalism

Cov: SS wave fraction = 0.292	975 Hel: SS wave fraction = 0.905202
Cov: SD wave fraction = 0.500	984 Hel: SD wave fraction = 0.00800091
Cov: DS wave fraction = 0.074	4949 Hel: DS wave fraction = 0.0859795
Cov: DD wave fraction = 0.131	546 Hel: DD wave fraction = 0.000817607

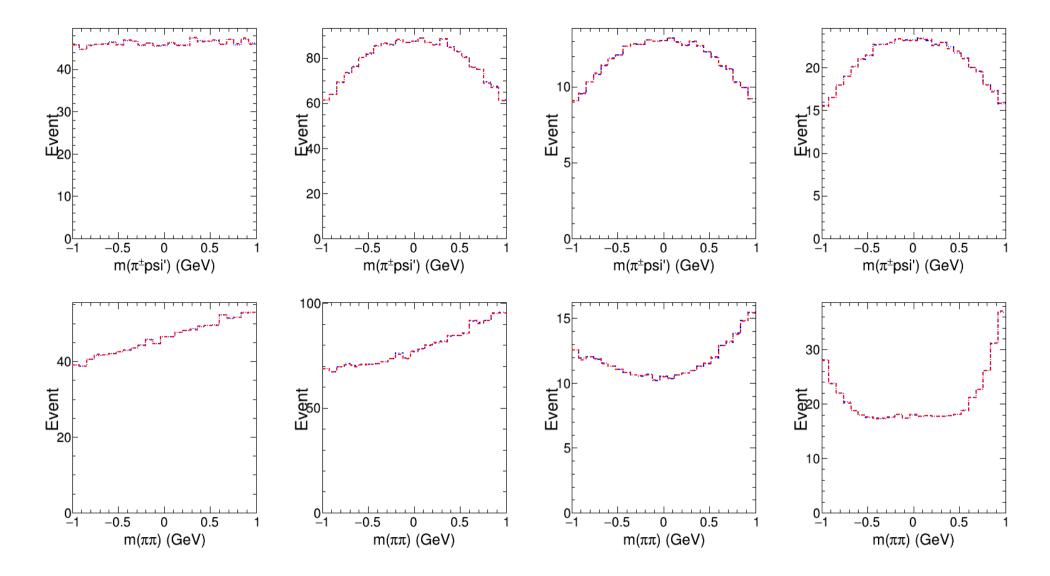
With $J/\psi \rightarrow l^+l^-$ in covariant tensor formalism

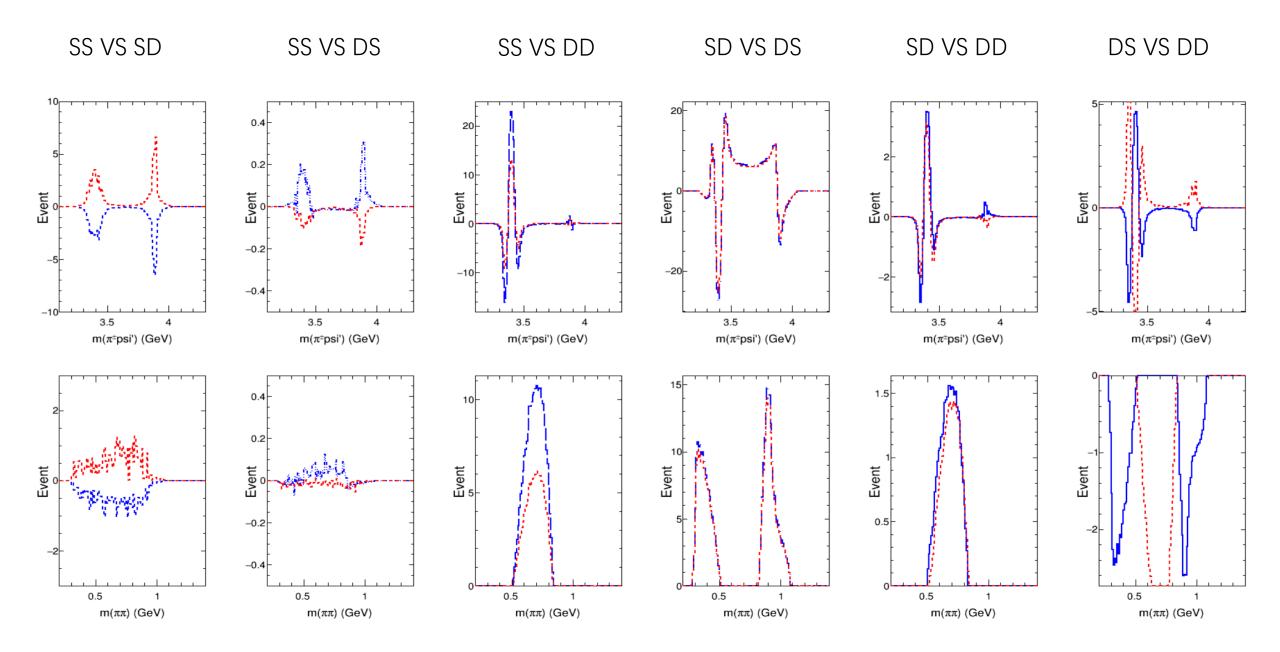
Cov: SS wave fraction = 0.930268	Hel: SS wave fraction = 0.905202
Cov: SD wave fraction = 0.00704404	Hel: SD wave fraction = 0.00800091
Cov: DS wave fraction = 0.062201	Hel: DS wave fraction = 0.0859795
Cov: DD wave fraction = 0.000487272	Hel: DD wave fraction = 0.000817607

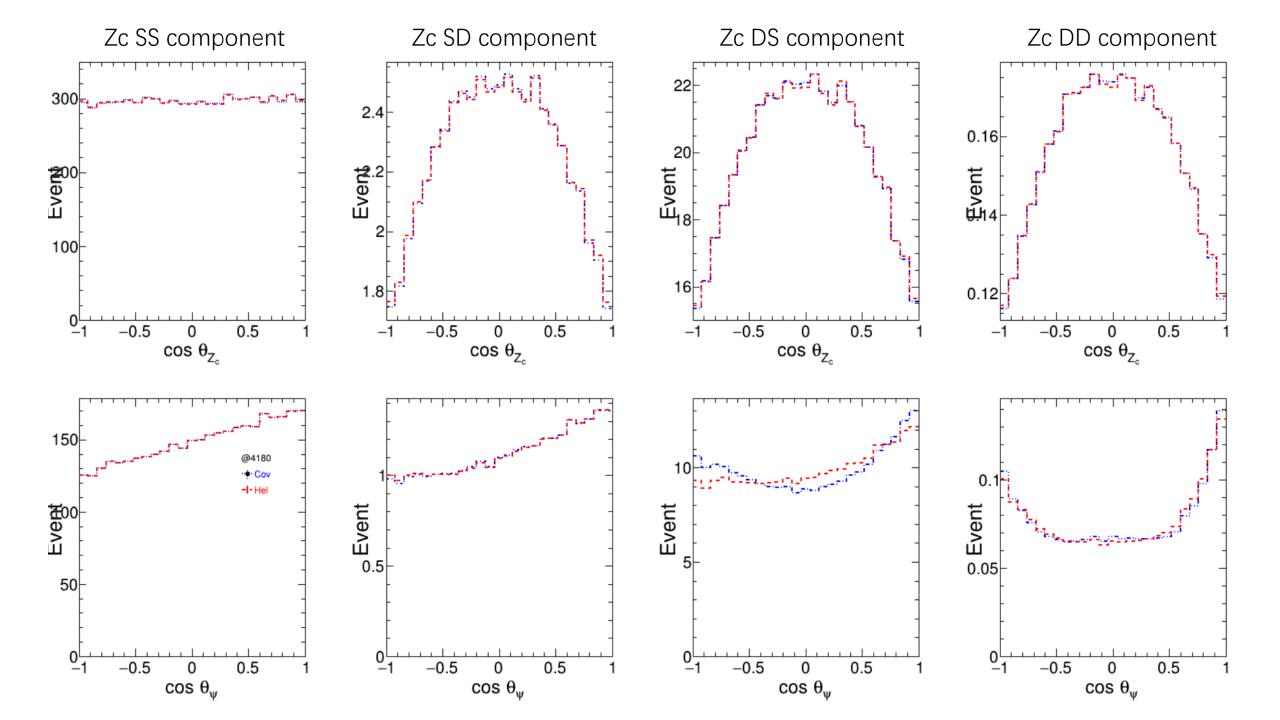
The comparison between cases with or without $J/\psi \rightarrow l^+l^-$ in covariant tensor formalism Red: with $J/\psi \rightarrow l^+l^-$ Blue: without $J/\psi \rightarrow l^+l^-$

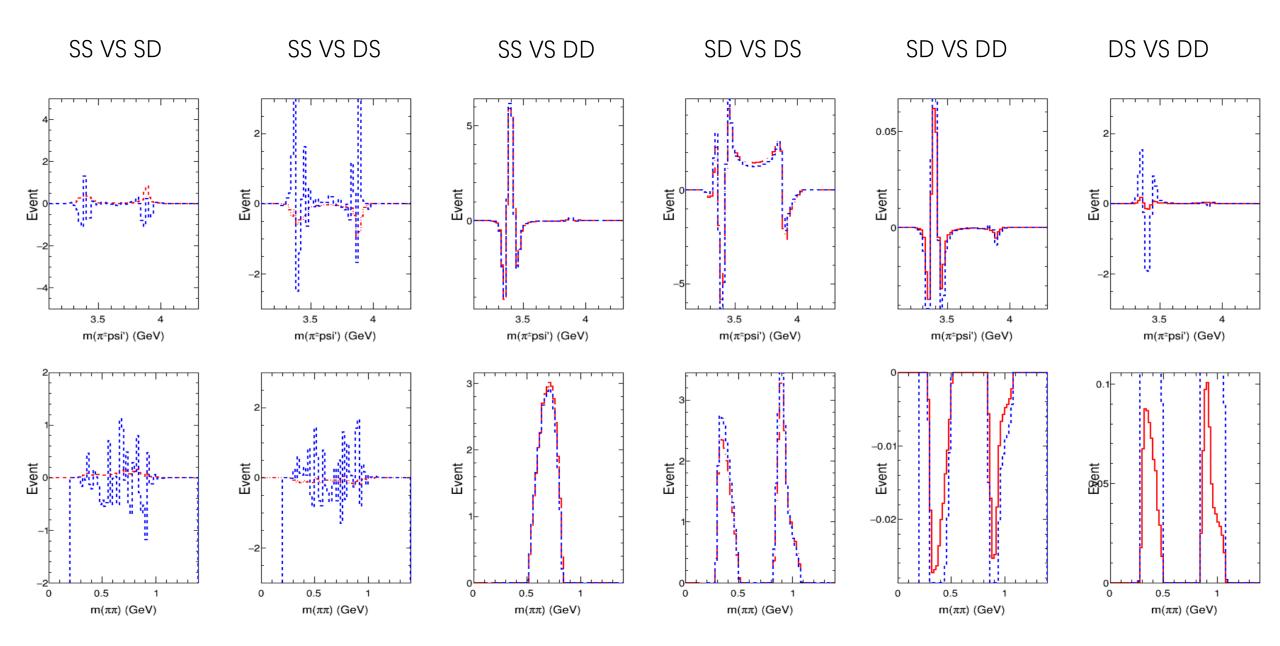


The comparison between cases with or without $J/\psi \rightarrow l^+l^-$ in covariant tensor formalism





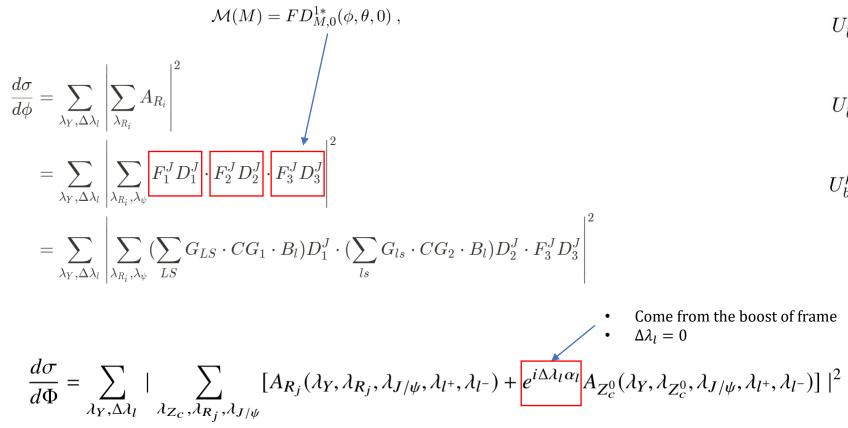




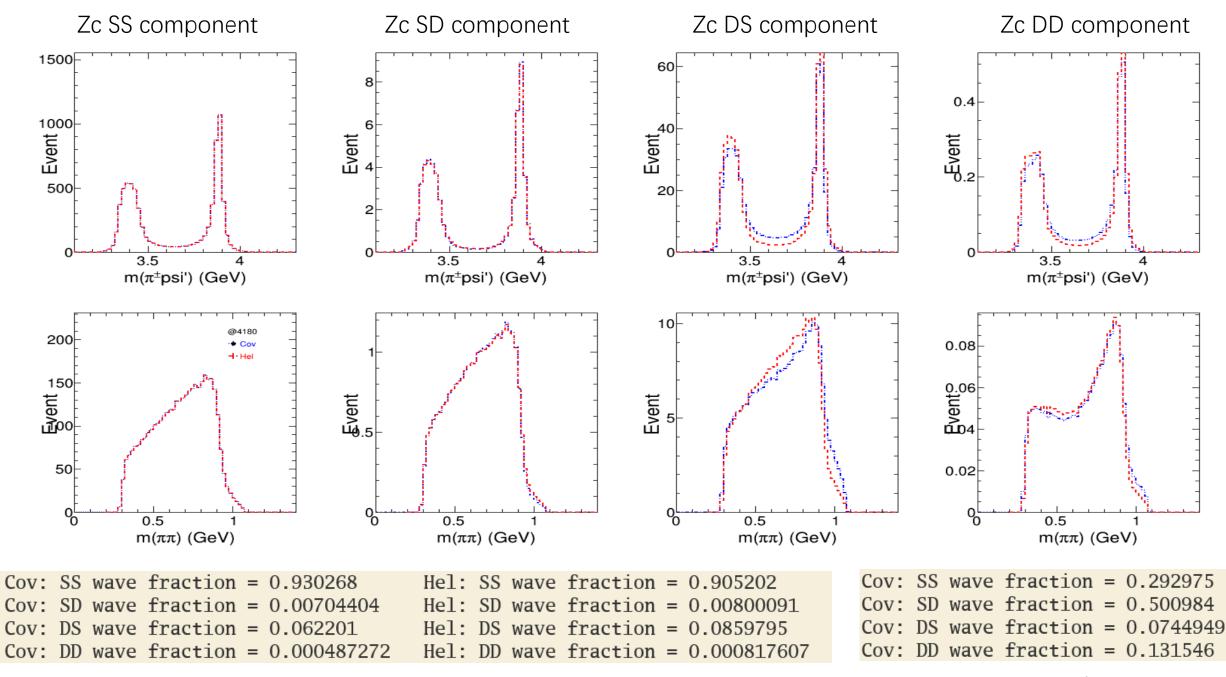
- Sequential decays: $Y \to \pi^+\pi^- J/\psi$, $J/\psi \to l^+l^-$
 - Assume leptons have 0 spin

$$\psi \to \pi^+ \pi^-$$
, $K^+ K^-$

这个过程是一个1⁻ \rightarrow 0⁻0⁻的过程,因为 $\eta_a\eta_b\eta_c = -1$ 和角动量守恒,因此只有1个 分波: P-wave,所以独立的Helicity振幅也只有一个。即 $F = F_{00}^1$,不变振幅:



```
In covariant tensor formalism,
just like:
        \psi \rightarrow \phi \pi^+ \pi^- \rightarrow \mathsf{K}^+ \mathsf{K}^- \pi^+ \pi^-
       U^{\mu}_{b_1SS} = \tilde{g}^{\mu\nu}_{(123)} \tilde{t}^{(1)}_{(12)\nu} f^{(\phi)}_{(12)} f^{(b_1)}_{(123)}
                                     +\tilde{g}^{\mu\nu}_{(124)}\tilde{t}^{(1)}_{(12)\nu}f^{(\phi)}_{(12)}f^{(b_1)}_{(124)},
      U^{\mu}_{b_1SD} = \tilde{t}^{(2)\mu\nu}_{(\phi3)} \tilde{t}^{(1)}_{(12)\nu} f^{(\phi)}_{(12)} f^{(b_1)}_{(123)}
                                     +\tilde{t}^{(2)\mu\nu}_{(\phi4)}\tilde{t}^{(1)}_{(12)\nu}f^{(\phi)}_{(12)}f^{(b_1)}_{(124)},
       U^{\mu}_{b_1DS} = \tilde{T}^{(2)\mu\lambda}_{(b_14)} \tilde{g}_{(123)\lambda\nu} \tilde{t}^{(1)\nu}_{(12)} f^{(\phi)}_{(12)} f^{(b_1)}_{(123)}
                                      +\tilde{T}^{(2)\mu\lambda}_{(b_{1}3)}\tilde{g}_{(124)\lambda\nu}\tilde{t}^{(1)\nu}_{(12)}f^{(\phi)}_{(12)}f^{(b_{1})}_{(124)},
     U^{\mu}_{b_1DD} = \tilde{T}^{(2)\mu\lambda}_{(b_14)} \tilde{t}^{(2)}_{(\phi3)\lambda\nu} \tilde{t}^{(1)\nu}_{(12)} f^{(\phi)}_{(12)} f^{(b_1)}_{(123)}
                                      +\tilde{T}^{(2)\mu\lambda}_{(b_13)}\tilde{t}^{(2)}_{(\phi4)\lambda\nu}\tilde{t}^{(1)\nu}_{(12)}f^{(\phi)}_{(12)}f^{(b_1)}_{(124)}
```



Without $J/\psi \rightarrow l^+l^-$

