

Measurements of light hypernuclei properties and production yields in Au+Au collisions from the STAR experiment

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Outline

Introduction

- Hypernuclei measurements in STAR BES-II
 - Internal structure
 - Branching ratios, lifetimes
 - Production mechanism

Yields, particle ratios, directed flow

Summary and outlook





Introduction: what and why

• What are hypernuclei?



- Why hypernuclei?
 - Probe hyperon-nucleon (Y-N) interaction lacksquare
 - Strangeness in high density nuclear matter
 - Equation-of-State (EoS) of neutron star





Marian Danysz (right) and Jerzy Pniewski (left) discovered hypernuclei in 1952



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 $^{4}_{\Lambda}$ He

 $^{4}_{\Lambda}$ He



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Introduction: how

- Experimentally, we can make measurements related to:
 - 1. Internal structure
 - Lifetime, binding energy, branching ratios etc. \bullet

Understanding hypernuclei structure can provide insights to the Y-N interaction

- 2. Production mechanism
 - Spectra, collectivity etc.

The process of hypernuclei formation in violent heavy-ion collisions is not well understood







Introduction: RHIC BES program

extends the energy reach below $\sqrt{s_{NN}}$ = 7.7 GeV, down to 3.0 GeV





• During the BES-II program, STAR utilized the fixed-target (FXT) setup, which



Introduction: hypernuclei and STAR BE

 Hypernuclei measurements are scarce in heavy-ion collision experiments



A. Andronic et al. PLB (2011) 697:203–207

- density

 \rightarrow A great opportunity to study hypernuclei production



• At low beam energies, hypernuclei production is expected to be enhanced due to high baryon

List of BES-II datasets:

Datasets with large statistics taken during **BES-II**

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Year	$\sqrt{s_{NN}}$ [GeV]	Events
2018	27	555 M
	<u>3.0</u>	258 M
	<u>7.2</u>	155 M
2019	19.6	478 M
	14.6	324 M
	<u>3.9</u>	53 M
	<u>3.2</u>	201 M
	<u>7.7</u>	51 M
2020	11.5	235 M
	<u>7.7</u>	113 M
	<u>4.5</u>	108 M
	<u>6.2</u>	118 M
	<u>5.2</u>	103 M
	<u>3.9</u>	117 M
	<u>3.5</u>	116 M
	9.2	162 M
	<u>7.2</u>	317 M
2021	7.7	101 M
	<u>3.0</u>	2103 M
	<u>9.2</u>	54 M
	<u>11.5</u>	52 M
	<u>13.7</u>	51 M
	17.3	256 M
	<u>7.2</u>	89 M

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fication and hypernuclei reconstruction







- Particle identification from energy loss measurement using TPC
- KF particle package^[1] is used for signal reconstruction
- Hypernuclei reconstructed via their weak decay channels: $^{3}_{\Lambda}H \rightarrow ^{3}He + \pi^{-}$ $^{3}_{\Lambda}H \rightarrow d + p + \pi^{-}$ $^{4}_{\Lambda}H \rightarrow ^{4}He +$

2023/4/24





$$\pi^{-}$$
 $^{4}_{\Lambda}\text{He} \rightarrow ^{3}\text{He} + p + \pi^{-}$

[1]Zyzak M, Kisel I, Senger P. Online selection of short-lived particles on many-core computer architectures in the CBM experiment at FAIR[R]. Collaboration FAIR: CBM, 2016.

Vertex





Hypernuclei signal reconstruction









Relative branching ratio: $R_3 =$



- Improved precision on R₃ lacksquare

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- F. Hildenbrand et al. PRC 102, 064002 (2020)
- Recent calculation shows that R_3 may be sensitive to the binding energy (B_{Λ}) of $^{3}_{\Lambda}H$
 - $B_{\Lambda} \rightarrow$ provide constraints to Y-N interaction
 - Using $\sqrt{s_{NN}} = 3.0$ GeV data:
 - $R_3 = 0.272 \pm 0.030(stat.) \pm 0.042(syst.)$
 - Model comparison suggesting a weakly-bounded state for ${}^{3}_{\Lambda}H$

• Stronger constraints on absolute B.R.s and $^{3}_{\Lambda}H$ internal structure models









${}_{\Lambda}^{3}$ H, ${}_{\Lambda}^{4}$ H and ${}_{\Lambda}^{4}$ He lifetimes



³_AH: ALICE(2022),arXiv:2209.07360 ⁴_AH: JPARC(2023),arXiv:2302.07443

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$^{\mathbf{4}}_{\Lambda}\mathbf{H}$

⁴_{\lambda}He

Using $\sqrt{s_{NN}}$ = 3.0 GeV and 7.2 GeV datasets:

- $^{3}_{\Lambda}$ H: $\tau = 221 \pm 15$ (stat.) ± 19 (syst.)[ps]
- $^{4}_{\Lambda}$ H: $\tau = 218 \pm 6$ (stat.) ± 13 (syst.)[ps]

⁴_AHe: $\tau = 229 \pm 23(\text{stat.}) \pm 20(\text{syst.})[\text{ps}]$

- Indication of shorter lifetimes for ${}^3_{\Lambda}H$, ${}^4_{\Lambda}H$ and ${}^4_{\Lambda}He$ than that of free Λ (with 1.8 σ , 3.0 σ , 1.1 σ respectively)
- Consistent with former measurements and world average values
- $\tau_{_{\Lambda}H}$: consistent with calculation including pion FSI^[1] and calculation with Λd 2-body picture^[2] within 1 σ
- au_{4H}^{4} and au_{4He}^{4} : consistent with expectations from isospin rule

Precision ${}^{3}_{\Lambda}$ H and ${}^{4}_{\Lambda}$ H measurements provide tighter constraints on models.

[1]A. Gal and H. Garcilazo, PLB 791, 48 (2019) [2]J.G. Congleton, J. Phys. G 18, 339 (1992)



Hypernuclei acoduction at 3 GeV



but fails to reproduce the trend of $^{3}_{\Lambda}$ H in 10-50%

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• Transport model (JAM) with coalescence approximately reproduces trends of $^{4}_{\Lambda}$ H rapidity distributions seen in data,









$^{3}_{\text{A}}\text{H}$ and $^{4}_{\text{A}}\text{H}$ directed flow at 3 GeV



- First observation of ${}^{3}_{\Lambda}$ H and ${}^{4}_{\Lambda}$ H directed flow (v₁) in mid-central 5-40% Au+Au collisions at 3 GeV
- Mid-rapidity v_1 slopes of ${}^3_{\Lambda}H$ and ${}^4_{\Lambda}H$ follow baryon mass scaling.

 \rightarrow Imply coalescence process to be the dominant formation mechanism for $^{3}_{\Lambda}H$ and $^{4}_{\Lambda}H$ production in the 3 GeV heavy-ion collisions

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arXiv:2211.16981 accepted by PRL

Xiujun Li, ATHIC 2023



Energy dependence of hypernuclei production in heavy-ion collisions



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- - Observed at both 0-10% and 10-40% centrality in Au+Au collisions at 3 GeV.
- The ${}^{4}_{\Lambda}H/{}^{4}He$ yield ratios are comparable to that of Λ/p
- Thermal model calculations including excited $^4_\Lambda H^*$ feed-down show a similar trend
 - Feed-down from excited state enhances ${}^{4}_{\Lambda}H$ production
 - Support creation of excited A=4 hypernuclei in heavy-ion collisions

A. Andronic et al, PLB 697 (2011) 203 (Thermal model)





at 3 GeV

Strangeness population factor $S_{\rm A}$

Relative suppression of hypernuclei production compared to light nuclei production

$$S_{A} = \frac{{}^{A}_{\Lambda}H}{{}^{A}_{He} \times \frac{\Lambda}{p}} = \frac{B_{A}({}^{A}_{\Lambda}H)(p_{T})}{B_{A}({}^{A}_{He})(p_{T})} o^{\triangleleft}$$

S.Zhang, PLB 684(2010)224

- B_A: Coalescence parameters

Expect ~1 if no suppression S₃< 1: relative suppression of ${}^{3}_{\Lambda}H/{}^{3}He$ compared to Λ/p $S_4 \sim 1$, $S_4 > S_3$: ${}^4_{\Lambda}H/{}^4He$ is comparable to Λ/p , possibly due to feed-down contributions from exited state that enhances ${}^4_{\Lambda}H$

No obvious kinematic and centrality dependence of $S_{3,4}$ is observed at 3 GeV.

 \rightarrow Coalescence parameter B_A of A_AH and AHe follows similar tendency versus p_T , rapidity and centrality,

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 10^{-1}





- indicates that N-N and Y-N interactions that drive coalescence dynamics in these collisions are similar





Energy dependence of S₃



- STAR, Science 328 (2010) 58 ALICE, PLB 754 (2016) 360 E864, PRC 70 (2004) 024902 NA49, J.Phys.Conf.Ser.110(2008)032010
- A. Andronic et al, PLB 697 (2011) 203 (Thermal (GSI))
- S. Zhang, PLB 684(2010)224 (Coal.+AMPT)
- T. Reichert, J. Steinheimer et al, arXiv:2210.11876(2022) (UrQMD, Thermal-FIST)

- Data show a hint of an increasing trend from $\sqrt{s_{NN}}$ = 3.0 GeV to 2.76 TeV
- For coalescence models, the energy dependence is sensitive to the source radius (Δr)
- Thermal-FIST, which includes feed-down to p and ${}^{3}\text{He}$ from unstable nuclei, describes the S_{3} data reasonably well







Summary



Presented measurements on hypernuclei production in the high-baryon-density region with

Xiujun Li, ATHIC 2023

Outlook

• e.g.
$${}^{4}_{\Lambda\Lambda}\text{He} \rightarrow {}^{4}_{\Lambda}\text{He}\pi, {}^{5}_{\Lambda\Lambda}\text{He} \rightarrow {}^{5}_{\Lambda}\text{He}\pi$$

14.6

Coalesc. (JAM)

30

- PHQMD

Outlook

• e.g.
$${}^{4}_{\Lambda\Lambda}\text{He} \rightarrow {}^{4}_{\Lambda}\text{He}\pi, {}^{5}_{\Lambda\Lambda}\text{He} \rightarrow {}^{5}_{\Lambda}\text{He}\pi$$

