

Heavy Flavor Physics Simulation For Hadronization at EicC

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Outline

- Introduction
- Simulation Setup
- Reconstruction
- Physics Projections
- Conclusions

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Introduction





- Why Eic?
 - Electro-magnetic probe for nuclear structure
 - Cleaner DIS background than HI collision
- Why EicC?
 - High luminosity and unique kinematic coverage
- Why heavy flavor?
 - $\gamma g \rightarrow c \bar{c}$: Sensitive to gluon distribution

Introduction





- Nuclear Physics A 740 (2004) 211–245 1.0 0.9 0.9 0.8 0.7 0.6 7 9 11 13 15 17 19 21 23 \mathcal{V} (GeV)
- Both hadronization in vacuum and nucleus medium is complex and poorly understood.
 - In the recent ALICE results, the baryon-to-meson ratio deviated from the results of pp, e⁺e⁻ and ep in the past.
 - Both parton energy loss and the hadron absorption model, two theories with different scales, can well describe the suppression of light hadrons.
- Our poor understanding of hadronization can be improved with precise measurement of heavy flavor production at Electron-ion collider in China (EicC).

Introduction

Particle	$\begin{array}{c} {\rm Momentum} \\ {\rm (GeV/c/u)} \end{array}$	${ m CM~energy}\ { m (GeV/u)}$	Average polarization	Luminosity at the nucleon level $(cm^{-2} \cdot s^{-1})$	Integrated luminosity (fb^{-1})
е	3.5		80%		
р	20	16.76	70%	2.00×10^{33}	50.5
d	12.90	13.48	Yes	8.48×10^{32}	21.4
$^{3}\mathrm{He}^{++}$	17.21	15.55	Yes	6.29×10^{32}	15.9
$^{7}\mathrm{Li}^{3+}$	11.05	12.48	No	9.75×10^{32}	24.6
${}^{12}C^{6+}$	12.90	13.48	No	8.35×10^{32}	21.1
$\rm ^{40}Ca^{20+}$	12.90	13.48	No	8.35×10^{32}	21.1
$^{197}{\rm Au}^{79+}$	10.35	12.09	No	9.37×10^{32}	23.6
$^{208}{\rm Pb}^{82+}$	10.17	11.98	No	9.22×10^{32}	23.3
$^{238}\mathrm{U}^{92+}$	9.98	11.87	No	8.92×10^{32}	22.5

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Detectors



	p-going	Barrel	e-going	
Silicon	5	5	5	ITS3, pixel size 10 μn
MPGD	1	4	0	1.71

Pitch Size(μm)

10

10

10

10

10

150(rp)x150(z)

150(rp)x150(z)

150(rp)x150(z)

150(rp)x150(z)

Barrel:

R(cm)

3.30

4.35

5.40

8.00

15.00

47.72

49.57

75.61

77.46

Length(cm)

28

28

28

28

38.70

127.47

127.47

201.98

201.98

End cap e going:

In R(cm)	Out R(cm)	Z(cm)
3.18	18.62	-25
3.18	36.50	-49
3.18	54.66	-73
3.95	77.46	-109.0
5.26	77.46	-145.0

End cap p going:

In R(cm)	Out R(cm)	Z(cm)
3.18	18.62	25
3.18	36.50	49
3.47	54.66	73
5.08	77.46	103.65
6.58	77.46	134.33
8.16	150.00	165.00

https://gitee.com/aiqiang-guo/EicC_Mvd_DP



Detector Parameters



 $p_{hadron} > 0.3 GeV/c_{10}$

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Generator and Reconstruction



- Event Generator
 - e(3.5 GeV)p(20 GeV)
 - $1 \, GeV^2 < Q^2 < 200 \, GeV^2$
- Pythiae6.428 $L_{int} = 8.06 \, f b^{-1}$

EicC run for 16 month

D^0/D^0 Reconstruction



0.5

 $\cos \theta_{xy}$

$D^0/\overline{D^0}$ Reconstruction





- no PID: all combination between opposite charge particles
- with PID: perfect PID at acceptance and abandoned if out of acceptance
- Vertex: topological cut applied to suppress background
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$\Lambda_c^{+/-}$ Reconstruction



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Baryon-to-meson Ratio vs p_T



- Fragmentation functions are thought to be universal.
- An enhancement of Λ_c^+/D^0 in pp system are found by ALICE in p+p at $\sqrt{s} = 5 TeV$ and $\sqrt{s} = 7 TeV$.
- Fragmentation can be well studied at wide kinematics region, especially at mid-rapidity, because of the high luminosity.



- Track multiplicity are in truth level(efficiency~1 applied).
- Particles from Λ_c^+ and D^0 candidates are removed from multiplicity counting.
- EicC can provide measurement at low multiplicity, a complementary to EIC-US high multiplicity measurements.
- Fragmentation can be well studied, especially at mid-rapidity, because of the high luminosity.

Double Ratio of D^0



- Double Ratio $R_A^h = \frac{N_{eA}^h(z, x, Q^2) / N_{eA}^{e_{DIS}}(x, Q^2)}{N_{ep}^h(z, x, Q^2) / N_{ep}^{e_{DIS}}(x, Q^2)}$
- Number of DIS electron divided to cancel the initial state effects
- Measurements at low ν and Q^2 regions
- Study nuclear matter effects at wider kinematic region.

z: hadron energy fraction of virtual photon at the target rest frame

New Configuration From Yuming

Det_v3_10_10_MMB	Det_v3_10_10_LMB
Det_v3_10_25_MMB	Det_v3_10_25_LMB
Det_v3_15_10_MMB	Det_v3_15_10_LMB
Det_v3_15_25_MMB	Det_v3_15_25_LMB
Det_v3_20_10_MMB	Det_v3_20_10_LMB
Det_v3_20_25_MMB	Det_v3_20_25_LMB
	Det_v3_10_10_MMB Det_v3_10_25_MMB Det_v3_15_10_MMB Det_v3_15_25_MMB Det_v3_20_10_MMB Det_v3_20_25_MMB

Pitch Size In Green Box: To be studied Low/Middle/High Materiel Budget



Without Any Cut

- Size2 is more important to the shape than Size1.
- The material budget will compromise the momentum resolution.
- When the material budget is low, the resolution of signal is almost equal to $48.5 MeV/c^2$.
- Significance isn't sensitive to pitch size. Actually, it becomes a little smaller as the pitch size decreases.

Topological Cut Comparison(Signal)



• Pitch sizes and material budget seem to have little effect to the topological distribution.

TMVA

TMVA



 $s/\sqrt{s+b} = 90.5$

 $s/\sqrt{s+b} = 68$

TMVA can't improve the performance. It may be caused by poor training.

Conclusion

- Heavy flavor physics with the latest detector performance parameters is simulated for EicC.
- Several Projections are done including:
 - Baryon-to-meson Ratio: Hadronization in vacuum at different collision system
 - D^0 double ratio vs. z: a complementary to EIC-US, wider kinematic region to study the nuclear medium of nuclear effect on hadronization
- EicC can provide precise measurements on open heavy flavor production and deepen our understanding of hadronization both in e+p system and in nuclear medium.

Thank you for attention!

BACKUP

pythiaeRHIC tune

MSEL=2 MSTP(13)=1 MSTP(14)=30 MSTP(15)=0 MSTP(16)=1 MSTP(17)=6 ! MSTP 17=(MSTP(18)=3	PARP(13)=1 PARP(18)=0.17 ! hermes P, PARP(81)=1.9 PARP(89)=1800 PARP(90)=0.16 PARP(91)=0.40 ! rms kt D: PARP(93)=5.	MSTJ(12)=1 MSTJ(45)=5 MSTU(16)=2 MSTU(112)=5 MSTU(113)=5 MSTU(114)=5 ! Now all	CKIN(32)=-1. CKIN(35)=0. CKIN(36)=-1 CKIN(37)=0. CKIN(38)=-1. CKIN(39)=4.
MSTP(19)=1 ! Hermes M!F	PARP(97)=6.0 ! D=1.0, t	CKIN(1)=1.	CKIN(40)=-1.
MSTP(20)=4 ! Hermes M:F	PARP(99)=0.40	CKIN(2)=-1.	CKIN(65)=1.e-09
MSTP(32)=8 F	PARP(100)=5	CKIN(3)=0.	CKIN(66)=-1. ! Ma
MSTP(38)=4 F	PARP(102)=0.28	CKIN(4)=-1.	CKIN(67)=0
MSTP(51)=10042 ! if p(F	PARP(103)=1.0	CKIN(5)=1.00	CVIN(68) = 1
!MSTP(51)=7 F	PARP(104)=0.8	CKIN(6)=1.00	CKIN(00) = I
MSTP(52)=2 !> p(F	PARP(111)=2.0	CKIN(7)=-10.	
MSTP(53)=3 F	PARP(161)=3.00	CKIN(8)=10.	CKIN(78)=-1.
MSTP(54)=1 F	PARP(162)=24.6	CKIN(9)=-40.	2
MSTP(55)=5	PARP(163)=18.8	CKIN(10)=40.	
MSTP(56)=1 F	PARP(164)=11.5	CKIN(11)=-40.	
MSTP(57)=1 F	PARP(165)=0.47679	CKIN(12)=40.	
MSTP(58)=5 F	PARP(166)=0.67597 ! PARP:	CKIN(13)=-40.	
MSTP(59)=1	! PARP(166)=0.5	CKIN(14)=40.	
MSTP(60)=7	! Now come a	CKIN(15)=-40.	
MSTP(61)=2	PARJ(1)=0.100	CKIN(16)=40.	
MSTP(71)=1 F	PARJ(2)=0.300	CKIN(17)=-1.	
MSTP(81)=0 F	PARJ(3)=0.4	CKIN(18)=1.	
MSTP(82)=1 F	PARJ(11)=0.5	CKIN(19)=-1.	
MSTP(91)=1 F	PARJ(12)=0.6	CKIN(20)=1.	
MSTP(92)=3 ! heri F	PARJ(21)=0.40 ! fragpt w:	CKIN(21)=0.	
MSTP(93)=1 F	PARJ(32)=1.0	CKIN(22)=1.	
MSTP(94)=2 ! D=3 F	PARJ(33)=0.80	CKIN(23)=0.	
MSTP(101)=1 F	PARJ(41)=0.30	CKIN(24)=1.	
MSTP(102)=1 F	PARJ(42)=0.58	CKIN(25)=-1.	
MSTP(111)=1 F	PARJ(45)=0.5	CKIN(26)=1.	
MSTP(121)=0 F	PARJ(170)=0.20 !pt for r	CKIN(27)=-1.	
! Now all !	!	CKIN(28)=1.	
PARP(2)=5. !min CMS erM	MSTJ(1)=1	CKIN(31)=2.	
PARP(13)=1	MSTJ(12)=1	CKIN(32)=-1.	
PARP(18)=0.17 ! herme: M	MSTJ(45)=5	CKIN(35)=0.	

Color Reconnection

the decay $B^+ = u\overline{b} \rightarrow u\overline{c}W^+ \rightarrow (u\overline{c})(c\overline{s}) \rightarrow (u\overline{s})(c\overline{c}) \rightarrow K^+ J/\psi \rightarrow K^+ \mu^+ \mu^-$ [4], where we have used brackets in intermediate states to delineate separate color singlet identities.

s10052-015-3674-4 (springer.com)

The MPI-based scheme <u>PYTHIA 8 online manual</u>

In this scheme partons are classified by which MPI system they belong to. The colour flow of two such systems can be fused, and if so the partons of the lower-pT system are added to the strings defined by the higher-pT system in such a way as to give the smallest total string length. The bulk of these lower-pT partons are gluons, and this is what the scheme is optimized to handle.

The newer scheme QCD-CR

The newer CR scheme builds on the minimization of the string length as well as the colour rules from QCD. A main feature of the new model is the introduction of junction structures. These are possible outcomes of the reconnection in addition to the more common string-string reconnections. The model works by constructing all pair of dipoles that are allowed to reconnect by QCD colour rules and switching if the new pair has a lower string length. Junctions are also allowed to be directly produced from three, and in some special cases, four dipoles. This is done iteratively until no further allowed reconnection lowers the total string length.

Color Reconnection for Beam Remnants



Figure 10: Examples of the formation of composite objects in a baryon beam remnant: (a) diquark, (b) baryon and (c) meson.