

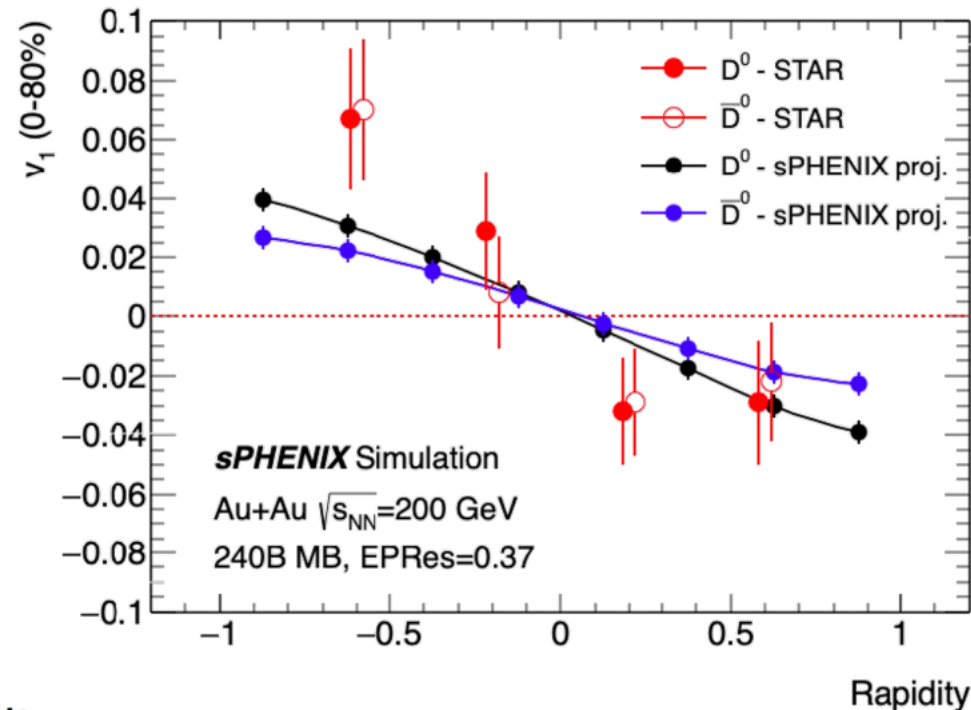
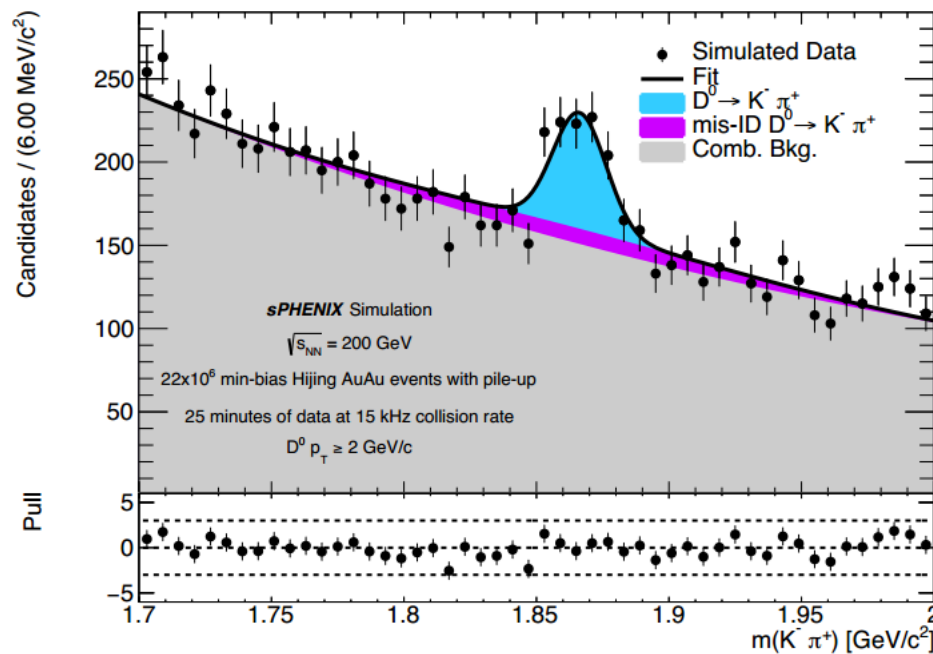
# Heavy flavor measurements at QM2023

HI-9.26-suyuan

# sPHENIX commissioning run started in May 18<sup>th</sup>

- First barrel HCal at RHIC for precise HF-jets
- Largest inclusive b-hadron sample at RHIC enabled by streaming readout

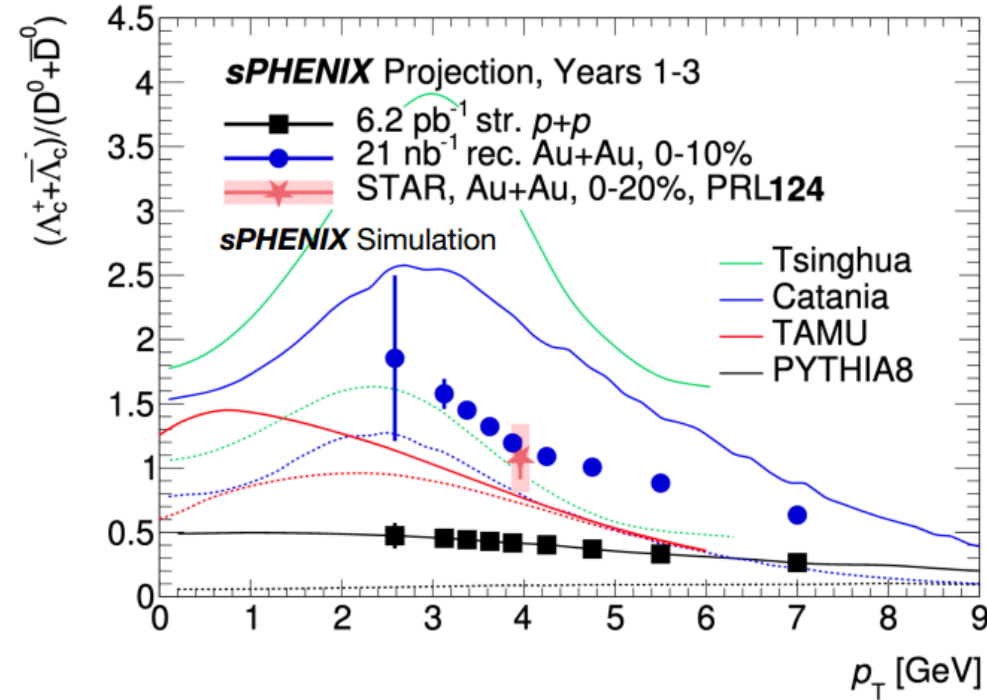
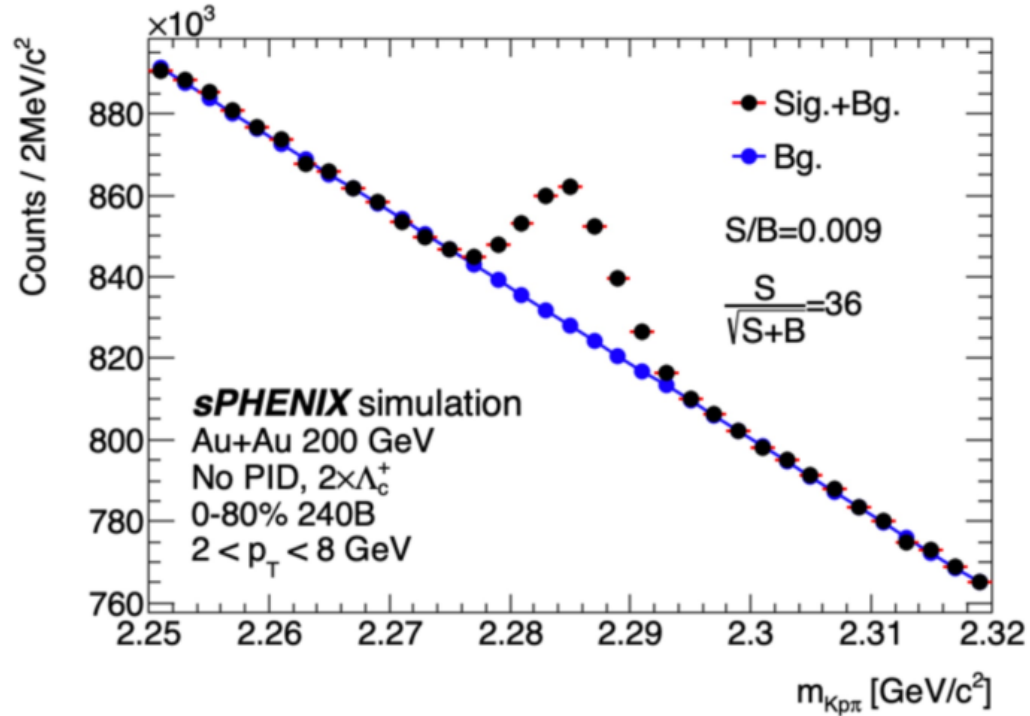
## Open charm physics simulation



- Prediction that transient mag. field alters  $v_1$
- This effect is odd under charge-conjugation, resulting in splitting
- $D^0 \rightarrow K^{\mp}\pi^{\pm}$  is mixed, requires good production knowledge

[PRL 123 \(2019\) 162301](#)  
[PRL 118 \(2017\) 212301](#)

# $\Lambda_c^+$ coalescence simulation

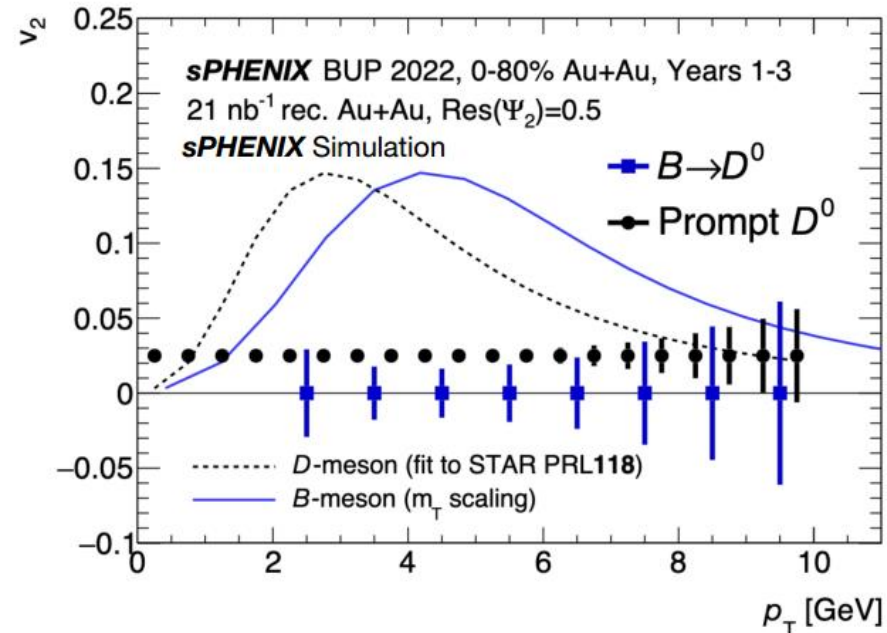
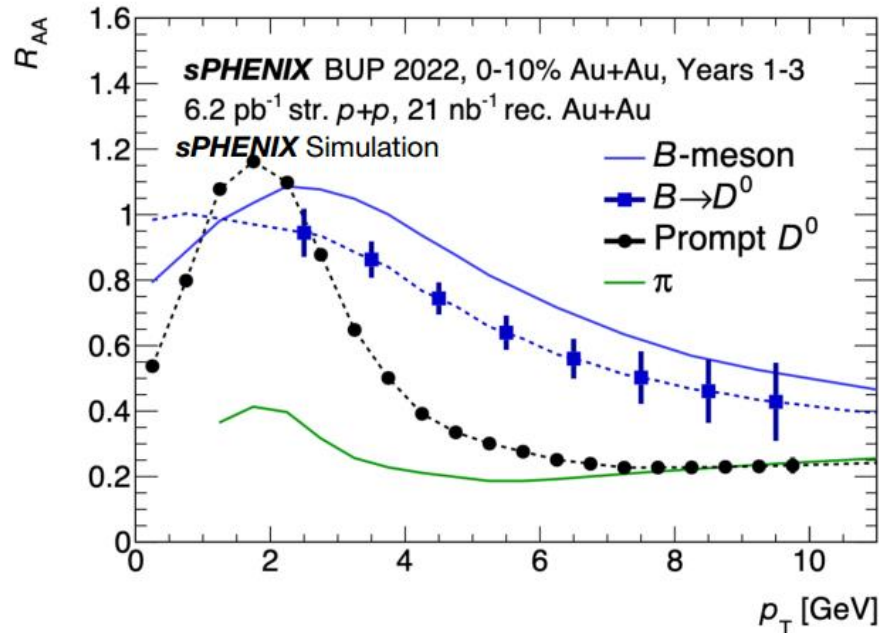


Open Charm Hadron	Constitution
$D^0$	$c\bar{u}$
$D^+$	$c\bar{d}$
$D_s^+$	$c\bar{s}$
$\Lambda_c^+$	$udc$

- Baryon/meson yields offer new insights in hadronization
- Models currently favor coalescence hadronization
- Low  $p_T$  region ( $< 8$  GeV) is key
- Huge benefit from streaming readout in  $pp$

[sPH-HF-2019-001](#)  
[PRL 124 \(2020\) 172301](#)

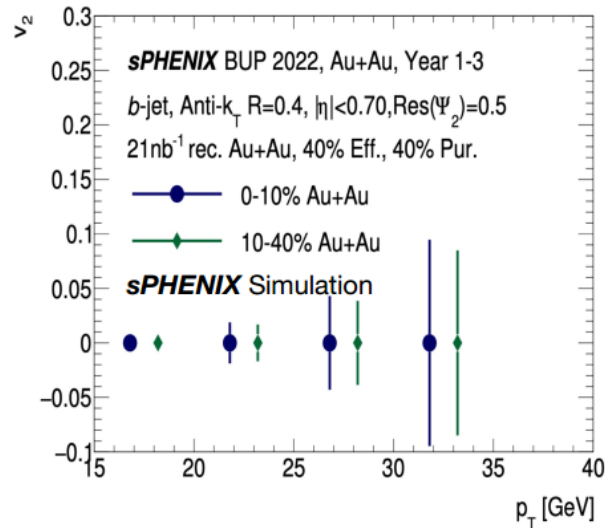
# b-hadron simulation



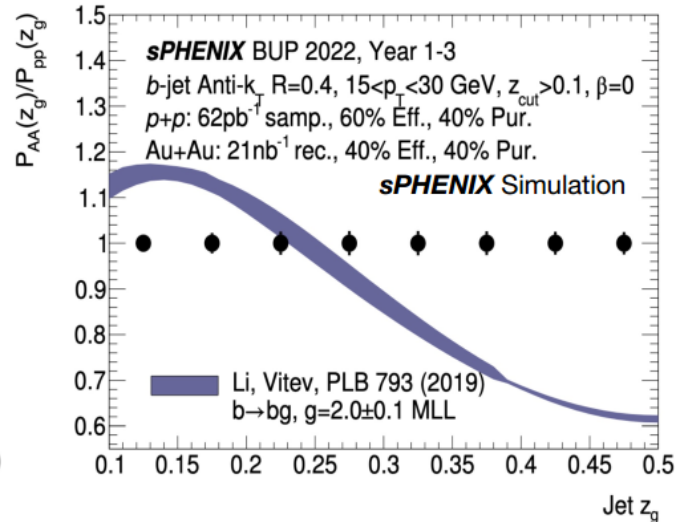
- Non-prompt *D*<sup>0</sup> reco. gives low *p*<sub>T</sub> b-hadron access
- sPHENIX aims to
  - constrain heavy quark diffusion coefficients
  - Parton energy loss mechanism

[PRL 118, 212301](#)

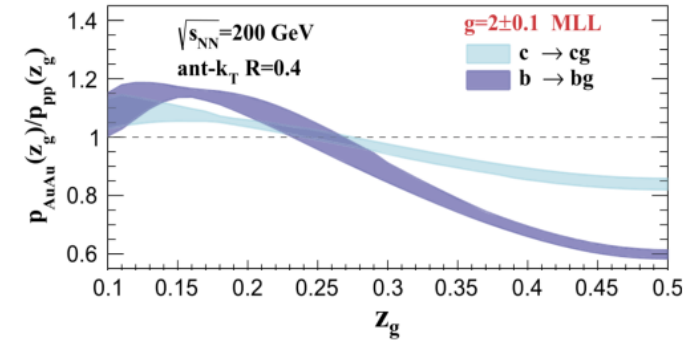
# Heavy flavor jet observable simulation



b-jet elliptical flow



Predicted b-jet subjet splitting sensitivity



Model comparison of b- and c-jet subjet splitting sensitivity

$$z_g = \frac{\min(p_{T1}, p_{T2})}{p_{T1} + p_{T2}}$$

$$p(z_g)|_j = \frac{1}{\sigma_j} \int dp_T d\eta \frac{d\sigma_j}{dp_T d\eta} \int_0^1 d\theta p(\theta, z_g)|_j$$

“We identify the region of jet transverse momenta where parton mass effects are leading and predict a unique reversal of the mass hierarchy of jet quenching effects in heavy ion relative to proton collisions. Namely, the momentum sharing distribution of prompt b-tagged jets is more strongly modified in comparison to the one for light jets.”  
- PLB 793 (2019)



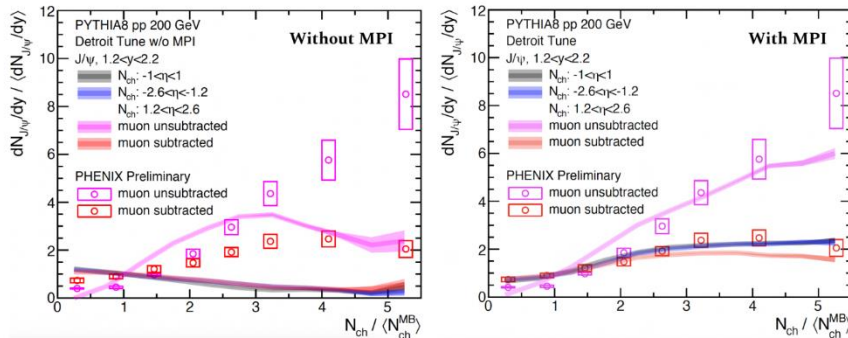
# Quarkonia & Heavy Flavor Overview

Three recent PHENIX analyses focus on the following collision systems and investigate the following:

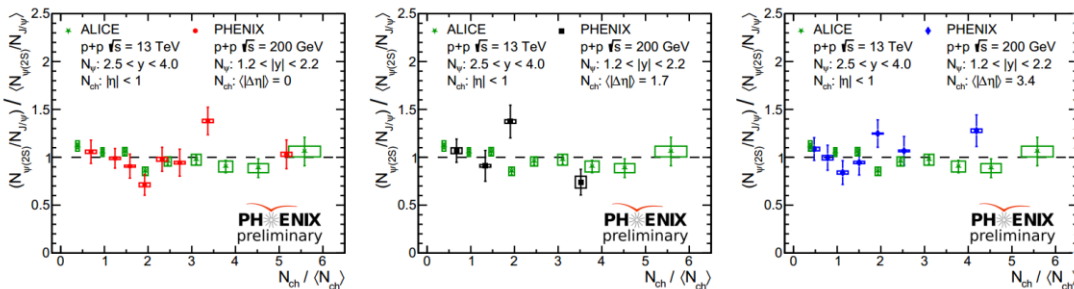
- 2013  $p+p$  at  $\sqrt{s} = 510$  GeV
- 2014  ${}^3\text{He}+\text{Au}$ ,  $\text{Au}+\text{Au}$  at  $\sqrt{s_{NN}} = 200$  GeV
- 2015  $p+p$ ,  $p+\text{Al}$ ,  $p+\text{Au}$  at  $\sqrt{s_{NN}} = 200$  GeV
- 2016  $d+\text{Au}$ ,  $\text{Au}+\text{Au}$  at  $\sqrt{s_{NN}} = 200$  GeV

- 1 Do we see evidence for multi-parton interactions at RHIC energies?
- 2 Are there final state effects on charmonium production in  $p+p$  collisions?
- 3 Is there evidence of mass ordering for charged hadron vs. open heavy flavor  $v_2$ ?

## Multiplicity Dependent $J/\psi$ Production



## Multiplicity Dependent $\psi(2S)$ to $J/\psi$



### SMALL SYSTEM COLLISIONS

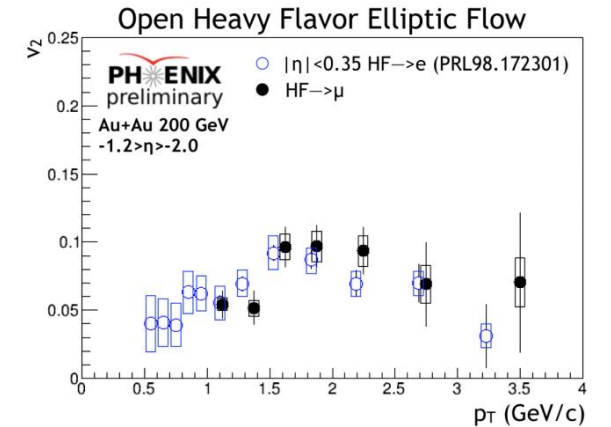
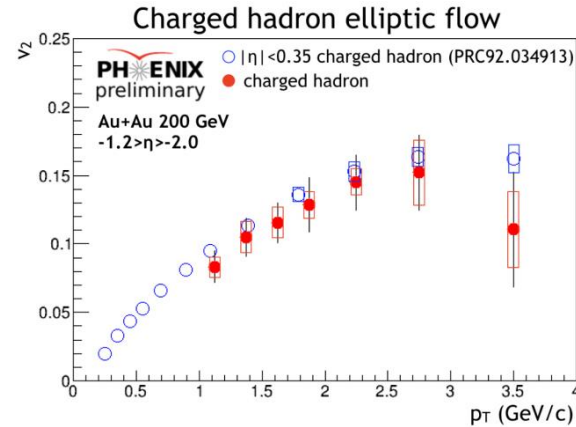
$p+p$

- Multiplicity dependent  $J/\psi$  varies based on  $\eta$  of charged particle tracks
    - PHENIX data well described by PYTHIA Detroit tune with MPI  $\Rightarrow$  Evidence for MPI at RHIC energies
  - $\psi(2S)$  to  $J/\psi$  ratio in  $p+p$  collisions shows weak dependence on multiplicity  $\Rightarrow$  No evidence for  $\psi(2S)$  final state effects in  $p+p$  collisions at RHIC
- Consistent with ALICE

### LARGE SYSTEM COLLISIONS

A+A

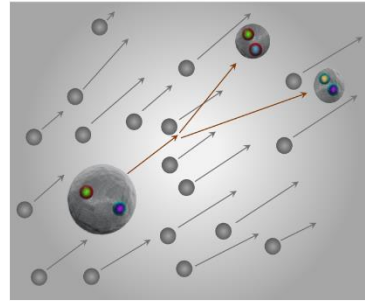
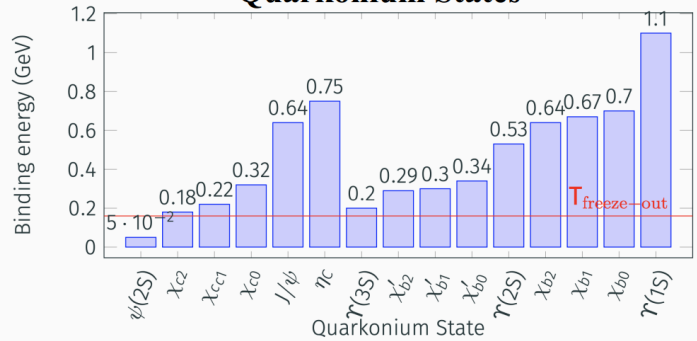
- First RHIC measurement of open heavy flavor  $v_2$  at forward rapidity
  - Results consistent with PHENIX mid-rapidity measurements, suggesting similar QGP effects (temperature/pressure gradients) at both rapidities
- $\Rightarrow$  Mass ordering observed at forward rapidity





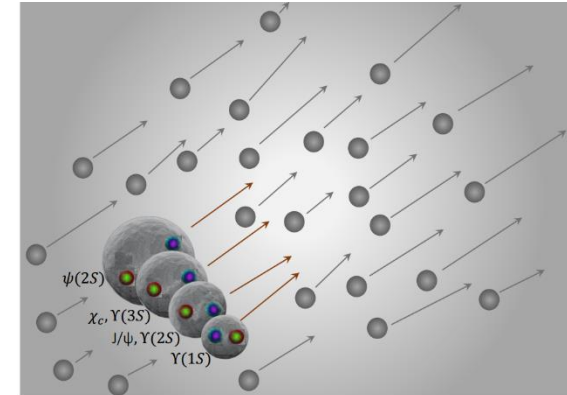
# New LHCb results on quarkonia production (and exotic hadron) in pp and pPb collisions.

## Quarkonium States



### Alternative way to break quarkonium states:

Large quarkonium states can break in high-multiplicity environment when interacting with co-moving particles [Ferrero, PLB749, 98 (2015)].



	$r \text{ (fm)}$
$J/\psi$	0.50
$\chi_c$	0.72
$\psi(2S)$	0.90
$\Upsilon(1S)$	0.28
$\chi_b$	0.44
$\Upsilon(2S)$	0.56
$\chi_b(2P)$	0.68
$\Upsilon(3S)$	0.78

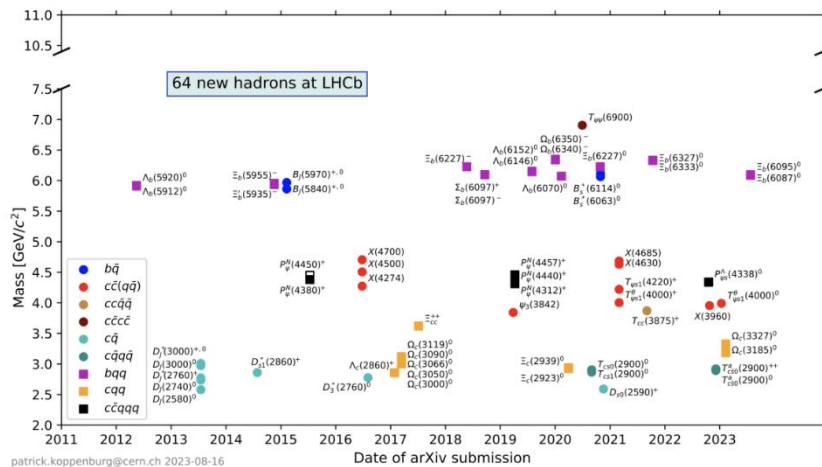
Non-Relativistic Potential Theory: Satz, J.Phys.G32:R25 (2006)

If medium is formed and temperature is above its binding energy, the quarkonium state is inhibited to be produced.

$T_{\text{freeze-out}} \sim 155\text{-}160 \text{ MeV}$

Lattice QCD A. Bazavov et al., PLB795 (2019) 15  
 Thermal Model N. Sharma et al. PRC 99 (2019) 044914  
 Thermal fits to ALICE data F. A. Flor, PLB 834 (2022) 137473

## Exotic Particles

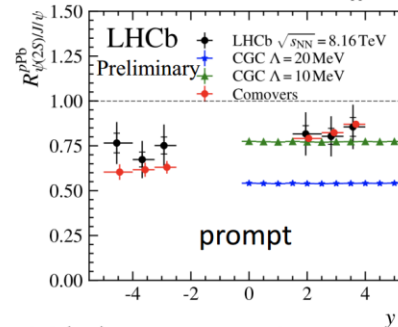


patrick.koppenburg@cern.ch 2023-08-16

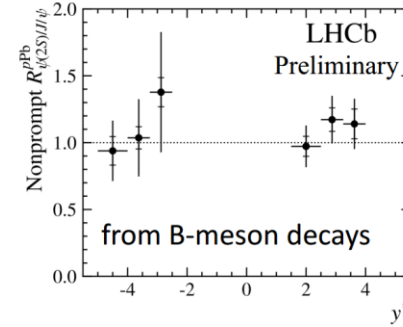
<https://www.nikhef.nl/~pkoppenb/particles.html>

## New $\psi(2S)$ Result at $\sqrt{s_{NN}}=8.16 \text{ TeV}$

$$R_{\psi(2S)/J/\psi}^{pPb} = \frac{R_{pPb}(\psi(2S))}{R_{pPb}(J/\psi)} = \frac{\frac{\sigma(\psi(2S))}{\sigma(J/\psi)}_{pPb}}{\frac{\sigma(\psi(2S))}{\sigma(J/\psi)}_{pp}}$$



LHCb-PAPER-2023-024 in preparation

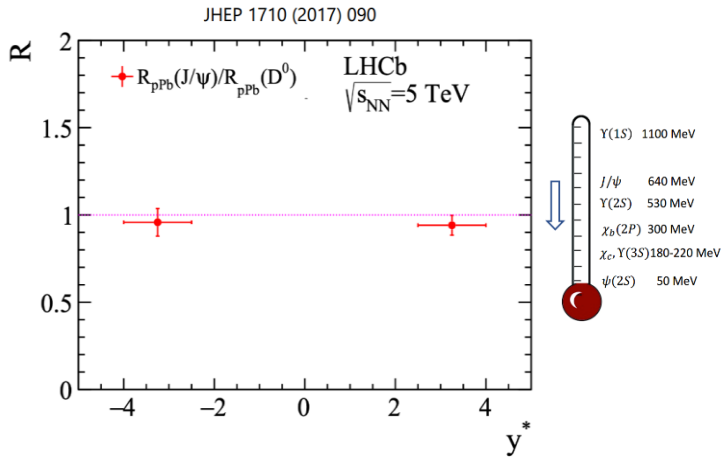


Additional  $\psi(2S)$  suppression only present in the prompt component, consistent with co-mover particle interactions [PLB749, 98 (2015)]

CGC : Factorization violating soft gluon exchanges PRC97, 014909 (2018)

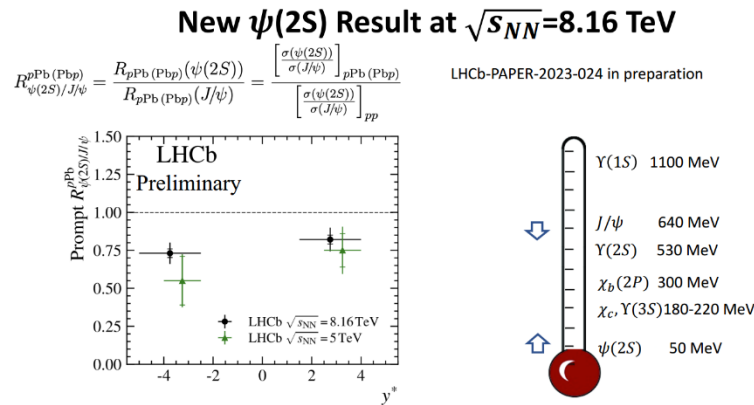
# “Now we know how to use the quarkonium thermometer”

Agnes Mocsy's thermometer



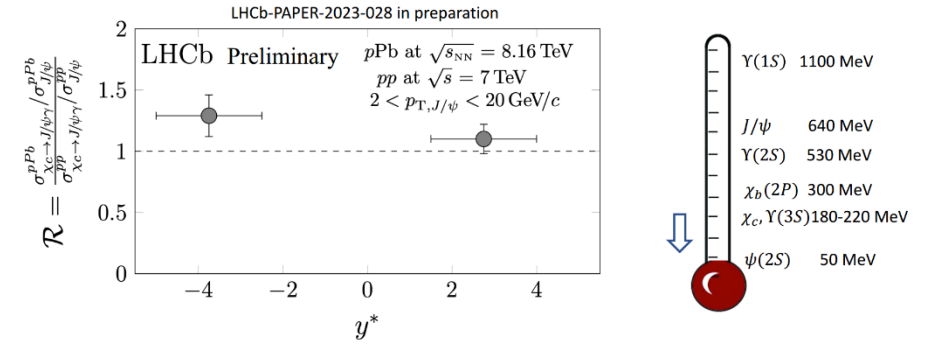
Initial-state effects are cancelled out in the  $R_{pA}(J/\psi)/R_{pA}(D^0)$  ratio.

$J/\psi$  yield is not affected by final-state effects.



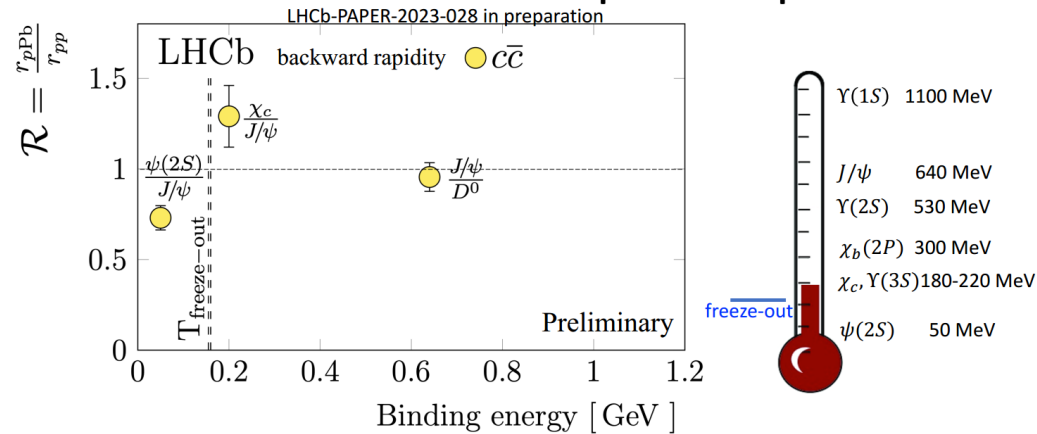
8.16 TeV result more precise and consistent with 5 TeV.

Confirming the existence of final-state effects on the  $\psi(2S)$  yields.



$\chi_c$  double ratio consistent with **NO final-state dissociation of  $\chi_c$  states in pPb collisions.**

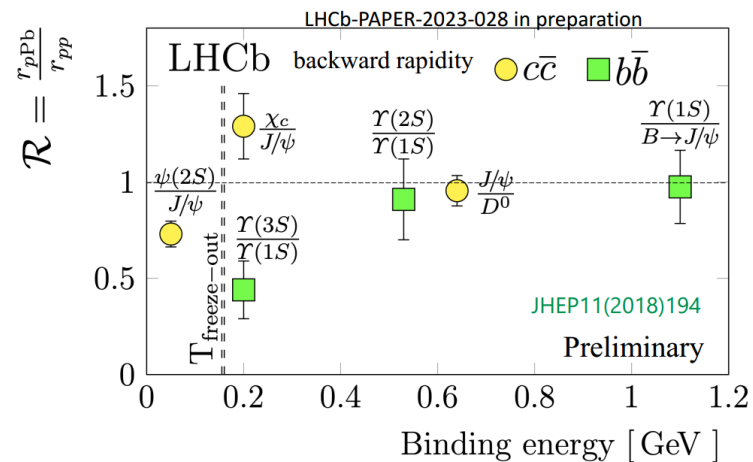
## Constraints to maximum medium temperature in pPb collisions.



The medium temperature hypothetically formed in pPb collisions cannot inhibit the formation of charmonium states with binding energy larger than 180 MeV, just 20 MeV above the estimated freeze-out temperature.

Limiting the medium temperature in this **small system close of a hadron environment**

18



Despite the similar binding energy and size with  $\chi_c$ ,  $Y(3S)$  is dissociated.

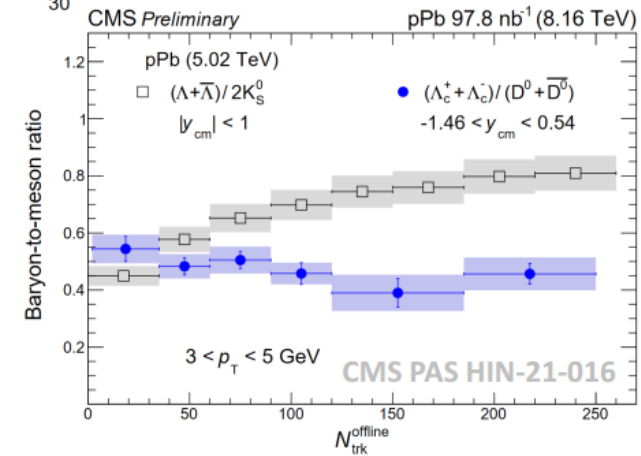
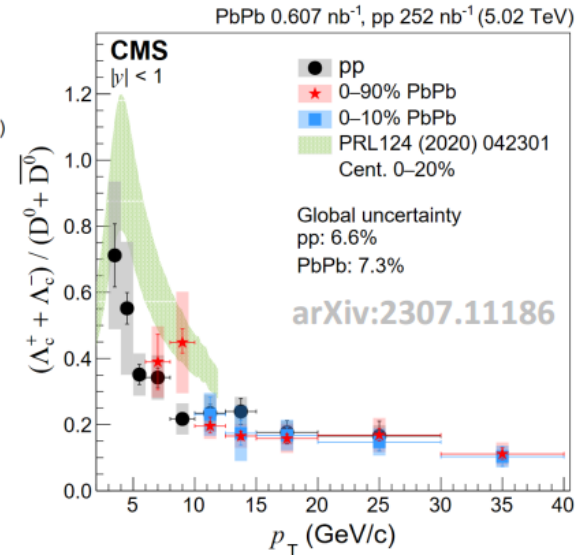
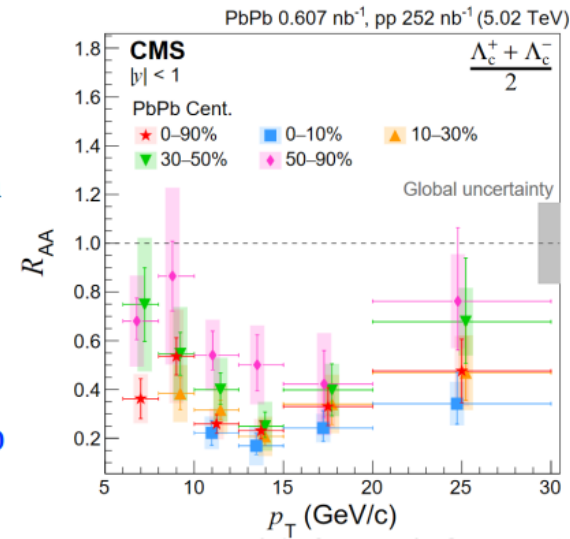
$Y(3S)$  is 2.9 x heavier and slower than  $\chi_c$ , more likely to interact with comoving particles. 8  
Theoretical input is welcome !!!

Cesar Luiz da Silva  
September 5th, 2023





- ❑  $\Lambda_c^+$  production significantly suppressed in PbPb  $\rightarrow$  charm quark E-loss
- ❑  $\Lambda_c^+$  production in pp collisions systematically higher than GM-VFNS calculation (hadronization tuned with  $e^+e^-$  data)
  - Breakdown of the universality of charm quark fragmentation functions?
- ❑ For  $p_T > 10$  GeV/c, the  $\Lambda_c^+/D^0$  ratios for pp and PbPb collisions are consistent, suggesting no significant contribution from coalescence to  $\Lambda_c^+$  hadronization.
  - $\Lambda_c^+/D^0$  ratios for pp and PbPb converge with  $e^+e^-$  for  $p_T > 10$  GeV/c
- ❑ No significant multiplicity dependence is observed for charm hadron production in pPb collisions
  - Different from strange quark, suggests coalescence processes of heavy quarks saturate earlier



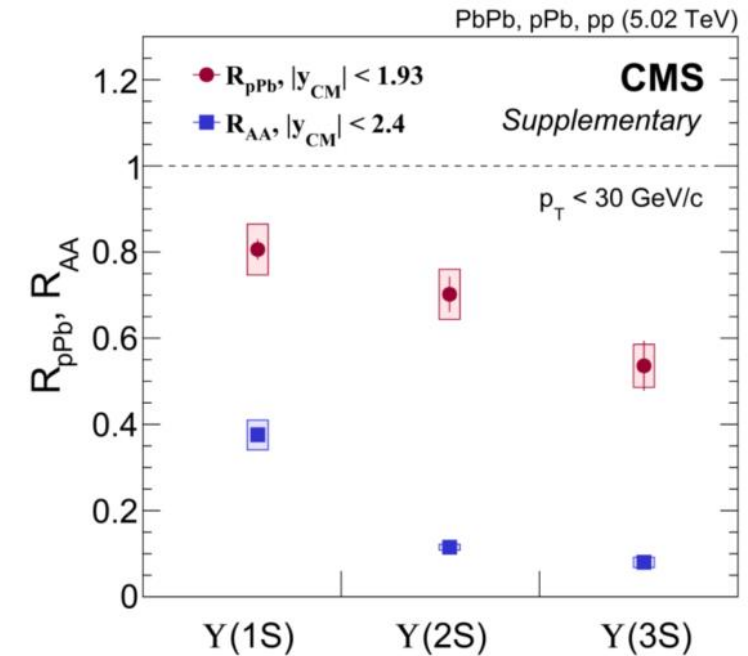


- Excellent performance for muon detection by the CMS experiment
  - Across pp, pPb and PbPb collisions, and wide range of detector occupancies

CMS-PAS-MUO-21-001

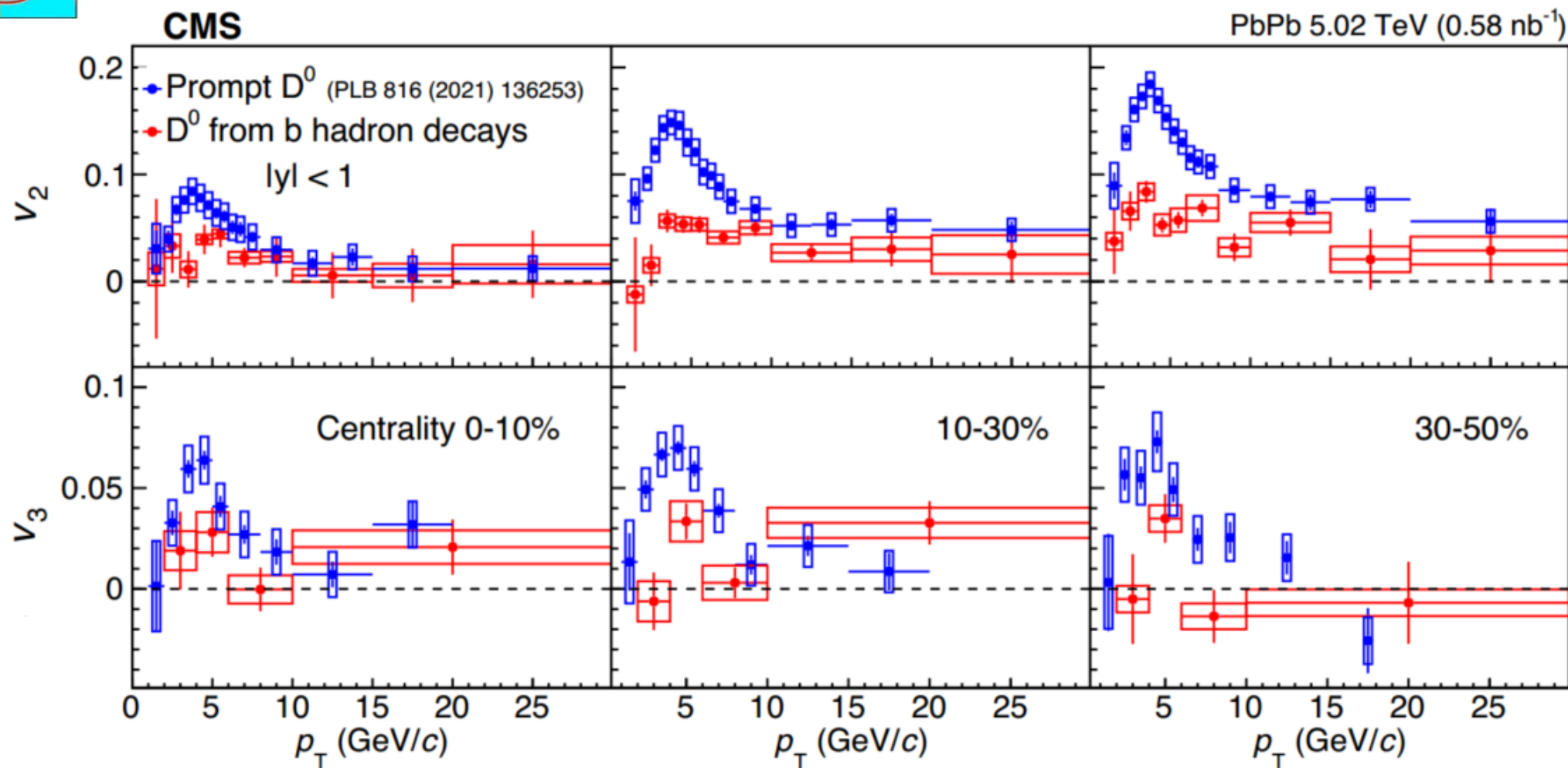
- Y states in pPb collisions
  - Sequential ordering of suppression
    - $R_{pPb}(1S) > R_{pPb}(2S) > R_{pPb}(3S)$
    - Follows their binding energies
    - Challenging to describe with initial state effects alone

- First measurement of Y(3S) in PbPb collisions
  - $R_{PbPb}(3S) = 0.080 \pm 0.014$  (stat)  $\pm 0.012$  (syst)
  - Same ordering as in pPb:  $R_{PbPb}(1S) > R_{PbPb}(2S) > R_{PbPb}(3S)$
  - Overall suppression much larger than in pPb



PbPb (2S) and (3S) [arXiv 2303.17026](https://arxiv.org/abs/2303.17026)

pPb: [PLB 835\(2022\), 137397](https://arxiv.org/abs/2208.137397)



- Strong  $p_T$  and centrality dependence
- Significant nonzero  $v_3$  up to  $\sim 10$  GeV

**VS**

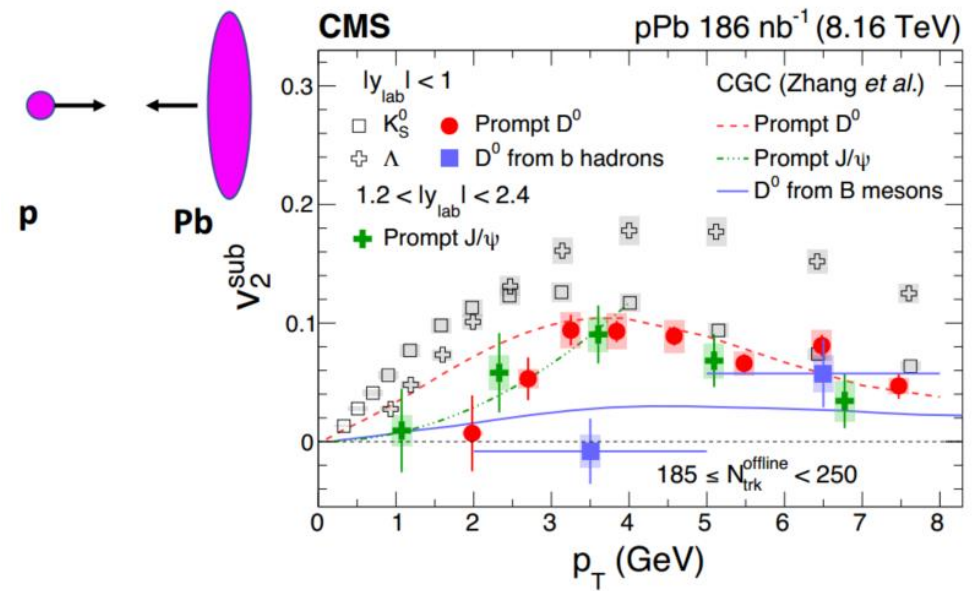
- Strong Weak  $p_T$  and centrality dependence
- Significant Indication of nonzero  $v_3$

□ Mass ordering of flow magnitudes

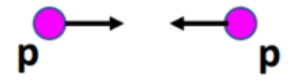


# D<sup>0</sup> flow in small systems

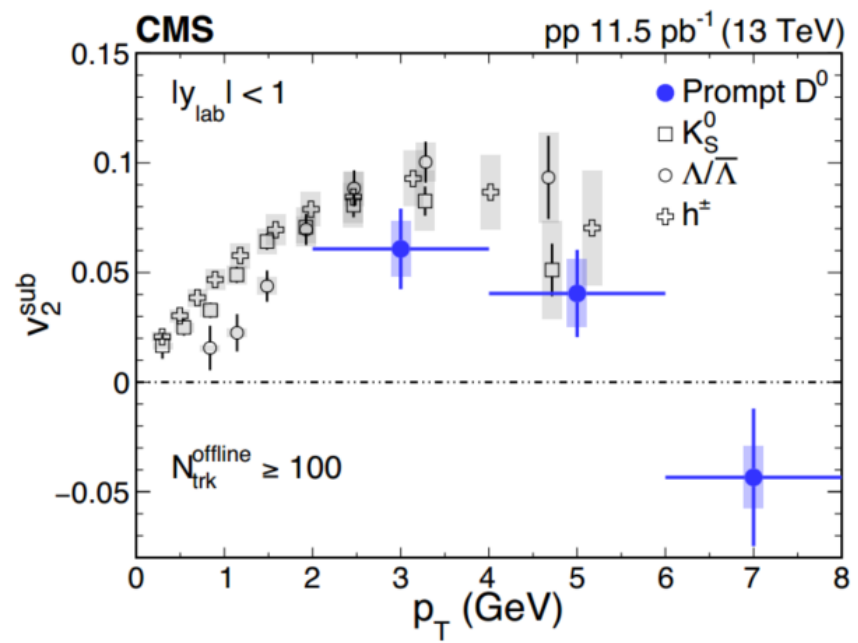
- Open charm (prompt D<sup>0</sup>)  $v_2 \approx$  hidden charm (prompt J/ $\psi$ )
- Flavor hierarchy prompt D<sup>0</sup>  $v_2 >$   $b \rightarrow D^0$   $v_2 \approx 0$
- In agreement with CGC model



○ Significant prompt D<sup>0</sup> flow in small systems



- prompt D<sup>0</sup>  $v_2$  compatible with light hadrons



PLB 813 (2021) 136036

Mass ordering observed for  $p_T < 10$  GeV

- Both flow and  $R_{AA}$
- Both PbPb and pPb

PLB 782 (2018) 474  
PRL 123 (2019) 022001





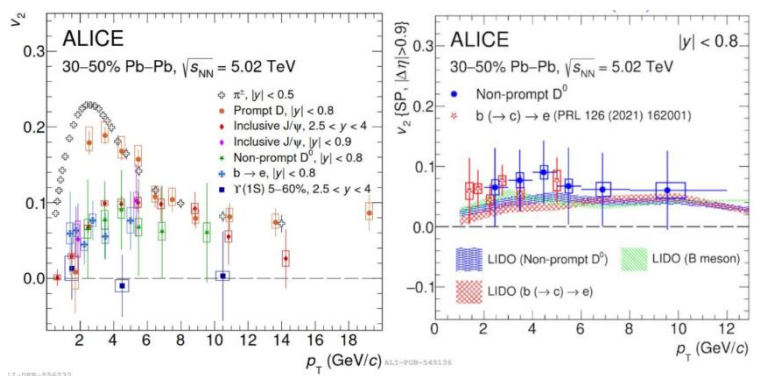
● **Mass-dependent energy loss**

- Observed hierarchy of charm hadron  $R_{AA}$ :  $4 < p_T < 10 \text{ GeV}/c$ 
  - $R_{AA}(\Lambda_c^+) > R_{AA}(D_s^+) > R_{AA}(D) \rightarrow$  recombination and radial flow
- $R_{AA}(\text{non-prompt } D^0) > R_{AA}(\text{prompt } D^0) \rightarrow$  dead-cone effect

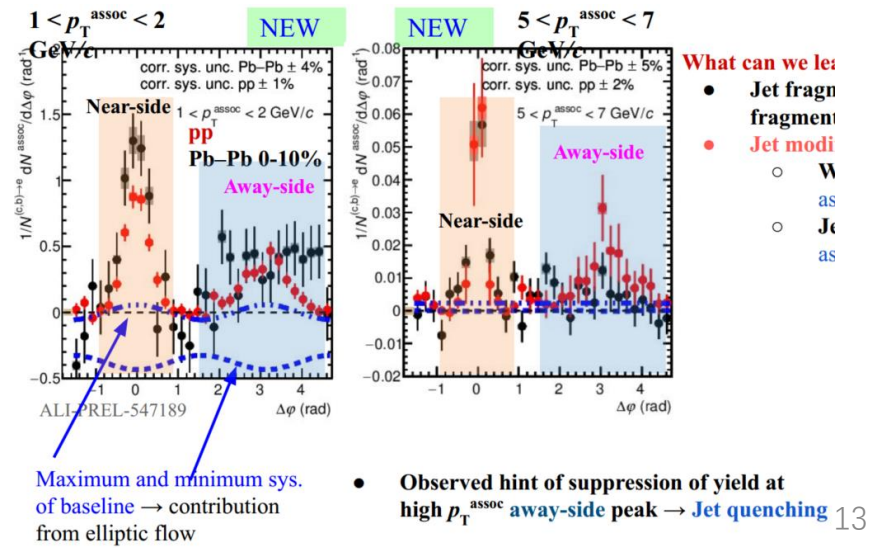
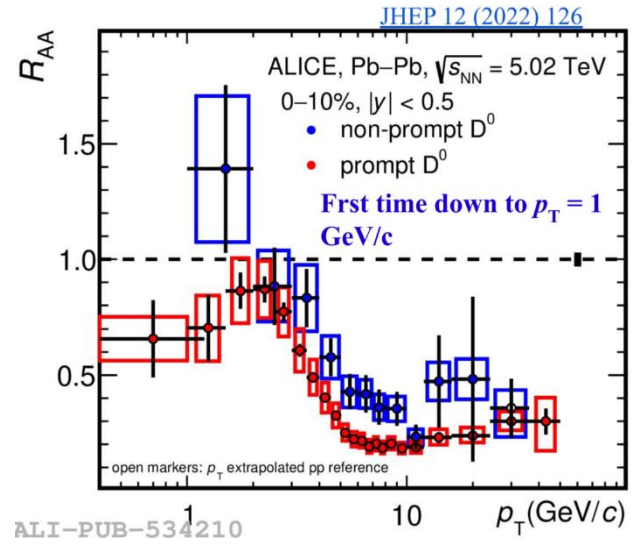
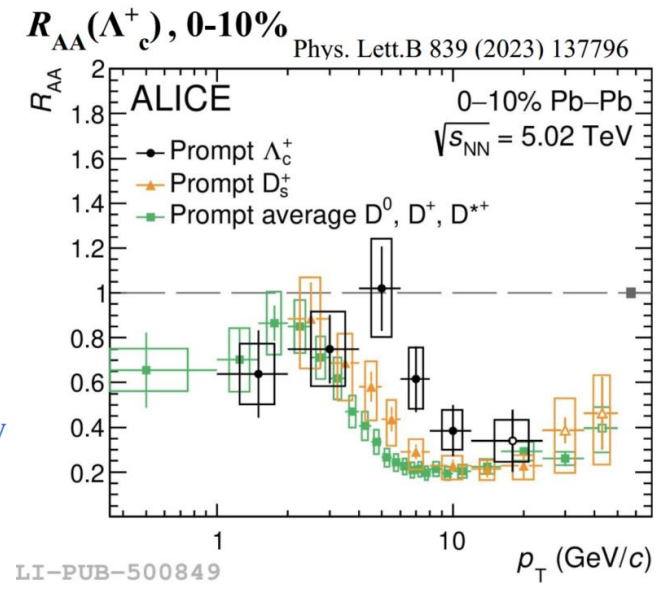
● **Collectivity in heavy-ion collisions**

- Strong coupling of charm quark with QGP constituents at low  $p_T \rightarrow$  charm thermalization in the medium
- Small  $v_2$  for beauty  $\rightarrow$  weaker thermalization
- $v_2 > 0$  for inclusive muon in high multiplicity p-Pb collisions  $\rightarrow$  presence of collective-like effects in small systems also in the HF sector

● **Hint of suppression at high  $p_T$  in away-side HFe-h correlation peak  $\rightarrow$  hint of jet quenching**



- $v_2(\text{non-prompt } D^0) > 0$  with significance of  $2.7\sigma$
- $v_2(\text{non-prompt } D^0) < v_2(\text{prompt } D) >$  with significance of  $3.2\sigma$  in the interval  $2 < p_T < 8 \text{ GeV}/c$



- Observed hint of suppression of yield at high  $p_T^{\text{assoc}}$  away-side peak  $\rightarrow$  Jet quenching

# Heavy-flavour jet substructure for probing the flavour dependences of QCD parton showers with ALICE

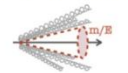
## QCD shower flavour dependence

### Casimir colour factors

Different emission properties due to the different amounts of colour charge carried by quarks and gluons

### The dead-cone effect

A suppression of emissions in a cone of size  $m/E$  around the direction of the emitter



### Quark-initiated shower

Narrower shower profile  
Fewer emissions in the shower



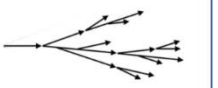
### Light-quark-initiated shower

Narrower shower profile  
Fewer emissions in the shower



### Gluon-initiated shower

Broader shower profile  
Higher number of emissions



### Heavy-quark-initiated shower

Suppression of small angle emissions  
Harder fragmentation

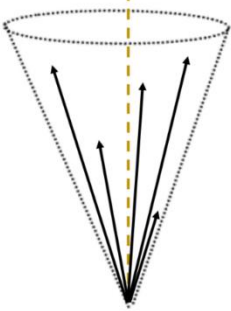


## Heavy-flavour jets

### Well controlled probe

Heavy-flavour jet production is perturbatively calculable down to low  $p_T$

### Charm quark



### Casimir effects

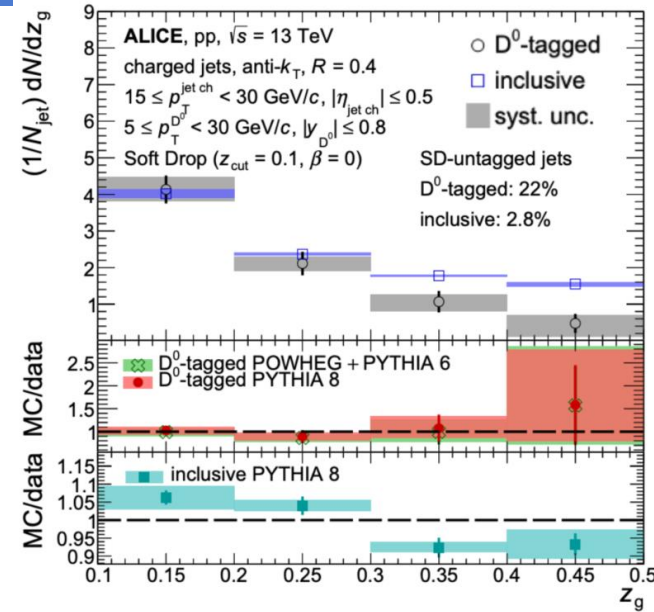
Heavy-flavour jets allow access to a high purity quark sample for jets and splittings

### Clean connection to the shower

The large mass of the heavy-quark suppresses thermal production and production during the process of hadronisation

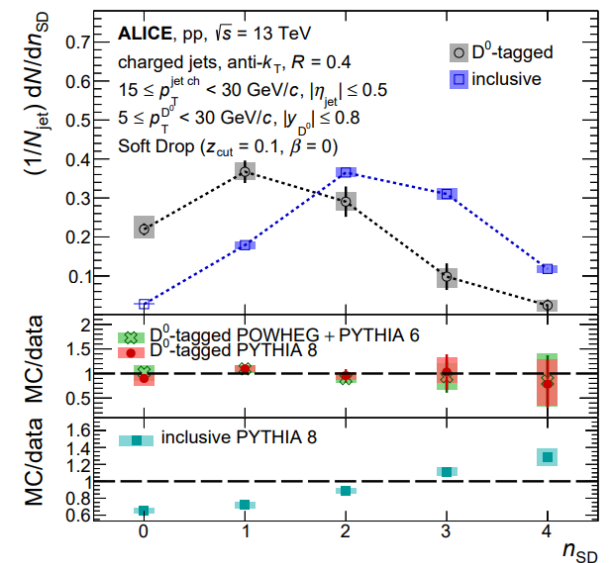
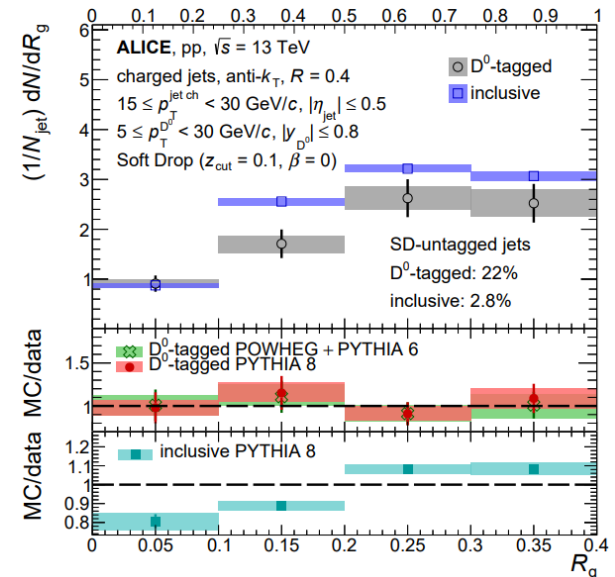
### Mass effects

At low energies heavy-quarks provide unique access to mass effects in the shower



charm-tagged jets have significantly fewer symmetric splittings (large  $z_g$  values), compared to inclusive jets  
~ Expected mass effects in the QCD splitting function

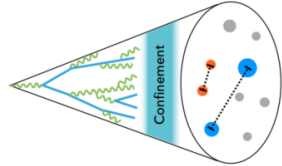
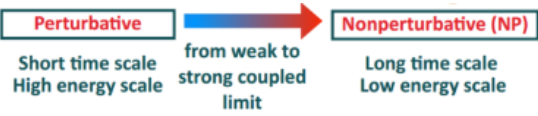
a reduction at large-angles compared to inclusive jets  
~ differences between quark and gluon fragmentation



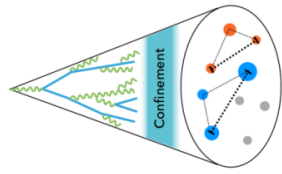
a significant shift to smaller values for the charm-tagged jets  
~ the presence of a dead cone for charm quarks (satisfying the Soft Drop condition)



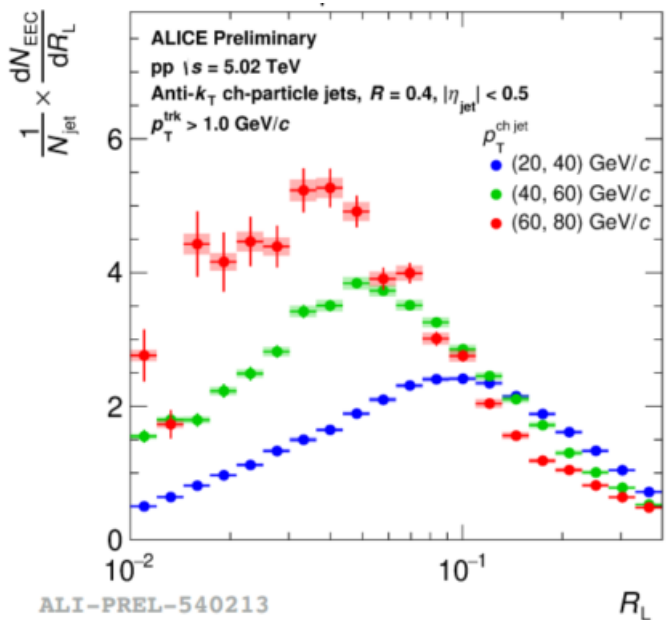
# First energy-energy correlators measurements for inclusive and heavy-flavour tagged jets with ALICE



▶ 2 point energy correlator results in pp at 5 TeV



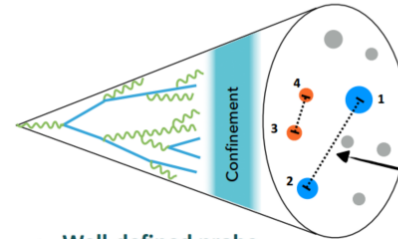
▶ 2 and 3 point energy correlator results in pp at 13 TeV



Different scaling behavior observed in the perturbative (large  $R_L$ ) and NP region (small  $R_L$ )

Transition position shifts to lower  $R_L$  for higher jet  $p_T$  range

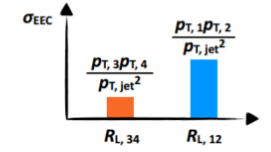
## A new jet substructure observable: 2 point energy correlators



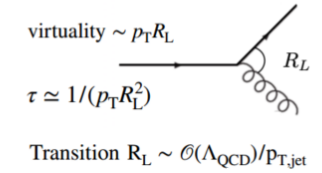
Count number of weighted pairs as function of  $R_L$

$$\frac{d\sigma_{EEC}}{dR_L} = \sum_{i,j} \int d\sigma(R'_L) \frac{p_{T,i} p_{T,j}}{p_{T,jet}^2} \delta(R'_L - R_{L,ij})$$

Energy weight



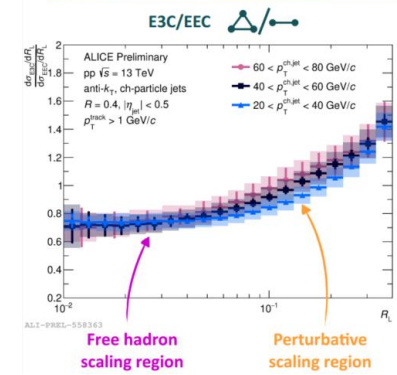
- ▶ Well-defined probe
  - ❖ IRC safe + pQCD calculation available: K. Lee, B. Mecaj, I. Moutl ([arXiv:2205.03414](https://arxiv.org/abs/2205.03414))
  - ❖ Soft contribution (MPI, UE) power suppressed by energy weight: no need for grooming when comparing to pQCD calculation
- ▶ Probing fixed scale with fixed  $R_L$ 
  - ❖ Large  $\rightarrow$  small angle: perturbative  $\rightarrow$  NP scales
  - ❖ When the virtuality approaches  $\mathcal{O}(\Lambda_{QCD})$ , EEC undergo transition into confinement region



$$\frac{d\sigma_{E3C}}{dR_L} = \sum_{i,j,k} \int d\sigma(R'_L) \frac{p_{T,i} p_{T,j} p_{T,k}}{p_{T,jet}^3} \delta(R'_L - R_{L,ijk})$$

Energy weight

## E3C/EEC ratio: isolate perturbative scaling behavior



- ▶ Cancellation of NP effects and systematic uncertainties
  - ❖ Small  $R_L$ : flat shape for free hadron region
  - ❖ Large  $R_L$ : slope sensitive to  $\alpha_s$
- E3C/EEC ratio  $\propto \alpha_s(Q) \ln R_L + \mathcal{O}(\alpha_s^2)$
- High precision constraint on  $\alpha_s$  from jet substructure
- ▶ Slope decrease towards higher jet  $p_T$ 
  - ❖ Higher jet  $p_T \rightarrow$  higher  $Q \rightarrow$  smaller  $\alpha_s \rightarrow$  smaller slope
- Consistent with running of the coupling!