

Helicity analysis of semileptonic Λ decays

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Semileptonic meeting

1 September 2020

Introduction

- Helicity method allows:
 - Compact calculations of the angular decay distributions
 - Analyze the semileptonic decays of polarized hyperon
 - Take into account the lepton mass effects
 $\Rightarrow$ vector and axial-vector currents
- **Aim:** a general modular method for the semileptonic hyperon decays
- Analysis is based on:
 - Helicity analysis for $\Xi^0 \rightarrow \Sigma^+ (\rightarrow p\pi^0) l^- \bar{\nu}_l$ ($l = e^-, \mu^-$) [[EPJ C59 \(2009\) 27](#)]
 - Polarization observables in e^+e^- annihilation to a $B\bar{B}$ pair [[PRD 99 \(2019\) 056008](#)]

Production process of two spin- $\frac{1}{2}$ baryons

- General framework of the $e^+e^- \rightarrow J/\psi \rightarrow \Lambda\bar{\Lambda}$ is described in [\[PRD99 \(2019\) 056008\]](#)
- Production process **doesn't depend** on the final states. It is the same for:
 - $e^+e^- \rightarrow J/\psi \rightarrow \Lambda\bar{\Lambda}$ with $\Lambda \rightarrow p\pi^-$ and $\bar{\Lambda} \rightarrow \bar{p}\pi^+$
 - $e^+e^- \rightarrow J/\psi \rightarrow \Lambda\bar{\Lambda}$ with $\Lambda \rightarrow pe^-\bar{\nu}_e$ and $\bar{\Lambda} \rightarrow \bar{p}\pi^+$

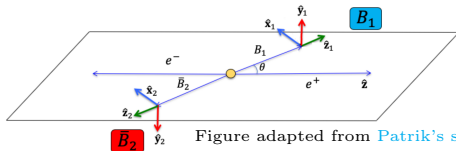


Figure adapted from [Patrik's slides](#)

- Spin density matrix of the production process:

$$\rho_{B_1, \bar{B}_2} = \frac{1}{4} \sum_{\mu, \bar{\nu}=0}^3 C_{\mu\bar{\nu}}(\theta_1) \sigma_{\mu}^{B_1} \otimes \sigma_{\bar{\nu}}^{\bar{B}_2}$$

$$C_{00} = 2(1 + \alpha_{\psi} \cos^2 \theta_1),$$

$$C_{02} = 2\sqrt{1 - \alpha_{\psi}^2} \sin \theta_1 \cos \theta_1 \sin(\Delta\Phi),$$

$$C_{11} = 2 \sin^2 \theta_1,$$

$$C_{13} = 2\sqrt{1 - \alpha_{\psi}^2} \sin \theta_1 \cos \theta_1 \cos(\Delta\Phi),$$

$$C_{20} = -C_{02},$$

$$C_{22} = \alpha_{\psi} C_{11},$$

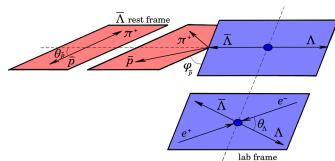
$$C_{31} = -C_{13},$$

$$C_{33} = -2(\alpha_{\psi} + \cos^2 \theta_1).$$

- Main parameters: α_{ψ} , $\Delta\Phi$

Decay chain: $\frac{1}{2} \rightarrow \frac{1}{2} + 0$ (1)

- Full determination in [\[PRD99 \(2019\) 056008\]](#)
- Decay matrix or transition matrix $a_{\mu\nu}$



$$\sigma_\mu \rightarrow \sum_{\nu=0}^3 a_{\mu\nu} \sigma_\nu^d$$

- $J = 1/2$ hyperon ($\bar{\Lambda}$) decays into a $J = 1/2$ baryon ($\bar{p}$) and a $J = 0$ pseudoscalar (π^+)

$$a_{\mu\nu} = \frac{1}{4\pi} \sum_{\lambda, \lambda'=-1/2}^{1/2} B_\lambda B_{\lambda'}^* \sum_{\kappa, \kappa'=-1/2}^{1/2} (\sigma_\mu)^{\kappa, \kappa'} (\sigma_\nu)^{\lambda', \lambda} \mathcal{D}_{\kappa, \lambda}^{1/2*}(\Omega) \mathcal{D}_{\kappa', \lambda'}^{1/2}(\Omega)$$

- $B_\lambda, B_{\lambda'}$ - helicity amplitudes, $\{B_{\frac{1}{2}}, B_{-\frac{1}{2}}\}$
- κ, κ' - index of mother hyperon ($\bar{\Lambda}$); λ, λ' - index of daughter baryon ($\bar{p}$)
- $\Omega = \{\phi, \theta, 0\}$

Decay chain: $\frac{1}{2} \rightarrow \frac{1}{2} + 0$ (2)

- Relation between helicity amplitudes and decay parameters

$$\text{normalization } |A_S|^2 + |A_P|^2 = |B_{-\frac{1}{2}}|^2 + |B_{\frac{1}{2}}|^2 = 1,$$

$$\alpha_D = -2\Re(A_S^* A_P) = |B_{1/2}|^2 - |B_{-1/2}|^2,$$

$$\beta_D = -2\Im(A_S^* A_P) = 2\Im(B_{1/2} B_{-1/2}^*),$$

$$\gamma_D = |A_S|^2 - |A_P|^2 = 2\Re(B_{1/2} B_{-1/2}^*),$$

$$\text{where } \beta_D = \sqrt{1 - \alpha_D^2} \sin \varphi_D \text{ and } \gamma_D = \sqrt{1 - \alpha_D^2} \cos \varphi_D$$

- Non-zero elements of the decay matrix $a_{\mu\nu}$

$$a_{00} = 1,$$

$$a_{21} = \beta_D \cos \varphi + \gamma_D \cos \theta \sin \varphi,$$

$$a_{03} = \alpha_D,$$

$$a_{22} = \gamma_D \cos \varphi - \beta_D \cos \theta \sin \varphi,$$

$$a_{10} = \alpha_D \cos \varphi \sin \theta,$$

$$a_{23} = \sin \theta \sin \varphi,$$

$$a_{11} = \gamma_D \cos \theta \cos \varphi - \beta_D \sin \varphi,$$

$$a_{30} = \alpha_D \cos \theta,$$

$$a_{12} = -\beta_D \cos \theta \cos \varphi - \gamma_D \sin \varphi,$$

$$a_{31} = -\gamma_D \sin \theta,$$

$$a_{13} = \sin \theta \cos \varphi,$$

$$a_{32} = \beta_D \sin \theta,$$

$$a_{20} = \alpha_D \sin \theta \sin \varphi,$$

$$a_{33} = \cos \theta$$

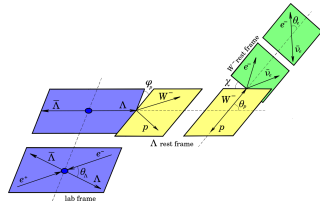
- Main parameters: α_D , ϕ_D

Decay chain: $\frac{1}{2} \rightarrow \frac{1}{2} + \{0, \pm 1\}$ (1)

- Decay matrix or transition matrix $b_{\mu\nu}$

$$\sigma_\mu \rightarrow \sum_{\nu=0}^3 b_{\mu\nu} \sigma_\nu^{d-W}$$

- $J = 1/2$ hyperon (Λ) decays into a $J = 1/2$ baryon (p) and a $J = \{0, \pm 1\}$ W^- -boson ($\rightarrow l\nu_l$)



$$b_{\mu\nu} = \frac{1}{4\pi} \sum_{\lambda_2=-1/2}^{1/2} \sum_{\lambda_W, \lambda'_W=-1}^1 H_{\lambda_2 \lambda_W} H_{\lambda_2 \lambda'_W}^* \sum_{\kappa, \kappa'=-1/2}^{1/2} (\sigma_\mu)^{\kappa, \kappa'} (\sigma_\nu)^{\lambda_2 - \lambda'_W, \lambda_2 - \lambda_W} \mathcal{D}_{\kappa, \lambda_2 - \lambda_W}^{1/2*}(\Omega) \mathcal{D}_{\kappa', \lambda_2 - \lambda'_W}^{1/2}(\Omega) \times$$

$$\sum_{\lambda_l, \lambda_\nu=-1/2}^{1/2} |h_{\lambda_l \lambda_\nu=\pm 1/2}^l|^2 \mathcal{D}_{\lambda_W, \lambda_l - \lambda_\nu}^{1*}(\Omega') \mathcal{D}_{\lambda'_W, \lambda_l - \lambda_\nu}^1(\Omega'),$$

- $H_{\lambda_2, \lambda_W}, H_{\lambda_2, \lambda'_W}$ - helicity amplitudes, $\{H_{\frac{1}{2}0}, H_{-\frac{1}{2}0}, H_{\frac{1}{2}1}, H_{-\frac{1}{2}-1}\}$
- κ, κ' - index of mother hyperon (Λ); λ_2 - index of daughter baryon (p)
- λ_W, λ'_W - index of W^- -boson; λ_l, λ_ν - index of lepton and neutrino
- $\Omega = \{\varphi, \theta, 0\}, \Omega' = \{\chi, \theta_l, 0\}$

Helicity amplitudes of the lepton pair $h_{\lambda_l \lambda_\nu}^l$

- Lepton and antineutrino spinors

$$\bar{u}_{l^-}(\mp \frac{1}{2}, p_{l^-}) = \sqrt{E_l + m_l} \left(\chi_{\mp}^{\dagger}, \frac{\pm |\vec{p}_{l^-}|}{E_l + m_l} \chi_{\mp}^{\dagger} \right),$$

$$v_{\bar{\nu}}(\frac{1}{2}) = \sqrt{E_{\nu}} \begin{pmatrix} \chi_{+} \\ -\chi_{+} \end{pmatrix},$$

where $\chi_{+} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ and $\chi_{-} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ are Pauli two-spinors

- SM form of the lepton current ($\lambda_W = \lambda_{l^-} - \lambda_{\bar{\nu}}$)

$$h_{\lambda_{l^-}=\mp 1/2, \lambda_{\bar{\nu}}=1/2}^l = \bar{u}_{l^-}(\mp \frac{1}{2}) \gamma^{\mu} (1 + \gamma_5) v_{\bar{\nu}}(\frac{1}{2}) \begin{Bmatrix} \epsilon_{\mu}(-1) \\ \epsilon_{\mu}(0) \end{Bmatrix}$$

where $\epsilon^{\mu}(0) = (0; 0, 0, 1)$ and $\epsilon^{\mu}(\mp 1) = (0; \mp 1, -i, 0)/\sqrt{2}$

- Moduli squared of the helicity amplitudes [\[EPJ C59 \(2009\) 27\]](#)

$$\text{nonflip}(\lambda_W = \mp 1) : |h_{\lambda_{l^-}=\mp \frac{1}{2}, \lambda_{\bar{\nu}}=\pm \frac{1}{2}}^l|^2 = 8(q^2 - m_l^2),$$

$$\text{flip}(\lambda_W = 0) : |h_{\lambda_{l^-}=\pm \frac{1}{2}, \lambda_{\bar{\nu}}=\pm \frac{1}{2}}^l|^2 = 8 \frac{m_l^2}{2q^2} (q^2 - m_l^2)$$

- Upper and lower signs refer to the configurations $(l^-, \bar{\nu}_l)$ ($\lambda_{\nu} = 1/2$) and (l^+, ν_l) ($\lambda_{\nu} = -1/2$), respectively
- In case of [the e-mode](#) only [nonflip transition](#) remains under assumption $\frac{m_e^2}{2q^2} \rightarrow 0$

Decay chain: $\frac{1}{2} \rightarrow \frac{1}{2} + \{0, \pm 1\}$ (2)

- Relation between helicity amplitudes and decay parameters

$$\text{normalization } \frac{1}{4}(1 \mp \cos \theta_l)^2 |H_{\frac{1}{2}1}|^2 + \frac{1}{4}(1 \pm \cos \theta_l)^2 |H_{-\frac{1}{2}-1}|^2 + \frac{1}{2} \sin^2 \theta_l (|H_{-\frac{1}{2}0}|^2 + |H_{\frac{1}{2}0}|^2) = 1,$$

$$\alpha_D^{sl} = \frac{1}{4}(1 \mp \cos \theta_l)^2 |H_{\frac{1}{2}1}|^2 - \frac{1}{4}(1 \pm \cos \theta_l)^2 |H_{-\frac{1}{2}-1}|^2 + \frac{1}{2} \sin^2 \theta_l (|H_{-\frac{1}{2}0}|^2 - |H_{\frac{1}{2}0}|^2),$$

$$\beta_D^{sl} = \frac{1}{\sqrt{2}} \sin \theta_l \sin \chi ((1 \pm \cos \theta_l) \Im(H_{-\frac{1}{2}0}^* H_{-\frac{1}{2}-1}) + (1 \mp \cos \theta_l) \Im(H_{\frac{1}{2}1}^* H_{\frac{1}{2}0})),$$

$$\gamma_D^{sl} = \frac{1}{\sqrt{2}} \sin \theta_l \cos \chi ((1 \pm \cos \theta_l) \Re(H_{-\frac{1}{2}0}^* H_{-\frac{1}{2}-1}) + (1 \mp \cos \theta_l) \Re(H_{\frac{1}{2}1}^* H_{\frac{1}{2}0}))$$

$$\text{where } \beta_D^{sl} = \sqrt{1 - (\alpha_D^{sl})^2} \sin \phi_D^{sl} \text{ and } \gamma_D^{sl} = \sqrt{1 - (\alpha_D^{sl})^2} \cos \phi_D^{sl}$$

- Non-zero elements of the decay matrix $b_{\mu\nu}$

$$b_{00} = 1, \quad b_{21} = \mp \gamma_D^{sl} \cos \theta \sin \phi \mp \beta_D^{sl} \cos \phi,$$

$$b_{03} = \alpha_D^{sl}, \quad b_{22} = \pm \beta_D^{sl} \cos \theta \sin \phi \mp \gamma_D^{sl} \cos \phi,$$

$$b_{10} = \alpha_D^{sl} \cos \phi \sin \theta, \quad b_{23} = \sin \theta \sin \phi,$$

$$b_{11} = \mp \gamma_D^{sl} \cos \theta \cos \phi \pm \beta_D^{sl} \sin \phi, \quad b_{30} = \alpha_D^{sl} \cos \theta,$$

$$b_{12} = \pm \beta_D^{sl} \cos \theta \cos \phi \pm \gamma_D^{sl} \sin \phi, \quad b_{31} = \pm \gamma_D^{sl} \sin \theta,$$

$$b_{13} = \sin \theta \cos \phi, \quad b_{32} = \mp \beta_D^{sl} \sin \theta,$$

$$b_{20} = \alpha_D^{sl} \sin \theta \sin \phi, \quad b_{33} = \cos \theta$$

- Each element of $b_{\mu\nu}$ is multiplied by $q^2 p$ where $p = \sqrt{M_+ M_-} / (2M_1)$
- Main parameters: α_D^{sl} , ϕ_D^{sl}

Semileptonic Λ decay

- Momenta and masses: $\Lambda(p_1, M_1) \rightarrow p(p_2, M_2) + e^-(p_e, m_e) + \bar{\nu}_e(p_{\bar{\nu}_e}, 0)$
- Standard set of invariant FF for the weak current-induced baryonic $1/2^+ \rightarrow 1/2^+$ transitions:

$$M_\mu = M_\mu^V + M_\mu^A = \langle p(p_2) | J_\mu^{V+A} | \Lambda(p_1) \rangle = \\ = \bar{u}(p_2) \left[\gamma_\mu (F_1^V(q^2) + F_1^A(q^2) \gamma_5) + \frac{i\sigma_{\mu\nu} q^\nu}{M_1} (F_2^V(q^2) + F_2^A(q^2) \gamma_5) + \frac{q_\mu}{M_1} (F_3^V(q^2) + F_3^A(q^2) \gamma_5) \right] u(p_1)$$

where $q_\mu = (p_1 - p_2)_\mu$

- For $\Lambda \rightarrow p e^- \bar{\nu}_e$ at $\mathcal{O}(\frac{m_e^2}{2q^2}) \rightarrow 0 \Rightarrow F_3^{V,A} \rightarrow 0$

Semileptonic Λ decay

- Momenta and masses: $\Lambda(p_1, M_1) \rightarrow p(p_2, M_2) + e^-(p_e, m_e) + \bar{\nu}_e(p_{\bar{\nu}_e}, 0)$
- Standard set of invariant FF for the weak current-induced baryonic $1/2^+ \rightarrow 1/2^+$ transitions:

$$M_\mu = M_\mu^V + M_\mu^A = \langle p(p_2) | J_\mu^{V+A} | \Lambda(p_1) \rangle =$$

$$= \bar{u}(p_2) \left[\gamma_\mu (F_1^V(q^2) + F_1^A(q^2) \gamma_5) + \frac{i\sigma_{\mu\nu} q^\nu}{M_1} (F_2^V(q^2) + F_2^A(q^2) \gamma_5) + \frac{q_\mu}{M_1} (F_3^V(q^2) + F_3^A(q^2) \gamma_5) \right] u(p_1)$$

where $q_\mu = (p_1 - p_2)_\mu$

- For $\Lambda \rightarrow p e^- \bar{\nu}_e$ at $\mathcal{O}(\frac{m_e^2}{2q^2}) \rightarrow 0 \Rightarrow F_3^{V,A} \rightarrow 0$
- $\Lambda \rightarrow p e^- \bar{\nu}_e \Rightarrow \Lambda \rightarrow p W^- (\rightarrow e^- \bar{\nu}_e)$
- Helicity amplitude is $H_{\lambda_2 \lambda_W} = (H_{\lambda_2 \lambda_W}^V + H_{\lambda_2 \lambda_W}^A)$ with $(\lambda_2 = \pm 1/2; \lambda_W = 0, \pm 1)$:

$$\begin{array}{l} \text{vector} \left(\begin{array}{l} H_{\frac{1}{2}1}^V = \sqrt{2M_-} \left(-F_1^V - \frac{M_1 + M_2}{M_1} F_2^V \right), \\ H_{\frac{1}{2}0}^V = \frac{\sqrt{M_-}}{\sqrt{q^2}} \left((M_1 + M_2) F_1^V + \frac{q^2}{M_1} F_2^V \right), \end{array} \right. \end{array}$$

$$\begin{array}{l} \text{axial-vector} \left(\begin{array}{l} H_{\frac{1}{2}1}^A = \sqrt{2M_+} \left(F_1^A - \frac{M_1 - M_2}{M_1} F_2^A \right), \\ H_{\frac{1}{2}0}^A = \frac{\sqrt{M_+}}{\sqrt{q^2}} \left(-(M_1 - M_2) F_1^A + \frac{q^2}{M_1} F_2^A \right). \end{array} \right. \end{array}$$

$$\text{where } M_\pm = (M_1 \pm M_2)^2 - q^2; \quad H_{-\lambda_2, -\lambda_W}^{V,A} = \pm H_{\lambda_2 \lambda_W}^{V,A}$$

Joint angular distribution

- Process $e^+e^- \rightarrow J/\psi \rightarrow \Lambda \bar{\Lambda} \rightarrow pe^- \bar{\nu}_e \bar{p} \pi^+$

$$\text{Tr} \rho_{pW\bar{p}} \propto \mathcal{W}(\xi; \omega) = \sum_{\mu, \bar{\nu}=0}^3 C_{\mu\bar{\nu}} b_{\mu 0}^{\Lambda} a_{\bar{\nu} 0}^{\bar{\Lambda}}$$

- $C_{\mu\bar{\nu}}(\theta_{\Lambda}; \alpha_{\psi}, \Delta\Phi)$
- $b_{\mu 0}$ matrices for $1/2 \rightarrow 1/2 + \{0, \pm 1\}$ decays $\Leftrightarrow b_{\mu 0}^{\Lambda} \equiv b_{\mu 0}(\theta_p, \varphi_p, \theta_e, \chi, q^2; \alpha_{\Lambda}^{sl})$
- $a_{\bar{\nu} 0}$ matrices for $1/2 \rightarrow 1/2 + 0$ decays $\Leftrightarrow a_{\bar{\nu} 0}^{\bar{\Lambda}} \equiv a_{\bar{\nu} 0}(\theta_{\bar{p}}, \varphi_{\bar{p}}; \alpha_{\bar{\Lambda}})$
 - $\xi \equiv (\theta_{\Lambda}, \theta_p, \varphi_p, \theta_e, q^2, \chi, \theta_{\bar{p}}, \varphi_{\bar{p}})$
 - $\omega \equiv (\alpha_{\psi}, \Delta\Phi, \alpha_{\Lambda}^{sl}, \alpha_{\bar{\Lambda}})$

- Elements of decay matrix $b_{\mu 0}$ and $a_{\bar{\nu} 0}$

$$\left. \begin{aligned} b_{00} &= 1, \\ b_{10} &= \alpha_D^{sl} \cos \phi \sin \theta, \\ b_{20} &= \alpha_D^{sl} \sin \theta \sin \phi, \\ b_{30} &= \alpha_D^{sl} \cos \theta \end{aligned} \right\} \Lambda \rightarrow pe^- \bar{\nu}_e \text{ where } \alpha_D^{sl} \equiv \alpha_{\Lambda}^{sl} \quad \left. \begin{aligned} a_{00} &= 1, \\ a_{10} &= \alpha_D \cos \varphi \sin \theta, \\ a_{20} &= \alpha_D \sin \theta \sin \varphi, \\ a_{30} &= \alpha_D \cos \theta \end{aligned} \right\} \bar{\Lambda} \rightarrow \bar{p} \pi^+ \text{ where } \alpha_D \equiv \alpha_{\bar{\Lambda}}$$

* each $b_{\mu 0}$ are multiplied by $q^2 p$ where $p = \sqrt{M_+ M_-} / (2M_1)$

[PRD 99 (2019) 056008]

Next steps

- Verify the correctness of the modular method for the semileptonic hyperon decays
 - Plan to use Andrzej's method (some package of scripts)
- Check the direction of particles in the helicity frame of the initial hyperon
- Expand the modular method to the semileptonic cascade decays, as an example:
 - $\Xi \rightarrow \Lambda(\rightarrow p + \pi) + l + \nu_l$
 - $\Xi \rightarrow \Sigma(\rightarrow p + \pi) + l + \nu_l$
 - ...

Backups



"I ALWAYS BACK UP EVERYTHING."

Form factors

$$F_i^{V,A}(q^2) = F_i^{V,A}(0) \prod_{n=0}^{n_i} \frac{1}{1 - \frac{q^2}{m_{V,A}^2 + n\alpha'^{-1}}} \approx F_i^{V,A}(0) \left(1 + q^2 \sum_{n=0}^{n_i} \frac{1}{m_{V,A}^2 + n\alpha'^{-1}} \right)$$

	$F_i^{V,A}(0)$	$m_{V,A}$	α' [GeV ⁻²]	n_i
$F_1^V(q^2)$	$-\sqrt{\frac{3}{2}}$ ¹	$m_{K^*(892)}^0$ ($J^P = 1^-$)	0.9	$n_1 = 1$
$F_2^V(q^2)$	$\frac{M_\Lambda \mu_p}{2M_p} F_1^V(0)$ ²			$n_2 = 2$
$F_3^V(q^2)$	0 ⁴			$n_3 = 2$
$F_1^A(q^2)$	$0.719 F_1^V(0)$ ³	$m_{K^*(1270)}^0$ ($J^P = 1^+$)		$n_1 = 1$
$F_2^A(q^2)$	0 ⁴	—		$n_2 = 2$
$F_3^A(q^2)$	$\frac{M_\Lambda(M_\Lambda + M_p)}{(m_K^-)^2} F_1^A(0)$ ⁴	m_K ($J^P = 0^-$)		$n_3 = 2$

● ¹ [PR135(1964)B1483], [PRL13(1964)264]

● ² $\mu_p = 1.793$ [Lect.NotesPhys.222(1985)1], [Ann.Rev.Nucl.Part.Sci.53(2003)39], [JHEP0807(2008)132]

● ³ [PRD41(1990)780]

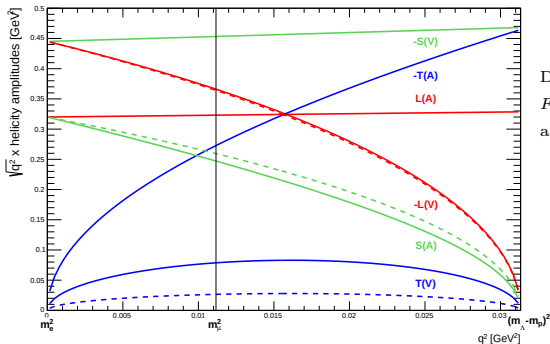
● ⁴ Vanish in the $SU(3)$ symmetry limit; Goldberger-Treiman relation [PR110(1958)1178], [PR111(1958)354]

Size estimations of helicity amplitudes

$$T(V,A) := \sqrt{q^2} H_{\frac{1}{2}1}^{V,A}$$

$$L(V,A) := \sqrt{q^2} H_{\frac{1}{2}0}^{V,A}$$

$$S(V,A) := \sqrt{q^2} H_{\frac{1}{2}t}^{V,A}$$



Dashed lines:
 $F_2^V(q^2)$ and $F_3^A(q^2)$
 are switched off

- If $q^2 = m_e^2 \Rightarrow H_{\frac{1}{2}0}^V$ and $H_{\frac{1}{2}0}^A$ are dominated
- If $q^2 = (M_\Lambda - M_p)^2 \Rightarrow H_{\frac{1}{2}1}^A = -\sqrt{2}H_{\frac{1}{2}0}^A$ are dominated
- Using data of the E-555 experiment (Fermilab) [\[PRD41 \(1990\) 780\]](#)

$$\Rightarrow \frac{F_1^A(0)}{F_1^V(0)} (\Lambda \rightarrow pe^- \bar{\nu}_e) = 0.731 \pm 0.016 \text{ and } \frac{F_2^V(0)}{F_1^V(0)} (\Lambda \rightarrow pe^- \bar{\nu}_e) = 0.15 \pm 0.30$$

$$\Rightarrow \frac{F_1^A(0)}{F_1^V(0)} (\Lambda \rightarrow pe^- \bar{\nu}_e) = 0.719 \pm 0.016 \text{ with constraint } \frac{F_2^V(0)}{F_1^V(0)} (\Lambda \rightarrow pe^- \bar{\nu}_e) \rightarrow 0.97 \text{ (CVC)}$$

Semileptonic Λ decay (2)

- Angular distributions for the $\Lambda \rightarrow pW^- (\rightarrow e^- \bar{\nu}_e)$ of a polarized Λ in the helicity formalism [EPJ C59 (2009) 27]

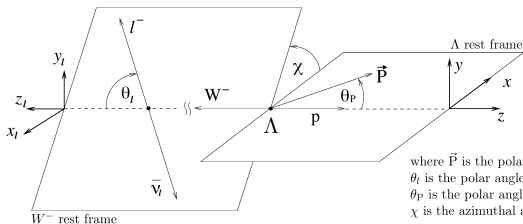
$$\frac{d\Gamma}{dq^2 d \cos \theta_l d\chi d \cos \theta_p} = \frac{1}{6} \frac{G_F^2}{(2\pi)^4} |V_{us}|^2 \frac{q^2 p}{8M_\Lambda^2} \times$$

$$\times \left[\frac{3}{8} (1 - \cos \theta_l)^2 |H_{\frac{1}{2}1}|^2 \rho_{-\frac{1}{2}-\frac{1}{2}} + \frac{3}{8} (1 + \cos \theta_l)^2 |H_{-\frac{1}{2}-1}|^2 \rho_{\frac{1}{2}\frac{1}{2}} + \frac{3}{4} \sin^2 \theta_l (|H_{-\frac{1}{2}0}|^2 \rho_{-\frac{1}{2}-\frac{1}{2}} + |H_{\frac{1}{2}0}|^2 \rho_{\frac{1}{2}\frac{1}{2}}) + \right.$$

$$\left. + \frac{3}{2\sqrt{2}} \rho_{\frac{1}{2}-\frac{1}{2}} \cos \chi \sin \theta_l \left(H_{-\frac{1}{2}0} H_{-\frac{1}{2}-1} (1 + \cos \theta_l) + H_{\frac{1}{2}0} H_{\frac{1}{2}1} (1 - \cos \theta_l) \right) \right]$$

where spin density matrix

$$\rho_{\lambda_p - \lambda_W, \lambda_p - \lambda'_W}(\underline{\theta_p}) = \frac{1}{2} \begin{pmatrix} 1 + P \cos \theta_p & P \sin \theta_p \\ P \sin \theta_p & 1 - P \cos \theta_p \end{pmatrix} \Rightarrow \rho_{\lambda_p - \lambda_W, \lambda_p - \lambda'_W} = \frac{1}{2} \begin{pmatrix} I_0 + I_z & I_x - iI_y \\ I_x + iI_y & I_0 - I_z \end{pmatrix}$$



where $\vec{P}$ is the polarization vector of the Λ
 θ_l is the polar angle of the lepton in the W^- decay plane
 θ_p is the polar angle of the Λ polarization vector in the Λ decay plane
 χ is the azimuthal angle between two decay planes

MC generator (step 1)

- Using `YYbar_example` package by Patrik, the presented process can be generated

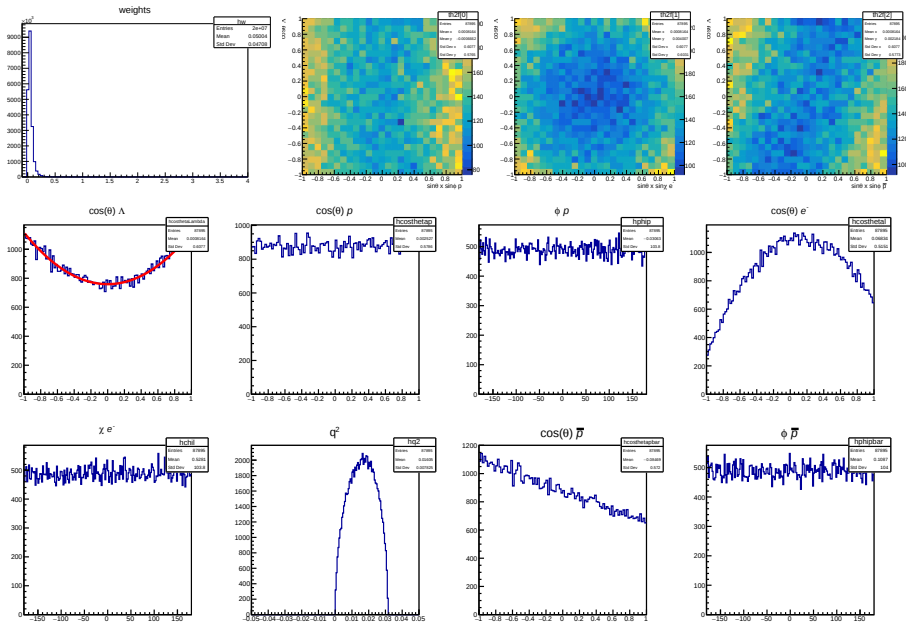
- Generate $N_{\text{evt}}^{\text{sig}}=10^6$ and $N_{\text{evt}}^{\text{phsp}}=10^7$
- Generate $q^2 \in (m_e^2, (M_\Lambda - M_p)^2)$

```
pp[0] = 0.461; // alpha_J/psi (arxiv:1808.08917)
pp[1] = 0.74; // Delta_Phi (arxiv:1808.08917, equal 42.4deg)
pp[2] = 0.719; // g_av=FA1_0/FV1_0 Lambda->p e- nu_ebar
pp[4] = 1.066; // g_w=FV2_0/FV1_0 Lambda->p e- nu_ebar
pp[4] = -0.758; // alpha_D LambdaBar->pbar pi+ (arxiv:1808.08917)
```

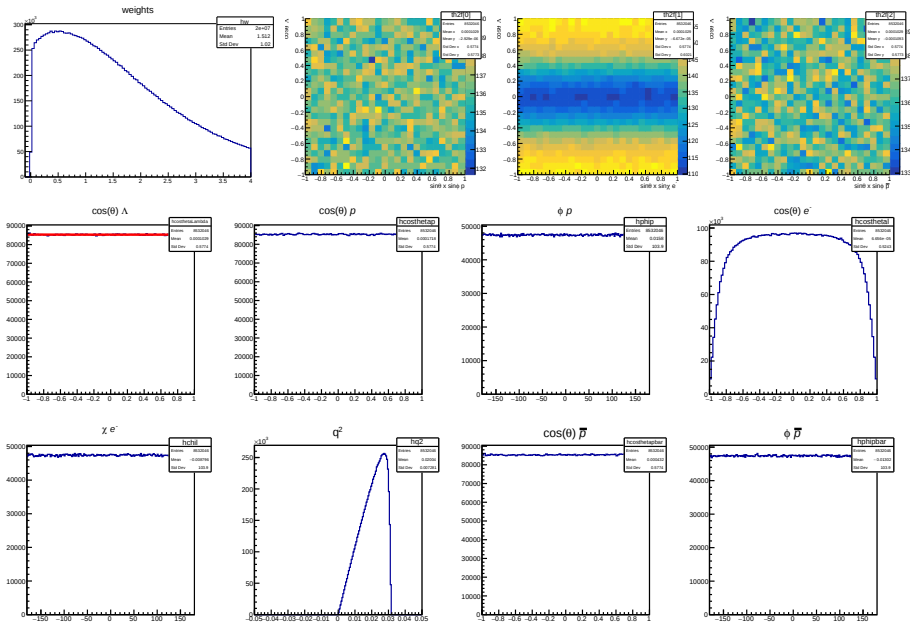
- No negative weights

```
Min: 1.09107e-06 Max:5.72174
nr of negative weights are: 0
FCN=91.6295 FROM MIGRAD STATUS=CONVERGED 104 CALLS 105 TOTAL
EDM=1.00386e-07 STRATEGY= 1 ERROR MATRIX ACCURATE
EXT PARAMETER STEP FIRST
NO. NAME VALUE ERROR SIZE DERIVATIVE
1 p0 7.59231e+02 4.25138e+00 1.20418e-02 -4.15449e-05
2 p1 4.69479e-01 1.55037e-02 4.39105e-05 -3.71647e-02
```

Random samples (signal) (step 1)



Random samples (phsp) (step 1)



Run through fit method (step 2)

- Set starting values in the fit

```
void mainMLL(){
    // instantiating the values to be measured
    Double_t pp[4];
    for( Int_t i = 0; i < 4; i++ ) pp[i]=0;

    // starting values for fit
    Double_t alpha_jpsi = 0.40; // alpha_J/Psi
    Double_t dphi_jpsi = 0.74; // relative phase, Dphi_J/Psi
    Double_t gav_lam_plnu = 0.73; // FA1_0/FV1_0 (Lam->p l nubar_l)
    Double_t gv_lam_plnu = 1.000; // FV2_0/FV1_0 (Lam->p l nubar_l)
    Double_t alpha_lam_pbarpip = -0.758; // alpha (Lambar->pbar pi+)

    alpha_jpsi = gRandom->Rndm();
    dphi_jpsi = gRandom->Rndm();
    gav_lam_plnu = 0.73;
    gv_lam_plnu = 1.000;
    alpha_lam_pbarpip = -0.758;

    ReadData();
    ReadMC();
}
```

$\Lambda \rightarrow pe \bar{\nu}_e$
 $\Lambda \rightarrow p \mu^- \bar{\nu}_\mu$
 V_{us}
 $\alpha_{\mu p} = 1.793$ (recent value is μ_p)
 helicity amplitudes [q]

- Fit method output

```
Loglike: -19457.7
FCN=-19458.2 FROM MINOS STATUS=SUCCESSFUL 254 CALLS 409 TOTAL
EDM=8.92012e-06 STRATEGY= 1 ERROR MATRIX ACCURATE
```

EXT NO.	PARAMETER NAME	VALUE	PARABOLIC ERROR	MINOS NEGATIVE	MINOS POSITIVE
1	alpha_jpsi	4.58133e-01	9.71892e-03	-9.67526e-03	9.76061e-03
2	dphi_jpsi	7.07653e-01	1.92031e-02	-1.90374e-02	1.93709e-02
3	gaLam	9.83858e-01	1.21829e-02	-1.20563e-02	1.23126e-02
4	gvLam	1.30270e+00	4.44752e-02	-4.44779e-02	4.44760e-02
5	aLambar	-7.59457e-01	8.10586e-03	-8.07343e-03	8.13534e-03

ERR DEF= 0.5

```
EXTERNAL ERROR MATRIX. NDIM= 25 NPAR= 5 ERR DEF=0.5
```

PARAMETER NO.	GLOBAL	1	2	3	4	5
1	0.35404	1.000	0.207	0.040	0.030	0.312
2	0.25216	0.207	1.000	-0.065	-0.047	0.173
3	0.08553	0.040	-0.065	1.000	0.011	-0.001
4	0.06216	0.030	-0.047	0.011	1.000	0.007
5	0.33109	0.312	0.173	-0.001	0.007	1.000

Fit results (step 2)

- $g_{av} = \frac{F_1^A(0)}{F_1^V(0)}, g_w = \frac{F_2^V(0)}{F_1^V(0)}$

Parameter	$N_{sig}^{MC} = 10^6$ $N_{phsp}^{MC} = 10^7$	$N_{sig}^{MC} = 10^5$ $N_{phsp}^{MC} = 10^6$
α_Ψ	0.4581 ± 0.0097	0.4867 ± 0.0196
$\Delta\Phi$	0.7077 ± 0.0192	0.7252 ± 0.0398
$\alpha_{\bar{\Lambda}}$	-0.7595 ± 0.0081	-0.762 ± 0.0163
g_{av}	0.9839 ± 0.0122	0.9793 ± 0.024
g_w	1.3027 ± 0.0445	1.2793 ± 0.0878

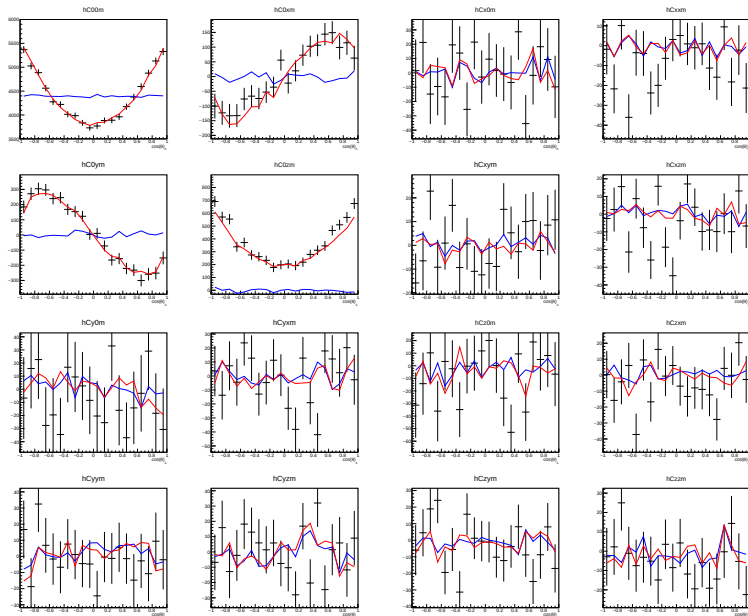
Parameter $N_{sig}^{MC} = 10^6$ $N_{phsp}^{MC} = 10^7$	e^- -mode	μ^- -mode
α_Ψ	0.4581 ± 0.0097	0.475 ± 0.003
$\Delta\Phi$	0.7077 ± 0.0192	0.6886 ± 0.0057
$\alpha_{\bar{\Lambda}}$	-0.7595 ± 0.0081	-0.7509 ± 0.0024
g_{av}	0.9839 ± 0.0122	1.1456 ± 0.0058
g_w	1.3027 ± 0.0445	1.6243 ± 0.0156

Fit method of moments (step 3) (1)

- Use fitted parameter values as input

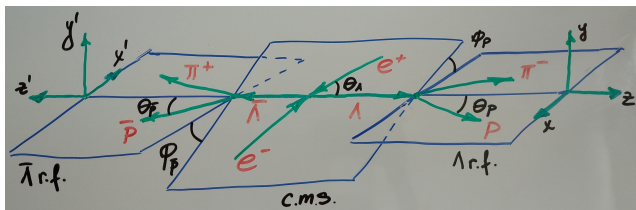
```
void momentsLLbar_seml(){  
  
    Double_t pp[5];  
    for( int i = 0; i < 5; i++ ) pp[i]=0.;  
  
    // nominal  
    pp[0] = 0.4581;    // alpha_J/psi  
    pp[1] = 0.7077;    // Delta_Phi  
    pp[2] = 0.9839;    // FA1_0/FV1_0 Lambda->p e- nu_ebar  
    pp[3] = 1.3027;    // FV2_0/FV1_0 Lambda->p e- nu_ebar  
    pp[4] = -0.7595;    // alpha_D LambdaBar->pbar pi+  
  
    AngDisLLbar angdis(pp[0],pp[1],pp[2],pp[3],pp[4]);  
}
```

Fit method of moments (step 3) (2)



Λ direction in c.m.s. (1)

- $e^+e^- \rightarrow \Lambda\bar{\Lambda} \rightarrow p\pi^-\bar{p}\pi^+$



- $\Lambda \rightarrow p\pi^-$ (c.m.s.)

- $x: \sin \theta_p \cos \phi_p$
- $y: \sin \theta_p \sin \phi_p$
- $z: \cos \theta_p$

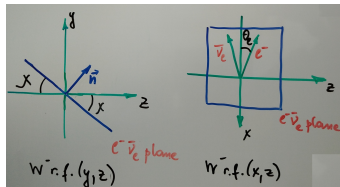
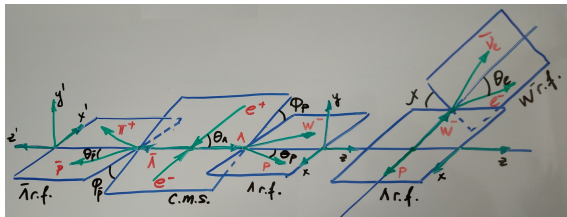
where $x = -x', y = y', z = -z'$

- $\bar{\Lambda} \rightarrow \bar{p}\pi^+$ (c.m.s.)

- $x': \sin \theta_{\bar{p}} \cos \phi_{\bar{p}}$
- $y': \sin \theta_{\bar{p}} \sin \phi_{\bar{p}}$
- $z': \cos \theta_{\bar{p}}$

Λ direction in c.m.s. (2)

• $e^+e^- \rightarrow \Lambda \bar{\Lambda} \rightarrow pl^- \bar{\nu}_l \bar{p} \pi^+$



• W^- rest frame

- $x: -\cos \theta_l$
- $y: \sin \theta_l \sin \chi$
- $z: \sin \theta_l \cos \chi$

• Λ rest frame

- $x: -\cos \theta_l \sin \theta_p$
- $y: \sin \theta_l \sin \chi$
- $z: \sin \theta_l \cos \chi \cos \theta_p$

• $\Lambda \rightarrow pl^- \bar{\nu}_l$ (c.m.s.)

- $x: -\cos \theta_l \sin \theta_p \cos \phi_p$
- $y: \sin \theta_l \sin \chi \sin \phi_p$
- $z: \sin \theta_l \cos \chi \cos \theta_p \cos \phi_p$

where $x = -x'$, $y = y'$, $z = -z'$

• $\bar{\Lambda} \rightarrow \bar{p} \pi^+$ (c.m.s.)

- $x': \sin \theta_{\bar{p}} \cos \phi_{\bar{p}}$
- $y': \sin \theta_{\bar{p}} \sin \phi_{\bar{p}}$
- $z': \cos \theta_{\bar{p}}$

$\Lambda\bar{\Lambda}$ direction in c.m.s. (3)

```
Double_t cosI(Int_t i,Int_t ip,Double_t z,Double_t ph)
{
    // calculate directional cos
    // i -> 0,x,y,z
    // ip -> 1,2 //Lambda,Lambda_bar
    if(i==0) return 1;
    if(i==2) return TMath::Sqrt(1-z*z)*TMath::Sin(ph);
    Double_t val = 0;
    if(i==1) val = TMath::Sqrt(1-z*z)*TMath::Cos(ph);
    if(i==3) val = z;
    if(ip==2) return -val;
    if(ip==1) return val;
    return 0;
}

Double_t cosI_semi(Int_t i,Int_t ip,Double_t z,Double_t ph,Double_t zl,Double_t phl)
{
    // calculate directional cos
    // i -> 0,x,y,z
    // ip -> 1,2 //Lambda,Lambda_bar
    if(i==0) return 1;
    if(i==2) return TMath::Sin(ph)*TMath::Sqrt(1-zl*zl)*TMath::Sin(phl);
    Double_t val = 0;
    if(i==1) val = -TMath::Sqrt(1-z*z)*TMath::Cos(ph)*zl;
    if(i==3) val = z*TMath::Cos(phl)*TMath::Sqrt(1-zl*zl);
    if(ip==2) return -val;
    if(ip==1) return val;
    return 0;
}
```

Some distributions (step 3)

