

Physics Beyond the Standard Model in Charmonium Sector

International Workshop on Physics at Future
High Intensity Collider @ 2-7 GeV in China



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Venue : Student Activity Center of USTC

Sponsors : University of Science and Technology, China

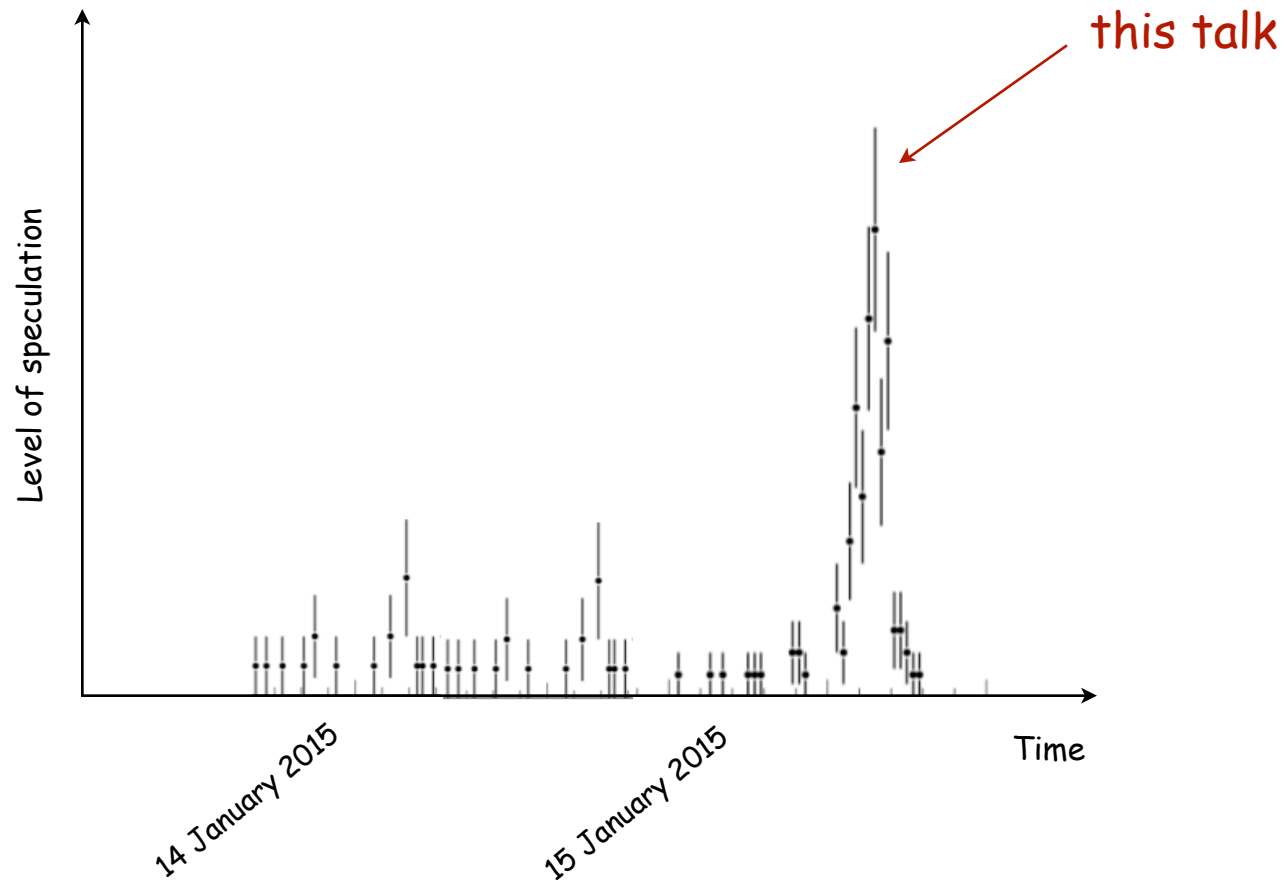
Collaboration Innovation Center for Particles and Interactions

State Key Laboratory of Particle Detection and Electronics

Detector



HIEPA-2015 presentations: level of speculation



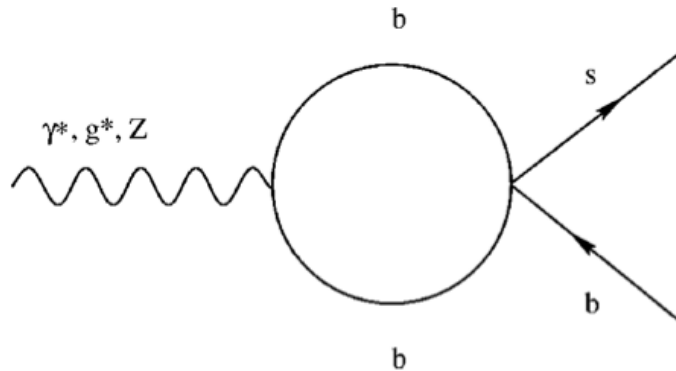
Introduction

- New physics can affect quarkonium decays
- ... but its contribution can be masked by QCD/QED processes
- We have to look at decays models where SM contribution is either small or under firm theoretical control
 - Processes with SM particles in the final state only ("heavy" NP)
 - quark FCNC transitions
 - lepton FCNC transitions
 - Processes with SM and NP particles in the final state ("light" NP)
 - light Dark Matter particles
 - dark photons
 - exotic light scalars, etc.

"Heavy" New Physics

Heavy New Physics

- FCNC processes can be probed in heavy quarkonium decays
 - FCNC in quark sector ($b \rightarrow s$): $\Upsilon \rightarrow BX_s$ or ($c \rightarrow u$): $J/\psi \rightarrow DX_u$



Datta, O'Donnell,
Pakvasa, Zhang (1999)

- This can be studied with a generic Lagrangian,

$$L_{new} = \frac{R_1}{\Lambda^2} \bar{s}(1 - \gamma^5)b \bar{b}(1 + \gamma^5)b + \frac{R_2}{\Lambda^2} \bar{s}(1 + \gamma^5)b \bar{b}(1 - \gamma^5)b$$

- ... or in particular models, i.e. FCNC Higgs, where
- DOPZ claim: $\text{Br}(J/\psi \rightarrow DX_u) \sim 10^{-5}$!

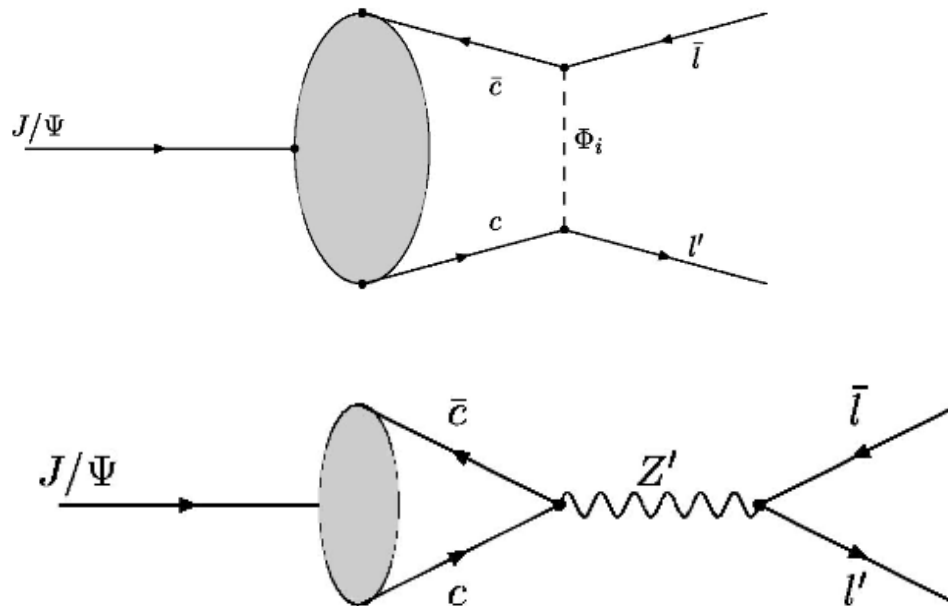
$$R_1 = \frac{1}{4} \frac{m_t^2}{f_\pi^2 m_H^2} |D_{Lbb}|^2 D_{Rbb} D_{Rbs}^*$$

$$R_2 = \frac{1}{4} \frac{m_t^2}{f_\pi^2 m_H^2} |D_{Rbb}|^2 D_{Lbb} D_{Lbs}^*$$

Heavy New Physics

- FCNC processes can be probed in heavy quarkonium decays
 - FCNC in lepton sector ($J/\psi \rightarrow e\mu$, etc.)

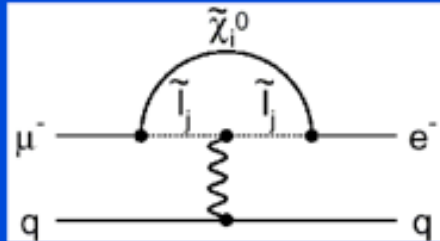
Hue, Yue, Feng (2003)



- Very much complimentary to muon conversion/FCNC tau decay experiments
- Recent BES constraint $\mathcal{B}(J/\psi \rightarrow e\mu) < 1.6 \times 10^{-7}$

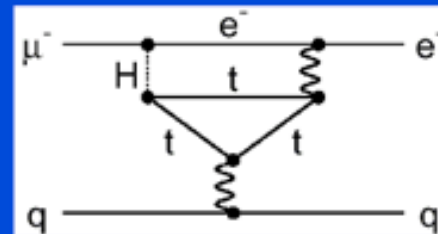
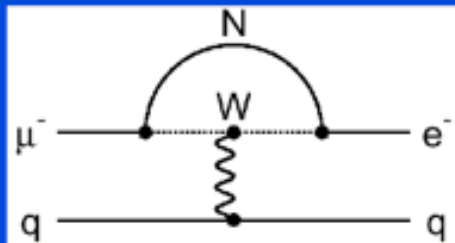
Sensitivity to Different Muon Conversion Mechanisms

Supersymmetry
Predictions at 10^{-15}



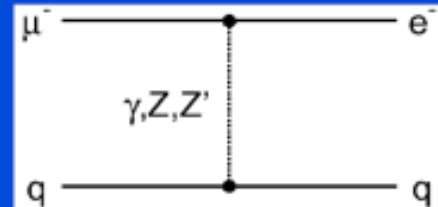
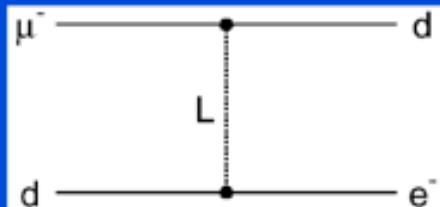
Compositeness
 $\Lambda_c = 3000 \text{ TeV}$

Heavy Neutrinos
 $|U_{\mu N}^* U_{eN}|^2 = 8 \times 10^{-13}$



Second Higgs
doublet
 $g_{H\mu e} = 10^{-4} \times g_{H\mu\mu}$

Leptoquarks
 $M_L = 3000 \sqrt{\lambda_{\mu d} \lambda_{e d}} \text{ TeV}/c^2$



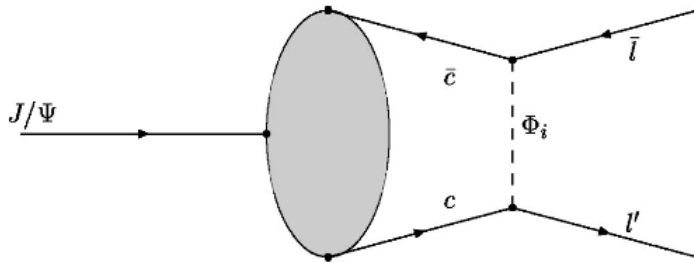
Heavy Z' ,
Anomalous Z
coupling
 $M_{Z'} = 3000 \text{ TeV}/c^2$
 $B(Z \rightarrow \mu e) < 10^{-17}$

After W. Marciano

James Miller, 2006

Heavy New Physics

- Example: leptoquarks



$$\begin{aligned} \text{Br}(J/\psi \rightarrow \mu \tau) &= \frac{9}{2^9 \pi^2 \alpha^2} m_\tau^2 M_{J/\psi}^2 \left(1 + 2 \frac{M_{J/\psi}^2}{m_\tau^2} \right) \\ &\times \left(1 - \frac{m_\tau^2}{M_{J/\psi}^2} \right)^2 \cdot \frac{|\lambda_L^{c\mu} \lambda_L^{c\tau}|^2 + |\lambda_R^{c\mu} \lambda_R^{c\tau}|^2}{M_\Phi^4} \\ &\times \text{Br}(J/\psi \rightarrow e^+ e^-). \end{aligned}$$

- From tau-decay:

$$\frac{|\lambda_L^{c\mu} \lambda_L^{c\tau}|^2 + |\lambda_R^{c\mu} \lambda_R^{c\tau}|^2}{M_\Phi^4} < 1.5 \times 10^{-10}.$$

- Prediction:

$$\text{Br}(J/\psi \rightarrow \mu \tau) < 3.0 \times 10^{-8}$$

Huo, Yue,
Feng (2003)

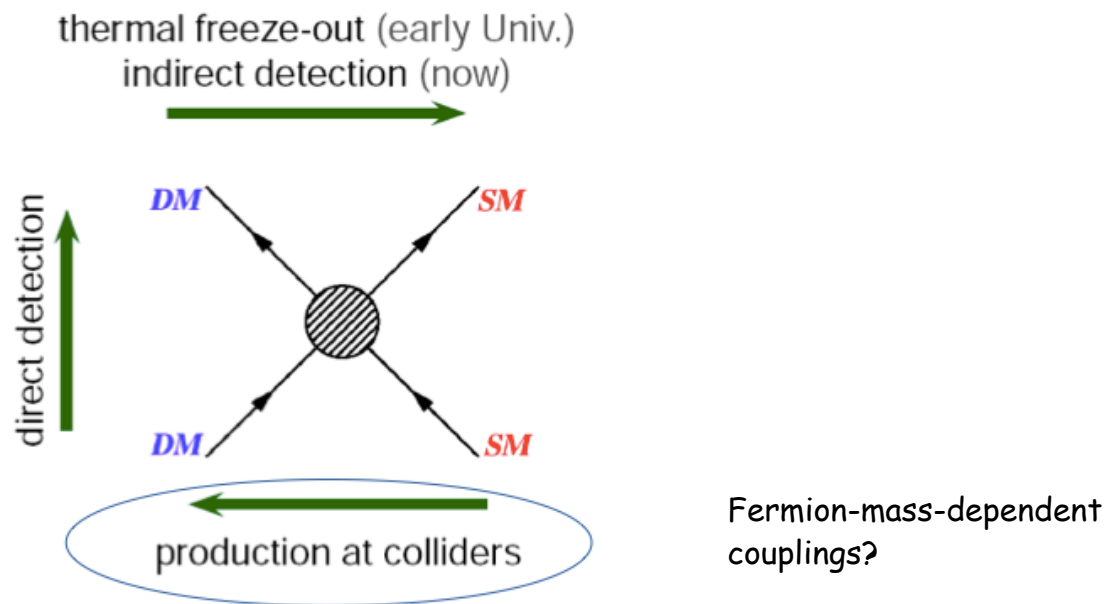
"Light" New Physics

Introduction: light Dark Matter

- DM does not have to be EW-scale (WIMP miracle)

$$\Omega_{DM} h^2 \sim \langle \sigma_{ann} v_{rel} \rangle^{-1} \propto \frac{M^2}{g^4}$$

- If DM is light, can use quarkonium decays to probe its properties



Backgrounds: decays to neutrino pairs

- Any studies of [cc] → missing energy + X would have SM background

- Invisible decays of heavy quarkonium (BR)

Chang, Lebedev, Ng (1998);
Yeghiyan (2009)

$$B(Y(1S) \rightarrow \nu\bar{\nu}) = \frac{\Gamma(Y(1S) \rightarrow \nu\bar{\nu})}{\Gamma_{Y(1S)}}$$

$$= \frac{N_\nu G_F^2}{48\pi} \left(1 - \frac{4}{3} \sin^2 \theta_W\right)^2 \frac{f_{Y(1S)}^2 M_{Y(1S)}^3}{\Gamma_{Y(1S)}}$$

Theory: $B(Y(1S) \rightarrow \nu\bar{\nu}) = 9.85 \times 10^{-6}$, Experiment: $B(Y(1S) \rightarrow \text{invisible}) < 3.0 \times 10^{-4}$,
 $B(J/\Psi \rightarrow \nu\bar{\nu}) = 2.70 \times 10^{-8}$. $B(J/\Psi \rightarrow \text{invisible}) < 7.2 \times 10^{-4}$.

BaBar (2009),
BES (2008)

- Decays with photons and missing energy (BR & γ -spectrum)

$$B(Y(3S) \rightarrow \nu\bar{\nu}\gamma) = \frac{N_\nu G_F^2}{243\pi} \frac{\alpha}{4\pi} \frac{f_{Y(3S)}^2 M_{Y(3S)}^3}{\Gamma_{Y(3S)}}$$

Theory: $B(Y(3S) \rightarrow \nu\bar{\nu}\gamma) = (3.14_{-0.32}^{+0.38}) \times 10^{-9}$
 Experiment: $B(Y(3S) \rightarrow \gamma + \text{invisible}) < (0.7-31) \times 10^{-6}$

Currently, SM "physics" background is not an issue.

Low-mass New Physics: light Dark Matter

➤ Effective Lagrangian analysis of low-mass Dark Matter

- Leading-order operators for DM of different spin for 1^{--} decays

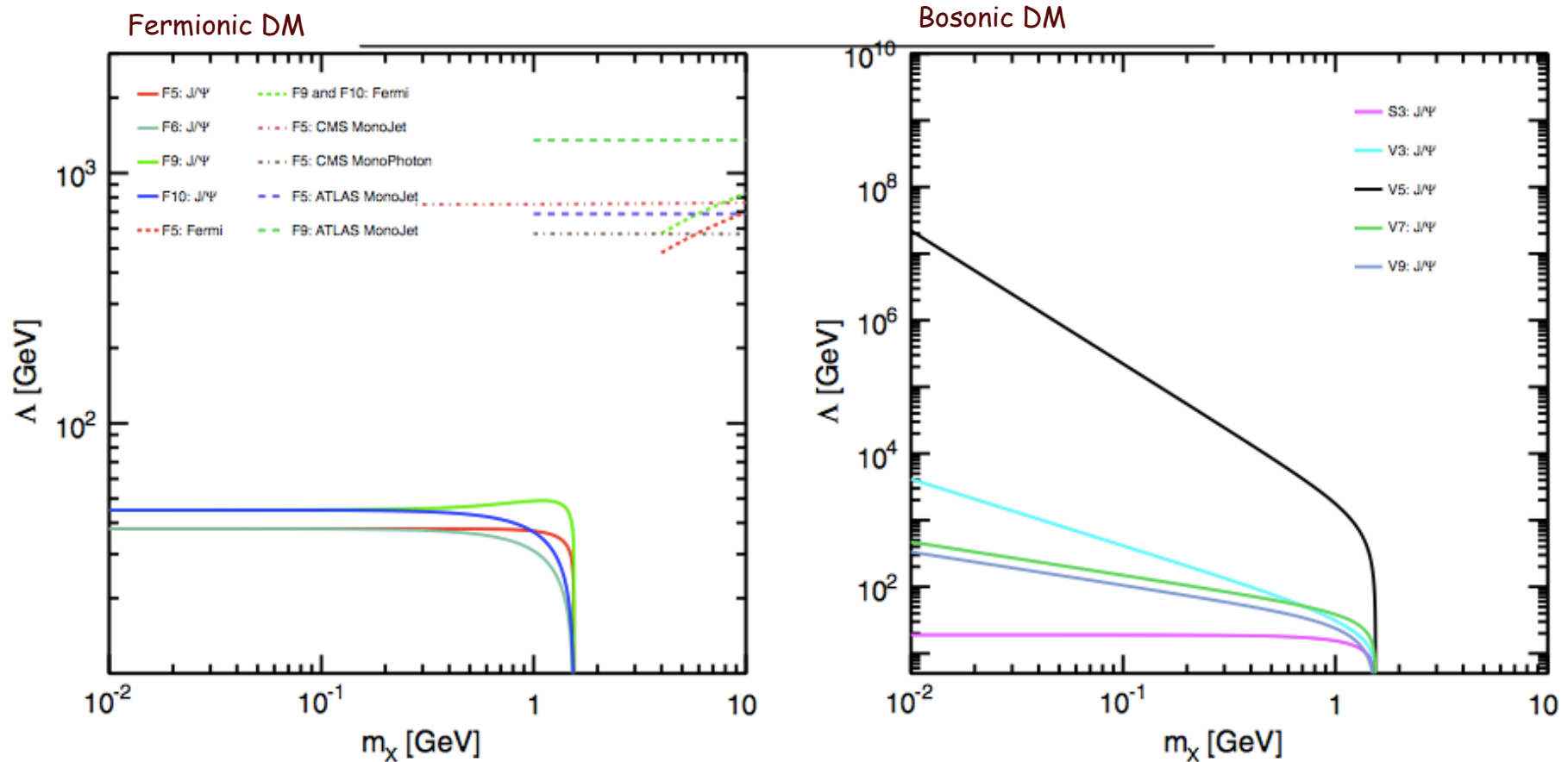
Name	Interaction structure	Annihilation	Scattering
F5	$(1/\Lambda^2)\bar{X}\gamma^\mu X\bar{q}\gamma_\mu q$	Yes	SI
F6	$(1/\Lambda^2)\bar{X}\gamma^\mu\gamma^5 X\bar{q}\gamma_\mu q$	No	No
F9	$(1/\Lambda^2)\bar{X}\sigma^{\mu\nu} X\bar{q}\sigma_{\mu\nu} q$	Yes	SD
F10	$(1/\Lambda^2)\bar{X}\sigma^{\mu\nu}\gamma^5 X\bar{q}\sigma_{\mu\nu} q$	Yes	No
S3	$(1/\Lambda^2)i\text{Im}(\phi^\dagger\partial_\mu\phi)\bar{q}\gamma^\mu q$	No	SI
V3	$(1/\Lambda^2)i\text{Im}(B_\nu^\dagger\partial_\mu B^\nu)\bar{q}\gamma^\mu q$	No	SI
V5	$(1/\Lambda)(B_\mu^\dagger B_\nu - B_\nu^\dagger B_\mu)\bar{q}\sigma^{\mu\nu} q$	No	SD
V7	$(1/\Lambda^2)B_\nu^{(\dagger)}\partial^\nu B_\mu\bar{q}\gamma^\mu q$	No	No
V9	$(1/\Lambda^2)\epsilon^{\mu\nu\rho\sigma}B_\nu^{(\dagger)}\partial_\rho B_\sigma\bar{q}\gamma_\mu q$	No	No

Fernandez, Kumar,
Seong, Stengel (2014)

- These operators are relevant only for invisible decays of quarkonia
- Note that Wilson coefficients are redefined in the scale Λ

Light Dark Matter I

- Constraints on NP scale Λ from J/ψ decays (use dipion decays of ψ'')



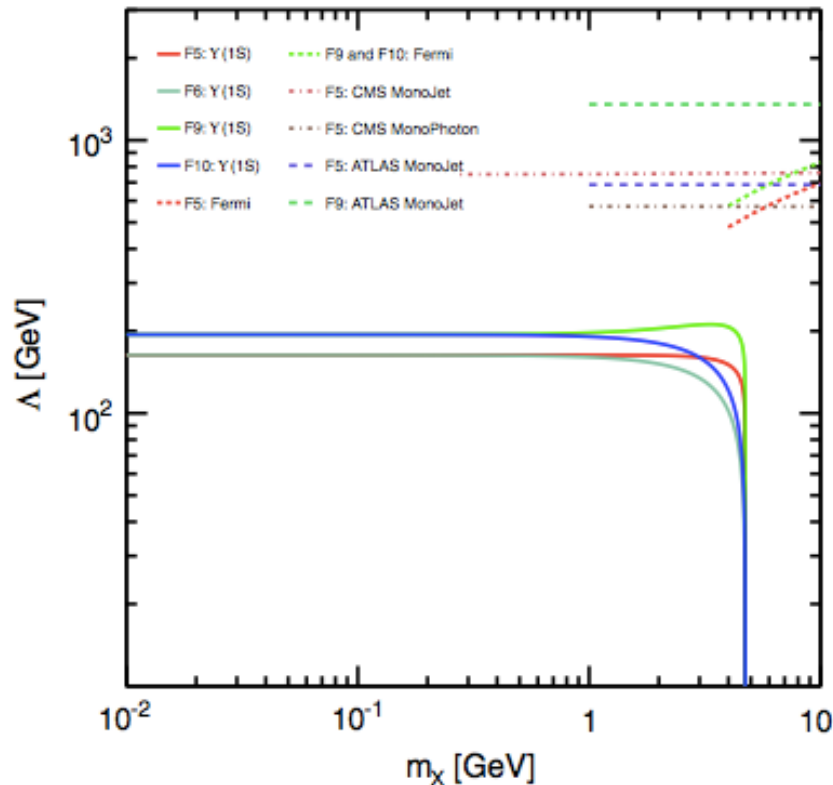
- Higher luminosities will provide more stringent constraints

Fernandez, Kumar,
Seong, Stengel (2014)

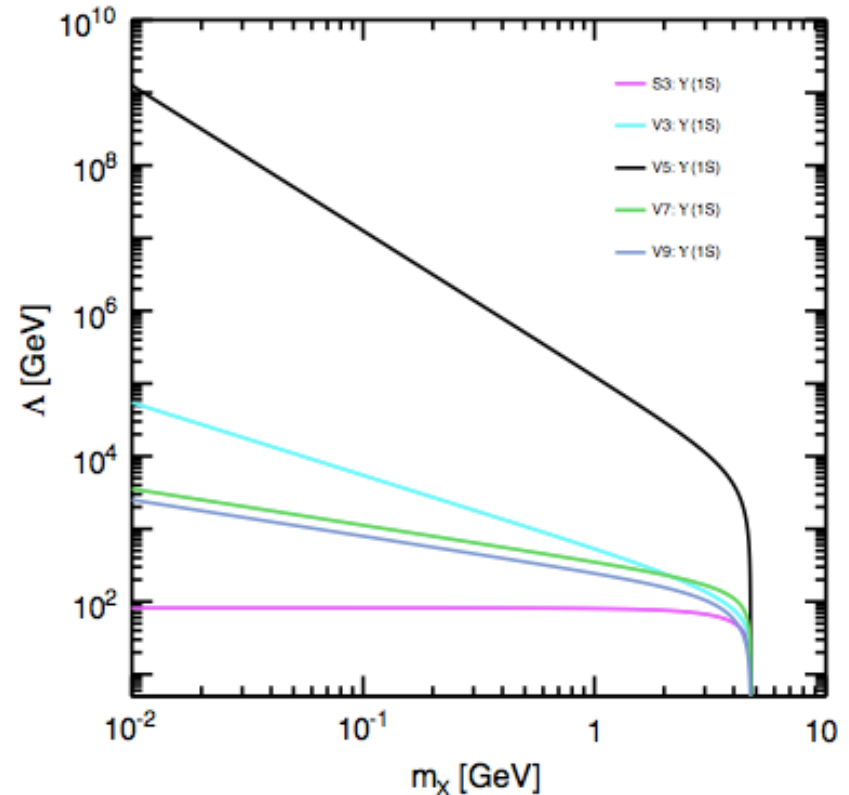
Light Dark Matter I

- Constraints on NP scale Λ from Υ decays (use dipion decays of $\Upsilon(3S)$)

Fermionic DM



Bosonic DM

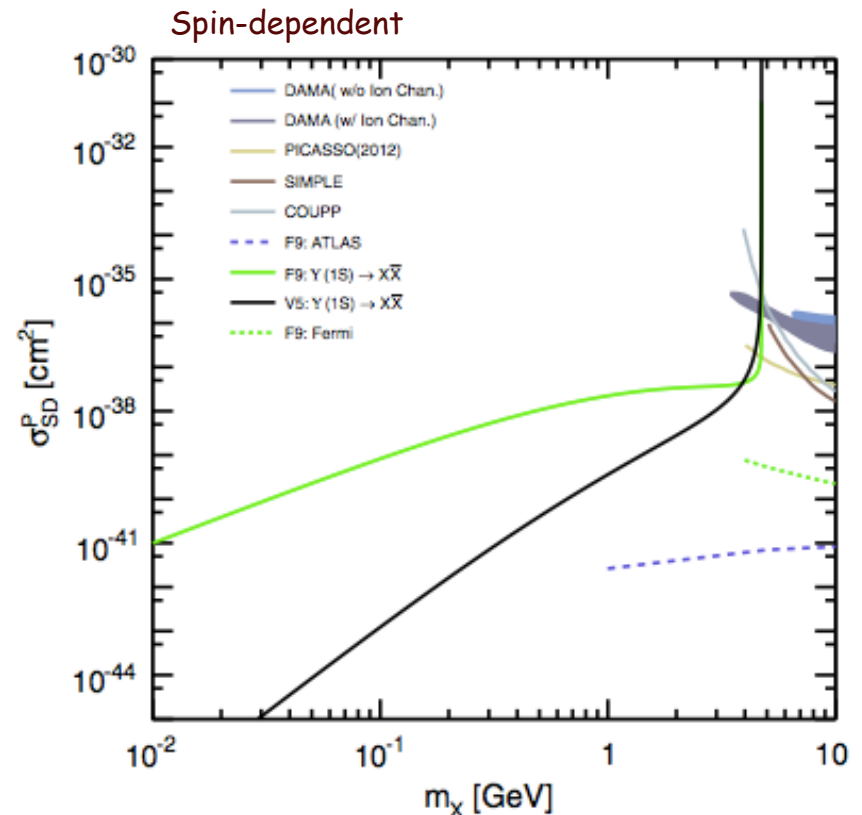
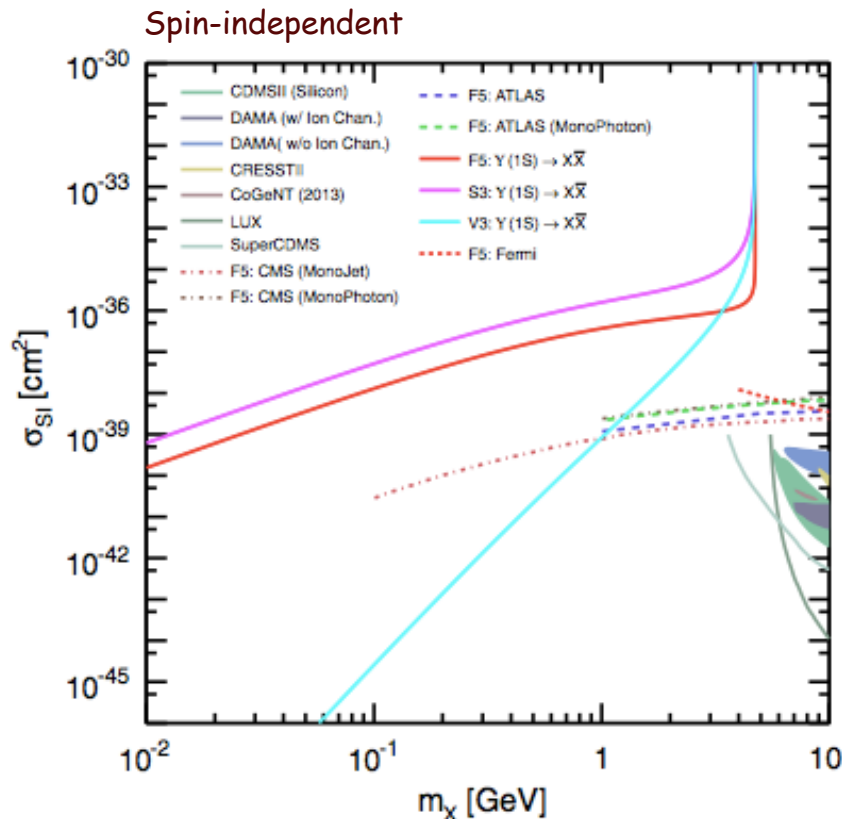


- Higher luminosities will provide more stringent constraints

Fernandez, Kumar,
Seong, Stengel (2014)

Light Dark Matter I

➤ Constraints on NP scale Λ from direct detection



- Complimentary constraints if coupling universality is assumed

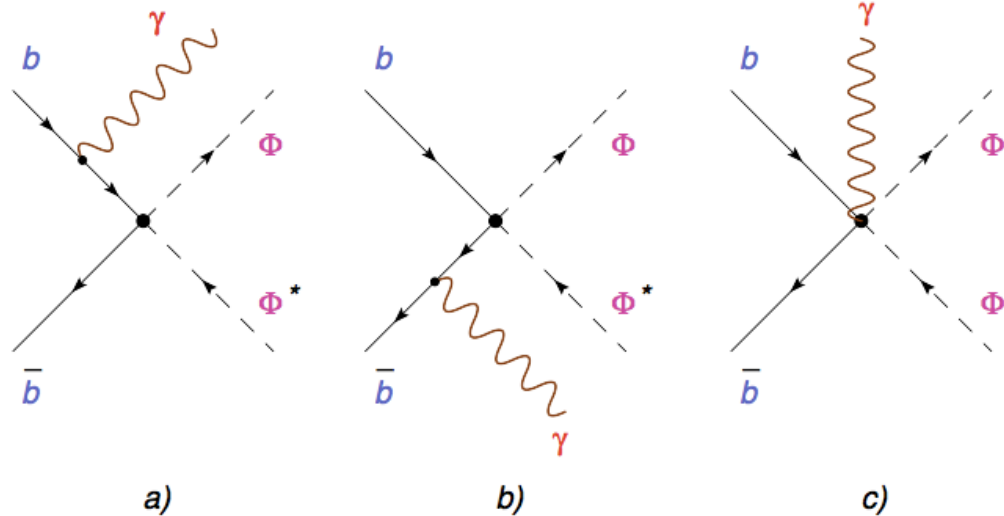
Fernandez, Kumar,
Seong, Stengel (2014)

Light Dark Matter II

➤ Effective Lagrangian analysis of low-mass Dark Matter

- New operators would appear for $\Upsilon(nS) \rightarrow \gamma + \text{missing energy}$

Yeghiyan (2009)



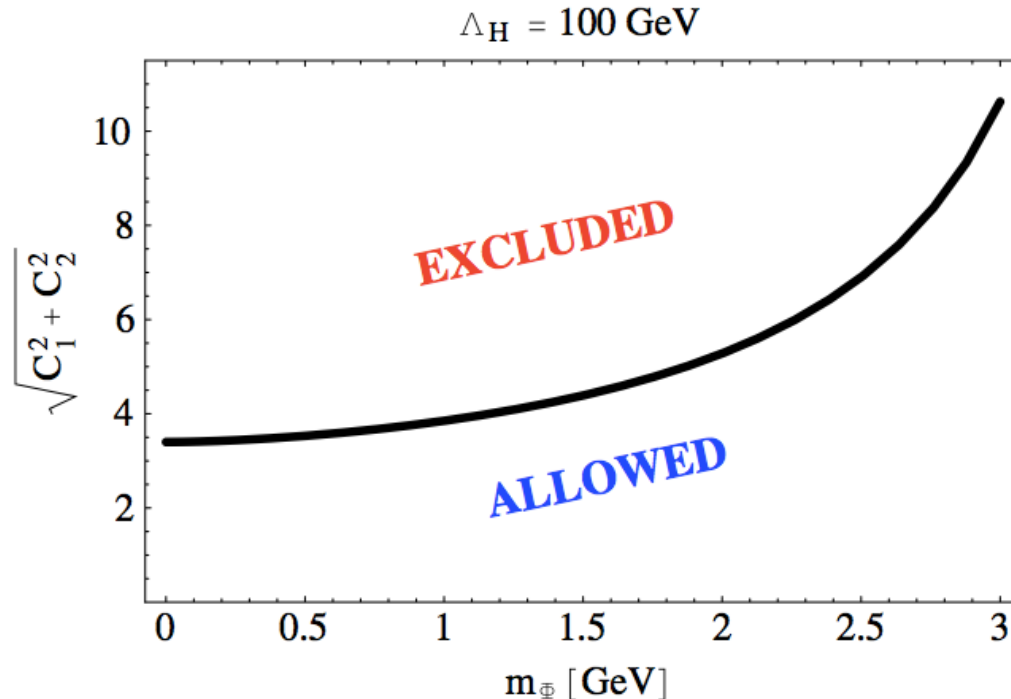
- For scalar DM only two new operators contribute, $H_{\text{eff}} = \frac{2}{\Lambda_H^2} \sum_i C_i O_i$

$$O_1 = m_b(\bar{b}b)(\Phi^*\Phi), \quad O_2 = im_b(\bar{b}\gamma_5 b)(\Phi^*\Phi),$$

$$O_3 = (\bar{b}\gamma^\mu b)(\Phi^*i\vec{\partial}_\mu\Phi), \quad O_4 = (\bar{b}\gamma^\mu\gamma_5 b)(\Phi^*i\vec{\partial}_\mu\Phi),$$

Light Dark Matter II

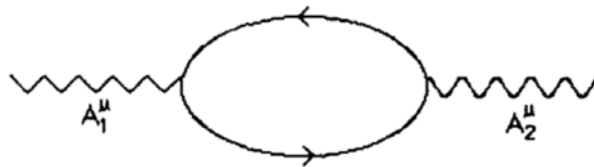
- Constraints on Wilson coefficients from Υ (ns) decays



Yeghiyan (2009)

- Higher luminosities will provide more stringent constraints
- Easy to provide specific models for DM, e.g. $-\mathcal{L} = \frac{m_0^2}{2} \Phi^2 + \lambda_1 \Phi^2 |H_1|^2 + \lambda_2 \Phi^2 |H_2|^2 + \lambda_3 \Phi^2 (H_1 H_2 + \text{H.c.})$
- ... in which case, $|\lambda_3| < \left(\frac{17.4}{\tan\beta}\right) \left(\frac{m_{H^0}}{160 \text{ GeV}}\right)^2 f^{-1/2}(x_\Phi)$.

Models of light Dark Matter and friends

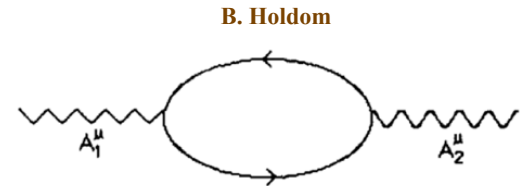


Low-mass New Physics: “dark photons”

➤ “Dark photons:” extra U(1) symmetries

- Assume there is an extra U(1) state V kinetically mixed with photons

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}V_{\mu\nu}V^{\mu\nu} - \frac{\kappa}{2}V_{\mu\nu}F^{\mu\nu} + \frac{m_V^2}{2}V_\mu V^\mu + \mathcal{L}_{h'}$$



Note: no direct couplings of V with fermions!

- Redefine (diagonalize) vector states...

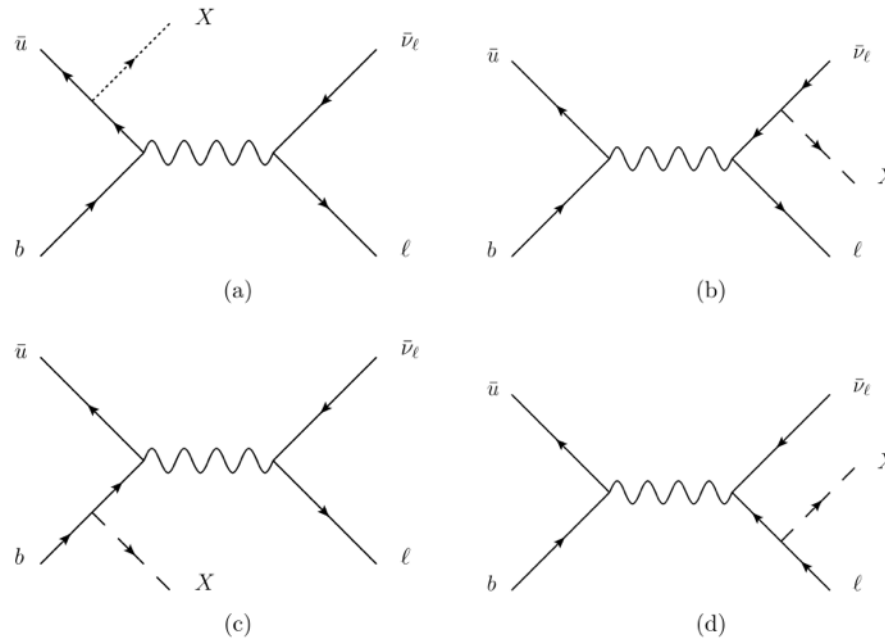
$$A \rightarrow A' - \frac{\kappa}{\sqrt{1-\kappa^2}}V', \quad V \rightarrow \frac{1}{\sqrt{1-\kappa^2}}V'$$

- ... to get a direct coupling of V' with fermion currents!

$$\mathcal{L}_f = -eQ_f A'_\mu \bar{\psi}_f \gamma^\mu \psi_f - \frac{\kappa e Q_f}{\sqrt{1-\kappa^2}} V'_\mu \bar{\psi}_f \gamma^\mu \psi_f$$

Low-mass New Physics: “dark photons”

- “Dark photons:” masses in the eV-MeV range
 - low-energy states (mesons/baryons) can decay into “dark photons”

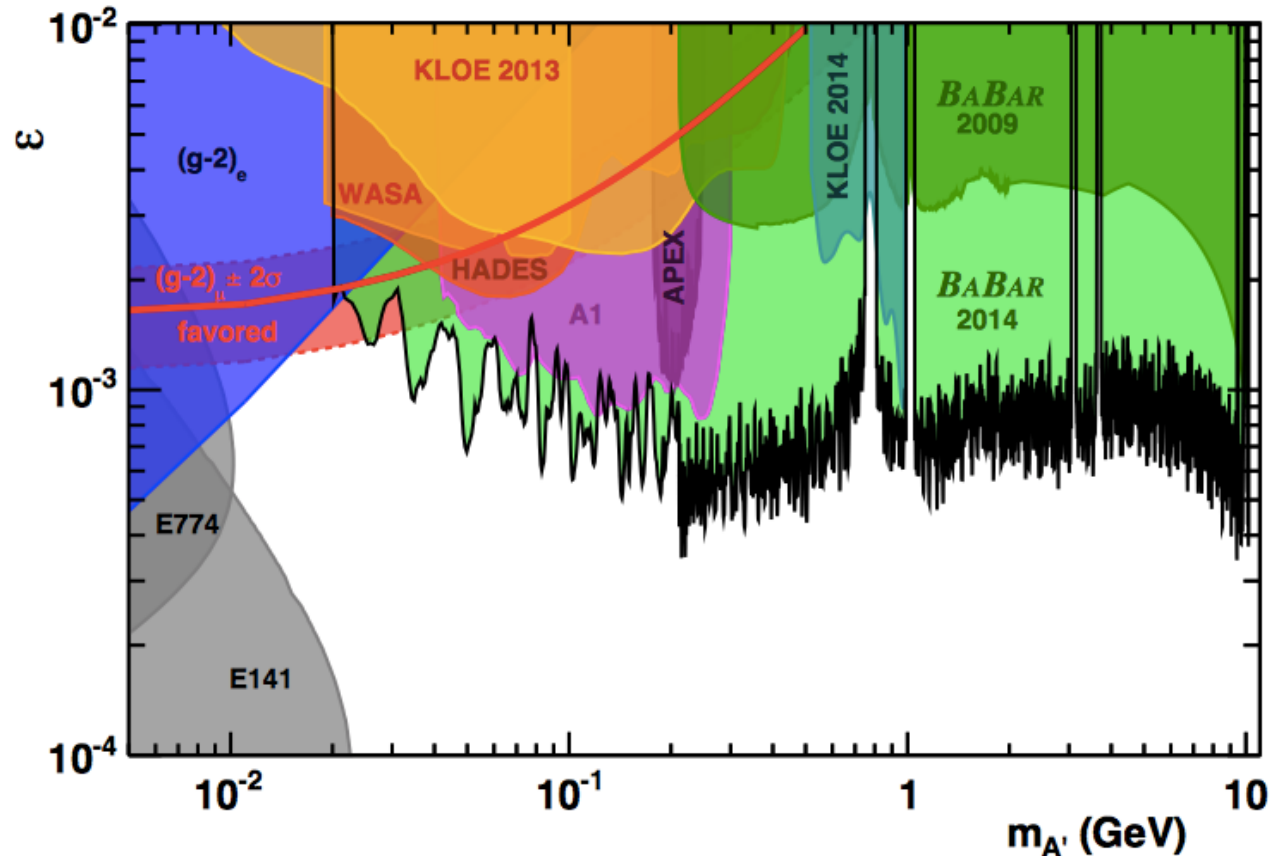


YG. Aditya,
K. Healey,
AAP

- Can be seen as “missing energy” in particle decays (e.g. $J/\psi \rightarrow \gamma\gamma + ME$ or $J/\psi \rightarrow ME + \text{light quark meson(s)}$)
- ... or in the lepton modes of V' decays into ee , etc.

Low-mass New Physics: “dark photons”

- BaBar recent search in $e^+e^- \rightarrow \gamma A', A' \rightarrow l^+l^-$ ($l = e, \mu$)

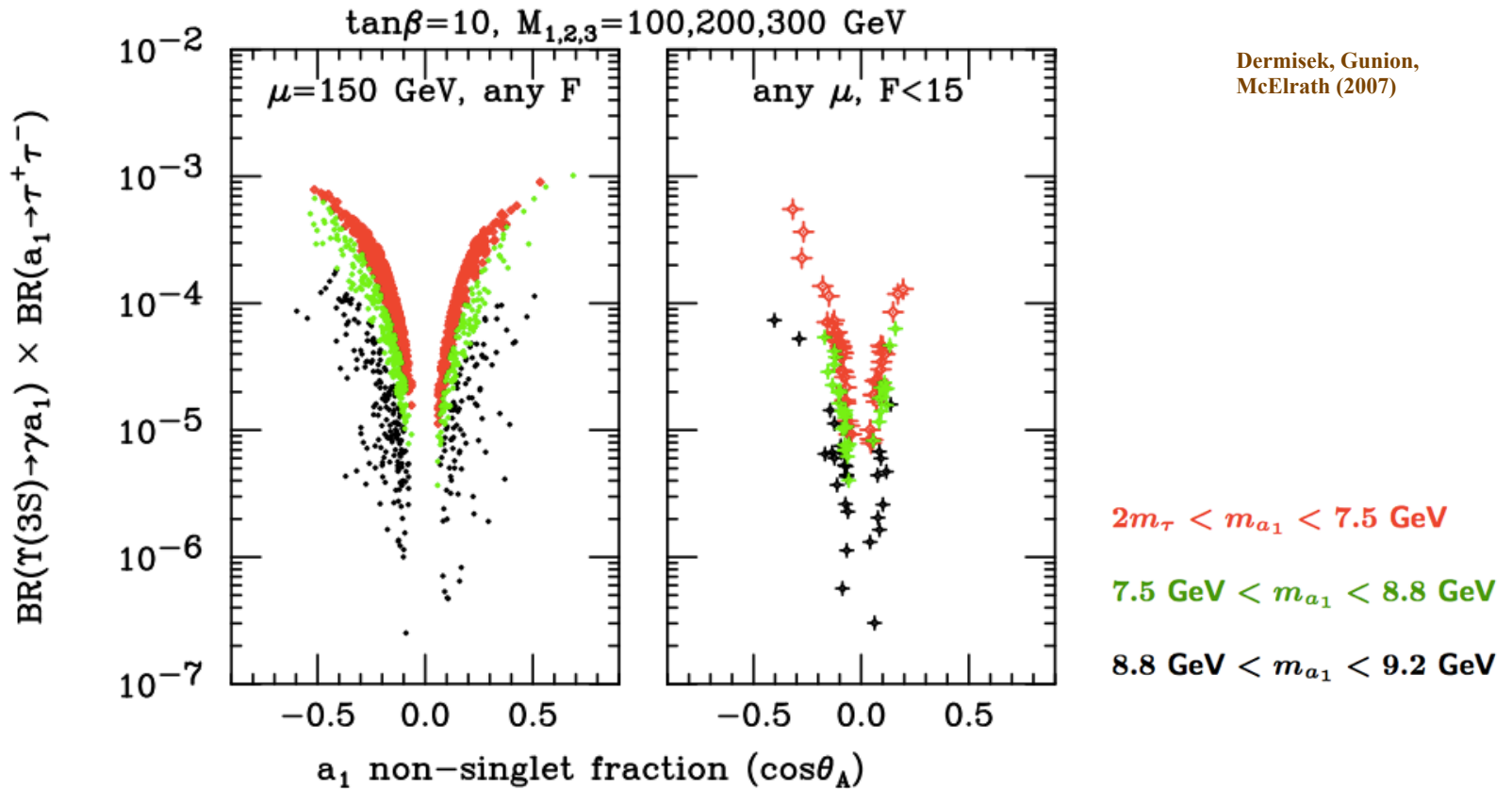


BaBar (2014)
PRL 113 (2014) 20, 201801

- can recast this search into search for low-mass BSM Higgs ($m \sim \text{several GeV}$)

Low-mass New Physics: pseudoscalar Higgs

- NMSSM: light scalar in $e^+e^- \rightarrow \gamma A', A' \rightarrow l^+l^-$ ($l = e, \mu$)



$$\mathcal{B}(\Upsilon(3S) \rightarrow \gamma A^0) \times \mathcal{B}(A^0 \rightarrow \tau^+ \tau^-) < (1.5 - 16) \times 10^{-5} \text{ at 90\% C.L.}$$

R. Godang (QWG13)

Low-mass New Physics: "axion-like" states

➤ "Axion-like" states

- Coupling similar to axion interactions (but can couple to leptons as well)

$$\mathcal{L}_a = -\frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \gamma_5 \psi + \frac{C_\gamma}{f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

- ... can be rewritten with

$$\mathcal{L}_a = -\left(\frac{1}{f_a} + \frac{4\pi C_\gamma}{f_a \alpha}\right) \partial_\mu a \bar{\psi} \gamma^\mu \gamma_5 \psi - im_\psi \left(\frac{8\pi C_\gamma}{f_a \alpha}\right) a \bar{\psi} \gamma_5 \psi$$

➤ New constraints from $J/\psi \rightarrow \gamma\gamma + \text{ME}$ or $J/\psi \rightarrow \text{ME} + \text{light quark meson(s)}$?

Conclusions

- New Physics particles can be searched for in quarkonium decays
 - SM final states: heavy new physics
 - SM+NP final states: light NP particles
- Interesting reach for a variety of models
 - complementarity to muon conversion experiments
 - complementarity to searches for Dark Matter
- Possible observations of exotic modes
 - small decay branching ratios...
 - ... but some models predict NP reach in the next round of experiments

7th International Workshop on Charm Physics
May 17-23, 2015
Wayne State University
Detroit, MI

<http://charm2015.wayne.edu>

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